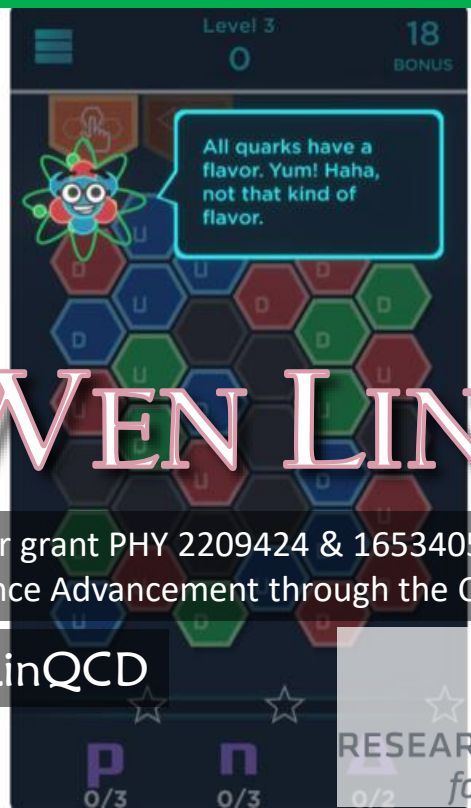


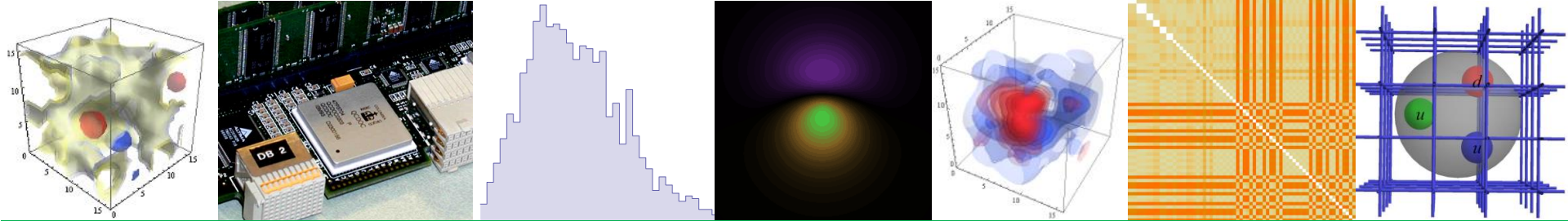
# QCD on the Lattice - II



HUEY-WEN LIN

This work of HL is supported by the NSF under grant PHY 2209424 & 1653405, DOE under DE-SC0024053 and the Research Corporation for Science Advancement through the Cottrell Scholar Award

 @LinQCD



# Probing the Heart of Matter with Supercomputers

**QUANTUM 3**

PLAY

NSF

U.S. DEPARTMENT OF ENERGY

Level 3  
3,000  
16 BONUS

u d

p 0/3 n 1/3 Δ 0/2

Level 3  
0  
18 BONUS

All quarks have a flavor. Yum! Haha, not that kind of flavor.

p 0/3 n 0/3

Level 8  
24,000  
11 BONUS

CTEQ

RESEARCH CORPORATION for SCIENCE ADVANCEMENT

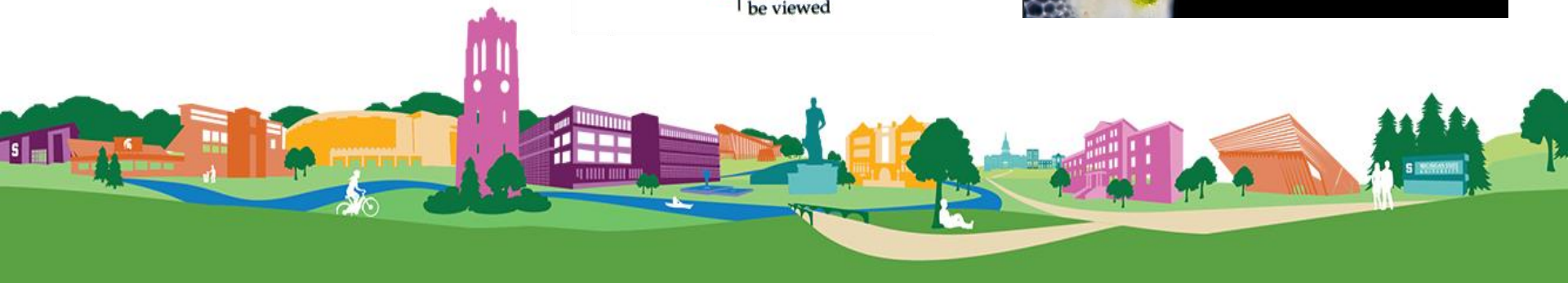
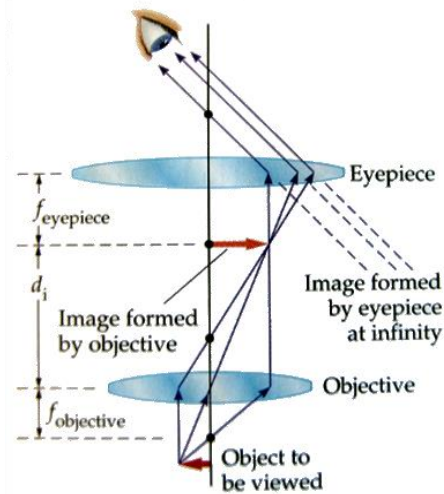
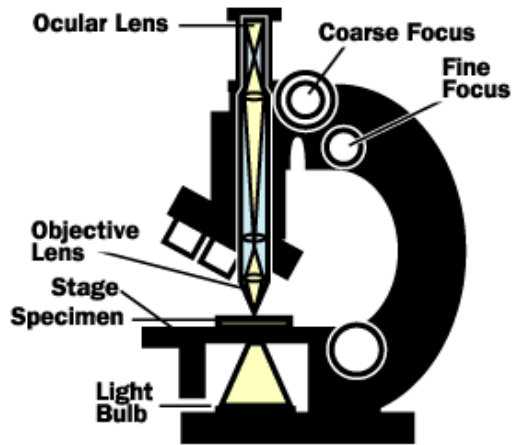
HUEY-WEN LIN

This work of HL is supported by the NSF under grant PHY 2209424 & 1653405, DOE under DE-SC0024053 and the Research Corporation for Science Advancement through the Cottrell Scholar Award

@LinQCD

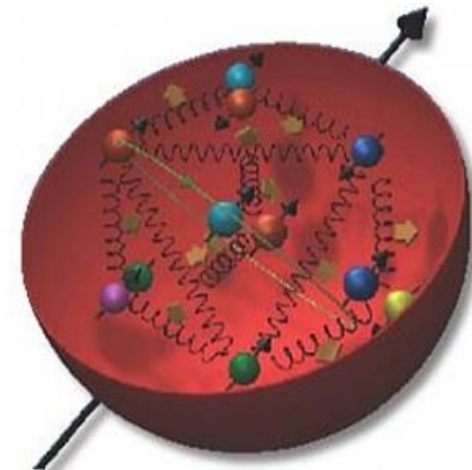
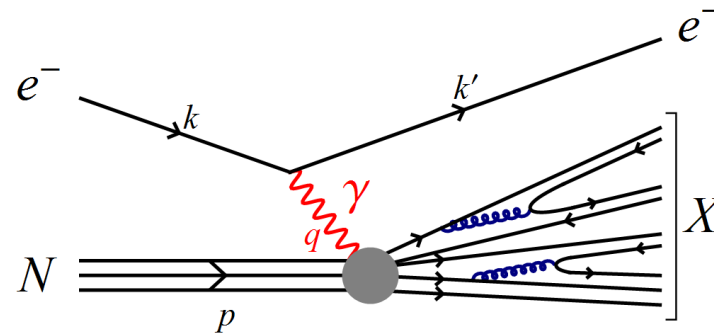
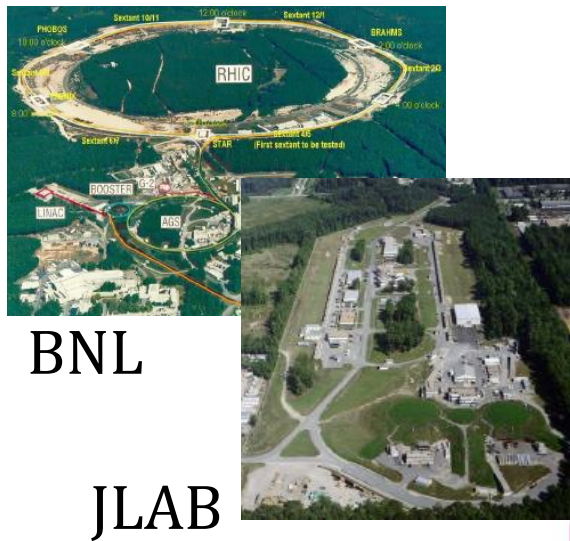
# Topic in QCD: Structure

∞ What is the structure of the nucleon?  
probing insights into nucleons

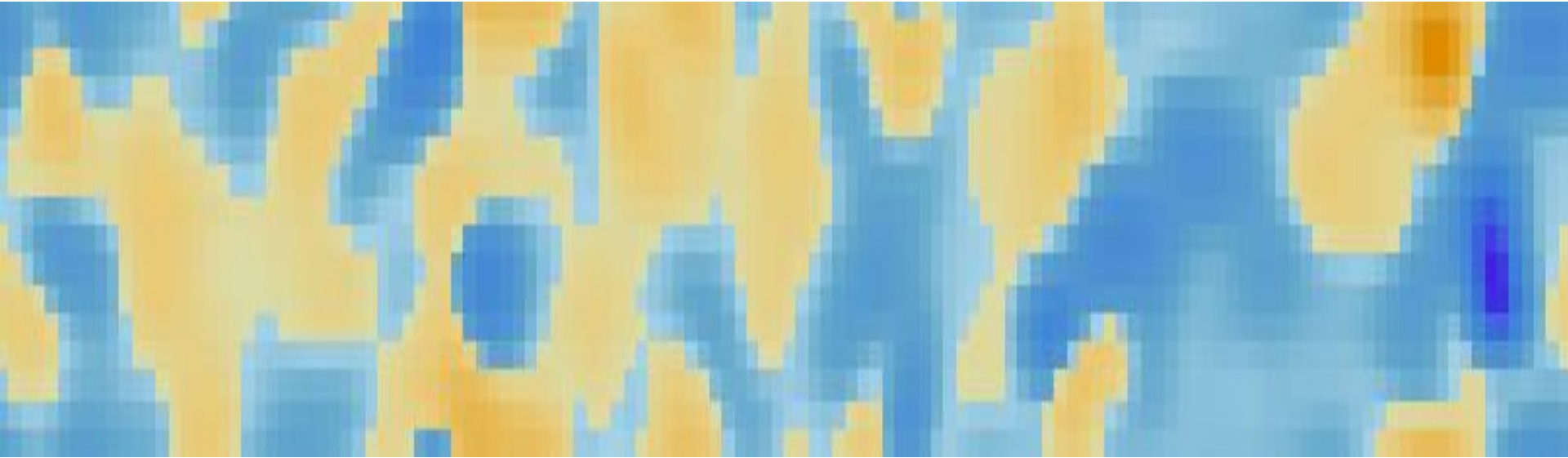


# Topic in QCD: Structure

∞ What is the structure of the nucleon?  
probing insights into nucleons



# Nucleon Matrix Elements

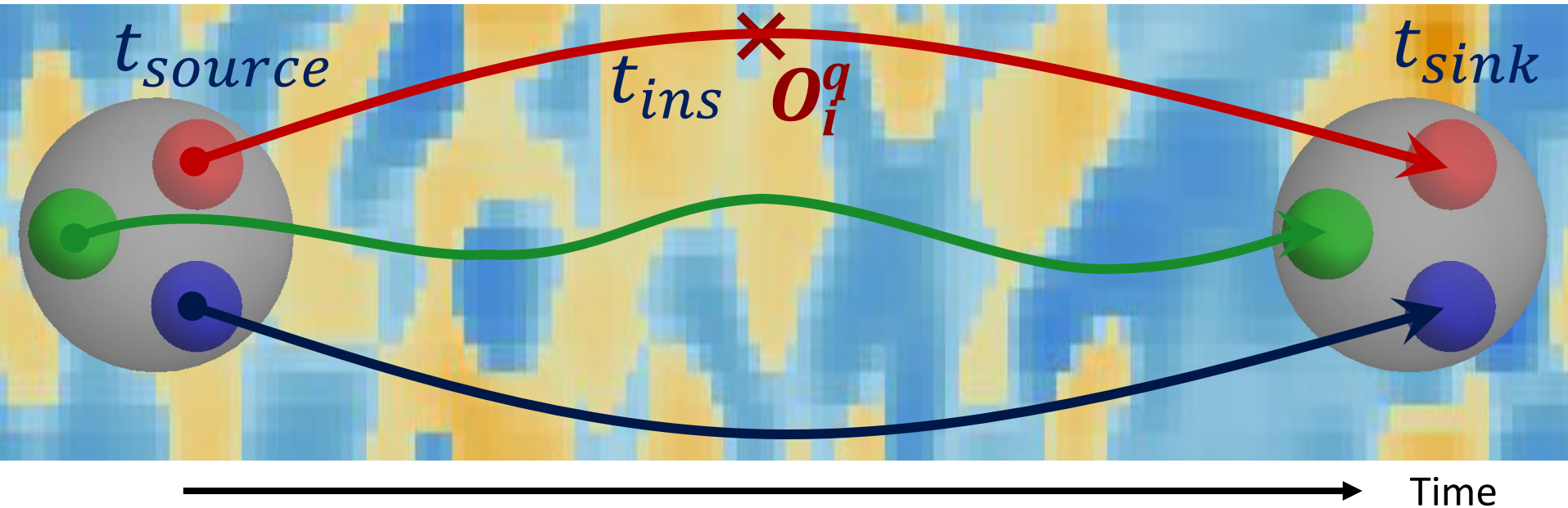


§ Pick a QCD vacuum

↪ Gauge/fermion actions, flavor (2, 2+1, 2+1+1),  $m_\pi$ ,  $a$ ,  $L$ , ...

# Nucleon Matrix Elements

Lattice-QCD calculation of  $\langle N | \bar{q} \Gamma q | N \rangle$



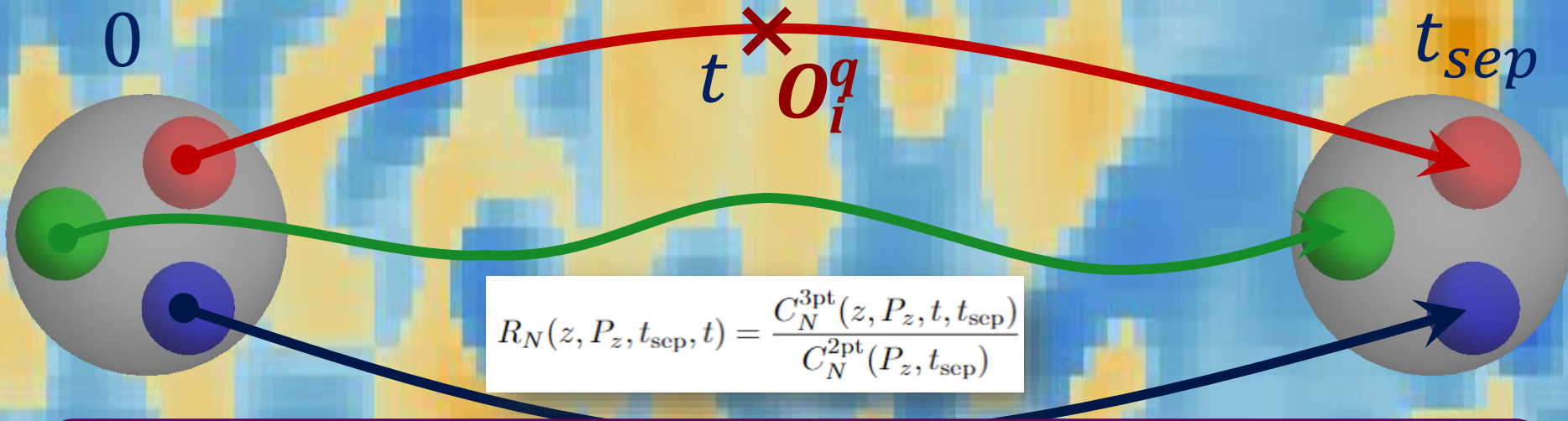
§ Construct correlators (hadronic observables)

⌘ Requires “quark propagator”

Invert Dirac-operator matrix (rank  $O(10^{12})$ )

# Nucleon Matrix Elements

Exercise: 3) Derive the time-dependent formula for the three-point correlator. Start with the pion case.



§ Careful analysis needed to remove systematics

⚡ Wrong results if **excited-state systematic** is not under control

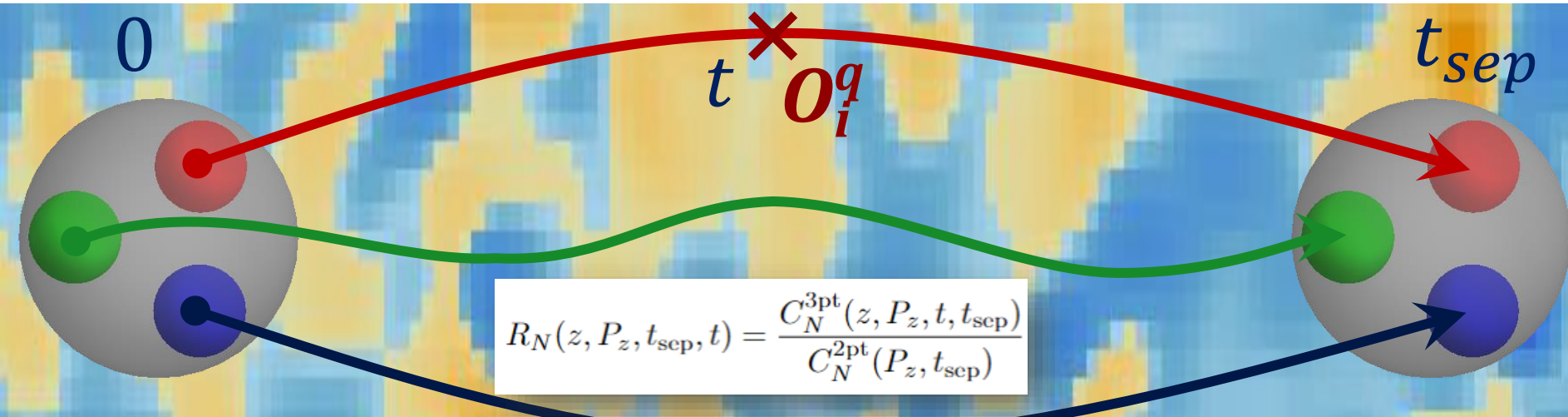
$$C_N^{3\text{pt}}(z, P_z, t, t_{\text{sep}}) \propto |A_{N,0}|^2 \langle 0|O_g|0\rangle e^{-E_{N,0}t_{\text{sep}}} + |A_{N,0}||A_{N,1}|\langle 0|O_g|1\rangle e^{-E_{N,1}(t_{\text{sep}}-t)} e^{-E_{N,0}t}$$

$$+ |A_{N,0}||A_{N,1}|\langle 1|O_g|0\rangle e^{-E_{N,0}(t_{\text{sep}}-t)} e^{-E_{N,1}t} + |A_{N,1}|^2 \langle 1|O_g|1\rangle e^{-E_{N,1}t_{\text{sep}}}$$

$$C_N^{2\text{pt}}(P_z, t) \propto |A_{N,0}|^2 e^{-E_{N,0}t} + |A_{N,1}|^2 e^{-E_{N,1}t} + \dots,$$

# Nucleon Matrix Elements

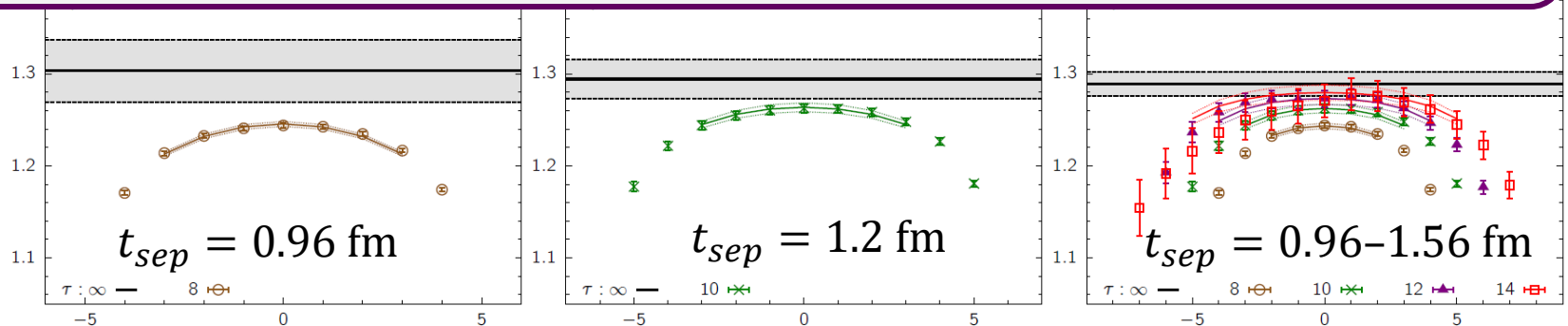
Lattice-QCD calculation of  $\langle N | \bar{q} \Gamma q | N \rangle$



§ Careful analysis needed to remove systematics

⚡ Wrong results if **excited-state systematic** is not under control

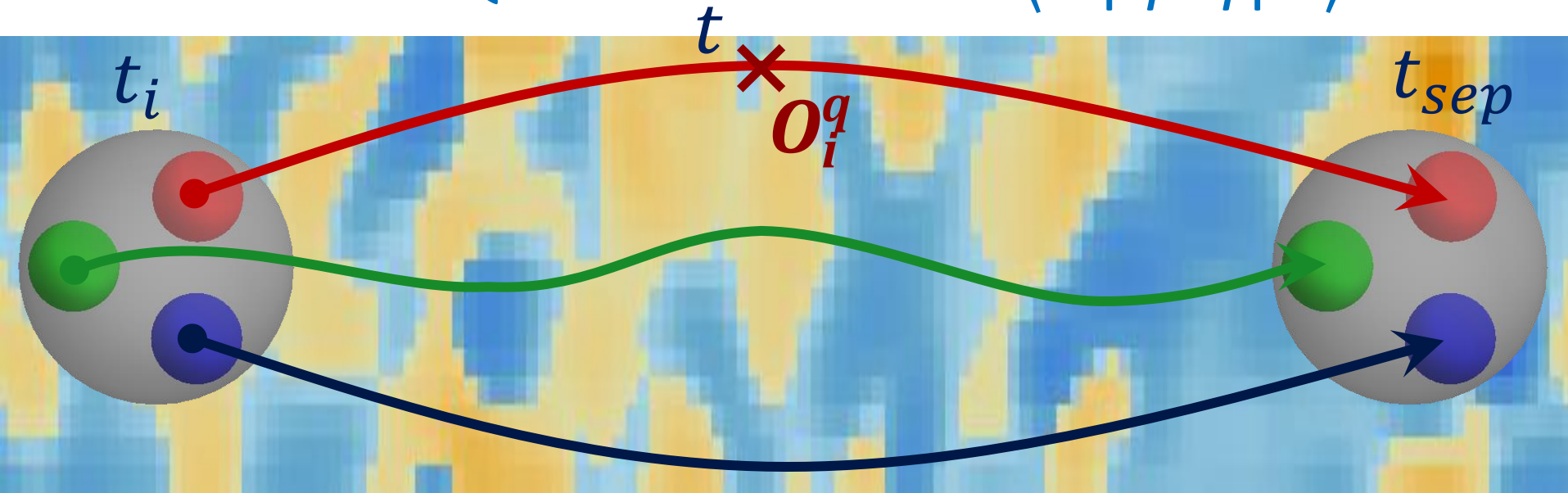
Ratio Plot





# Nucleon Matrix Elements

Lattice-QCD calculation of  $\langle N | \bar{q} \Gamma q | N \rangle$



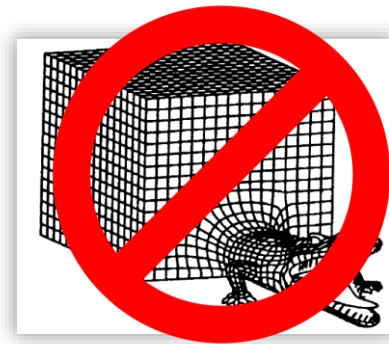
§ Systematic uncertainty (nonzero  $a$ , finite  $L$ , etc.)

∞ Nonperturbative renormalization

e.g. RI/SMOM scheme in  $\overline{\text{MS}}$  at 2 GeV

∞ Extrapolation to the continuum limit

$(m_\pi \rightarrow m_\pi^{\text{phys}}, L \rightarrow \infty, a \rightarrow 0)$

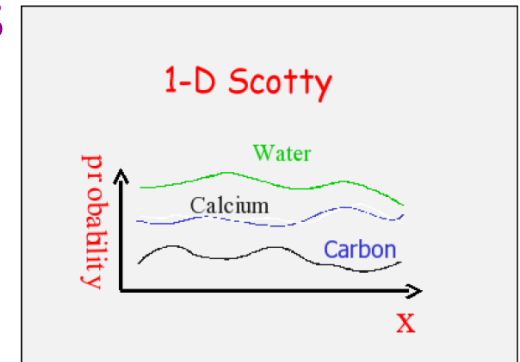


# Hadron Structure

## § Structure function/distribution functions

∞ deep inelastic scattering

∞  $\langle x^n \rangle_q$ ,  $\langle x^n \rangle_{\Delta q}$ ,  $\langle x^n \rangle_{\delta q}$



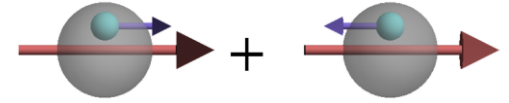
$$\mathcal{O}_{V^a}^{\mu\nu} = \bar{q} \gamma^{\{\mu} \overleftrightarrow{D}^{\nu\}} \tau^a q,$$

$$\mathcal{O}_{A^a}^{\mu\nu} = \bar{q} \gamma^{\{\mu} \overleftrightarrow{D}^{\nu\}} \gamma^5 \tau^a q,$$

$$\mathcal{O}_{T^a}^{\mu\nu\rho} = \bar{q} \sigma^{[\mu} \{\nu\} \overleftrightarrow{D}^{\rho\}} \tau^a q,$$

# Moments of PDFs

§ First moments are most commonly done

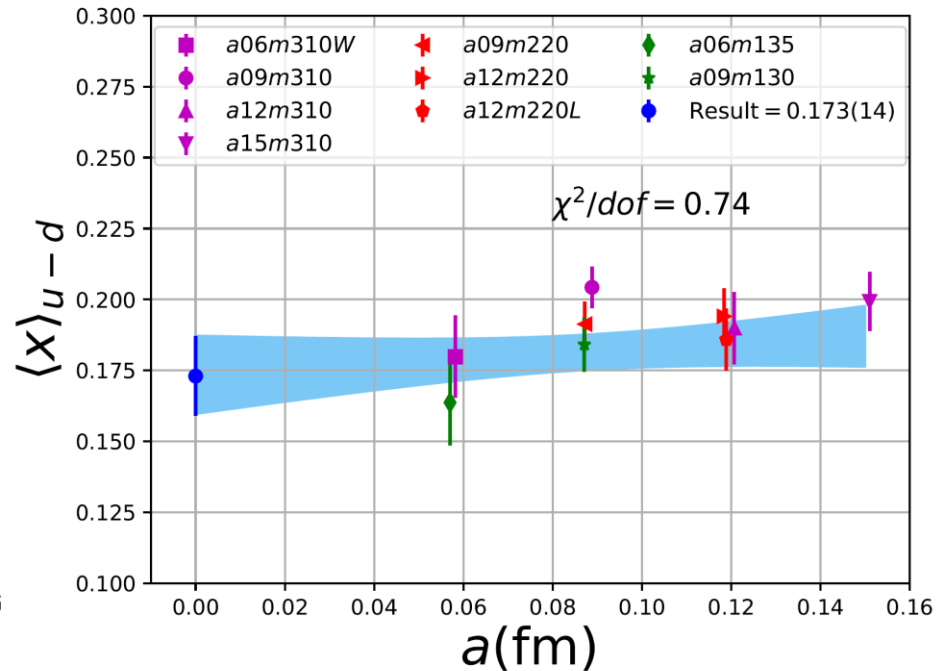
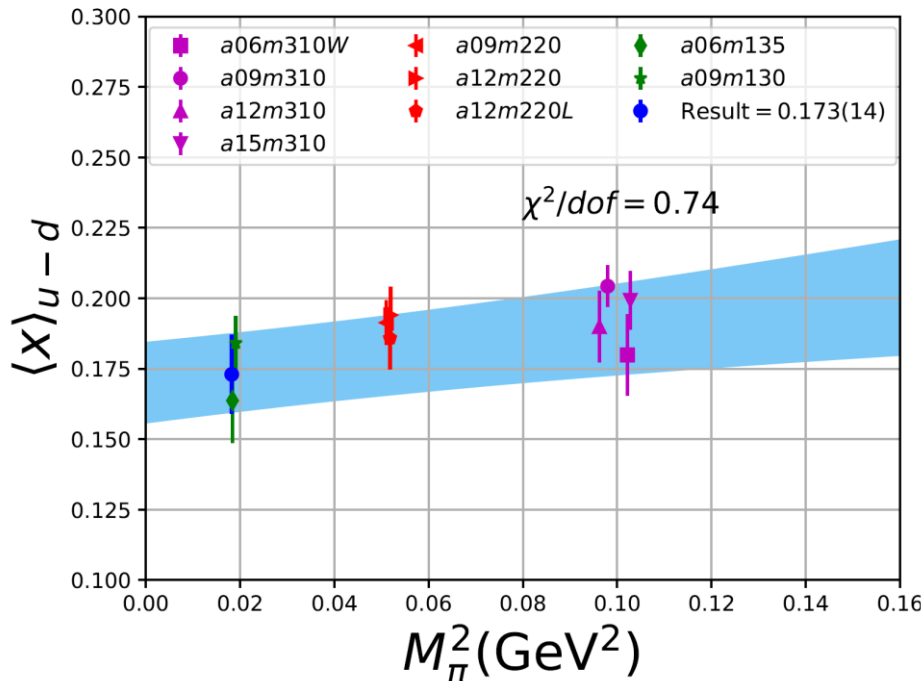


§ State-of-the art example

$$\langle x^{n-1} \rangle_q = \int_{-1}^1 dx x^{n-1} q(x)$$

∞ Extrapolate to the physical limit

Santanu Mondal et al (PNDME collaboration), 2005.13779



§ Usually more than one LQCD calculation

∞ Sometimes LQCD numbers do not even agree with each other...

# Moments of PDFs

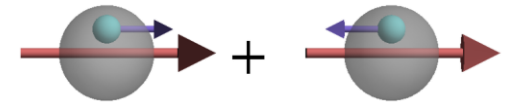
§ PDG-like rating system or average

§ LatticePDF Workshop

↻ Lattice representatives came together and devised a rating system

§ Lattice QCD/global fit status

$$\langle x^{n-1} \rangle_q = \int_{-1}^1 dx x^{n-1} q(x)$$



LatticePDF Report, 1711.07916, 2006.08636

Moment	Collaboraton	Reference	$N_f$	DE	CE	FV	RE	ES	Value	Global Fit
$\langle x \rangle_{u+-d+}$	ETMC 20	(Alexandrou <i>et al.</i> , 2020b)	2+1+1	■	★	○	★	★**	0.171(18)	0.161(18)
	PNDME 20	(Mondal <i>et al.</i> , 2020)	2+1+1	★	★	★	★	★	0.173(14)(07)	
	Mainz 19	(Harris <i>et al.</i> , 2019)	2+1	★	○	★	★	★	0.180(25)( $^{+14}_{-6}$ )	
	$\chi$ QCD 18	(Yang <i>et al.</i> , 2018b)	2+1	○	★	○	★	★	0.151(28)(29)	
	RQCD 18	(Bali <i>et al.</i> , 2019b)	2	★	★	○	★	★	0.195(07)(15)	
$\langle x \rangle_{u+}$	ETMC 20	(Alexandrou <i>et al.</i> , 2020b)	2+1+1	■	★	○	★	★**	0.359(30)	0.353(12)
	$\chi$ QCD 18	(Yang <i>et al.</i> , 2018b)	2+1	○	★	○	★	★	0.307(30)(18)	
$\langle x \rangle_{d+}$	ETMC 20	(Alexandrou <i>et al.</i> , 2020b)	2+1+1	■	★	○	★	★**	0.188(19)	0.192(6)
	$\chi$ QCD 18	(Yang <i>et al.</i> , 2018b)	2+1	○	★	○	★	★	0.160(27)(40)	
$\langle x \rangle_{s+}$	ETMC 20	(Alexandrou <i>et al.</i> , 2020b)	2+1+1	■	★	○	★	★**	0.052(12)	0.037(3)
	$\chi$ QCD 18	(Yang <i>et al.</i> , 2018b)	2+1	○	★	○	★	★	0.051(26)(5)	
$\langle x \rangle_g$	ETMC 20	(Alexandrou <i>et al.</i> , 2020b)	2+1+1	■	★	○	★	★**	0.427(92)	0.411(8)
	$\chi$ QCD 18	(Yang <i>et al.</i> , 2018b)	2+1	○	★	○	★	★	0.482(69)(48)	
	$\chi$ QCD 18a	(Yang <i>et al.</i> , 2018a)	2+1	■	★	★	★	■	0.47(4)(11)	

\*\* No quenching effects are seen.

# Moments of PDFs

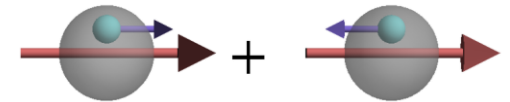
§ PDG-like rating system or average

§ LatticePDF Workshop

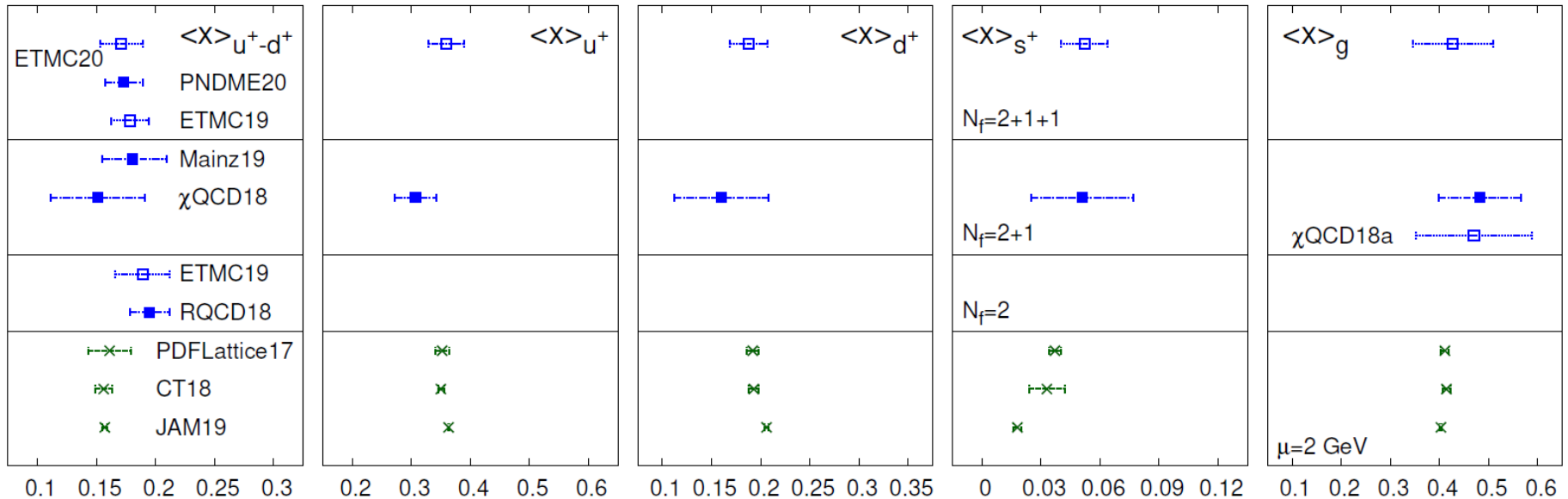
∞ Lattice representatives came together and devised a rating system

§ Lattice QCD/global fit status

$$\langle x^{n-1} \rangle_q = \int_{-1}^1 dx x^{n-1} q(x)$$



LatticePDF Report, 1711.07916, 2006.08636



# Moments of PDFs

§ PDG-like rating system or average

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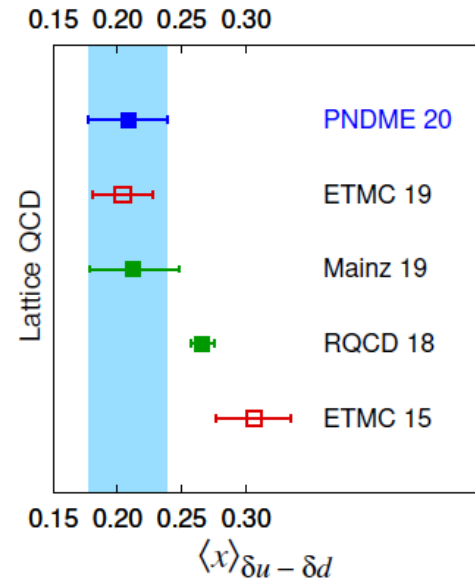
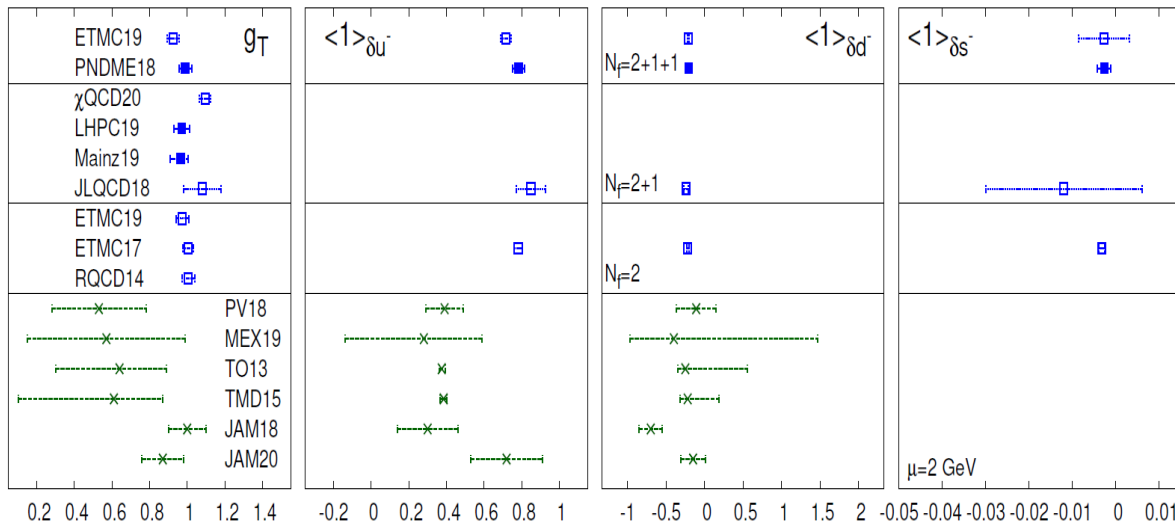
∞ Lattice representatives came together and devised a rating system

§ Recent lattice QCD/global fit status

$$\langle x^{n-1} \rangle_{\delta q} = \int_{-1}^1 dx x^{n-1} \delta q(x)$$



LatticePDF Report, 1711.07916, 2006.08636

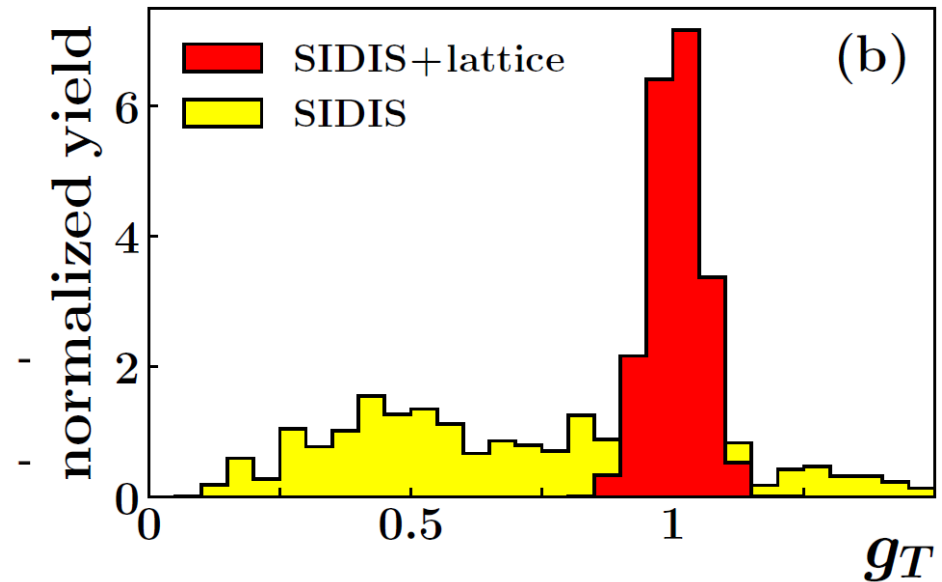
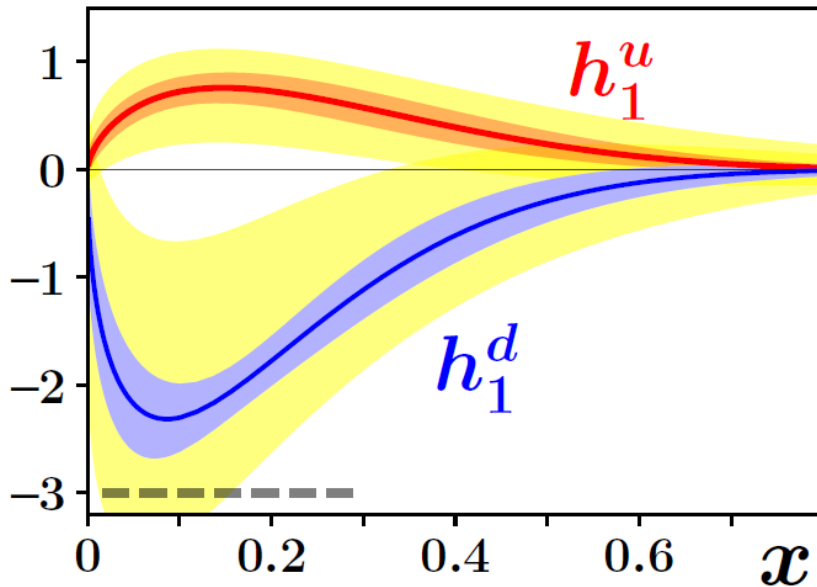


S. Mondal et al (PNDME), 2005.13779

# From Charges to PDFs

## § Improved transversity distribution with LQCD $g_T$

- ∞ Global analysis with 12 extrapolation forms:  $g_T = 1.006(58)$
- ∞ Use to constrain the global analysis fits to SIDIS  $\pi^\pm$  production data from proton and deuteron targets

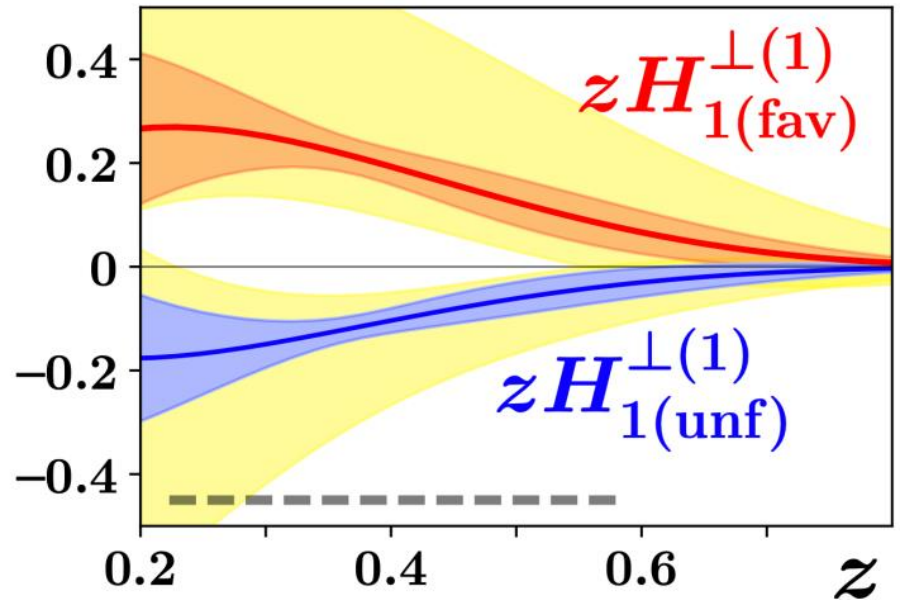
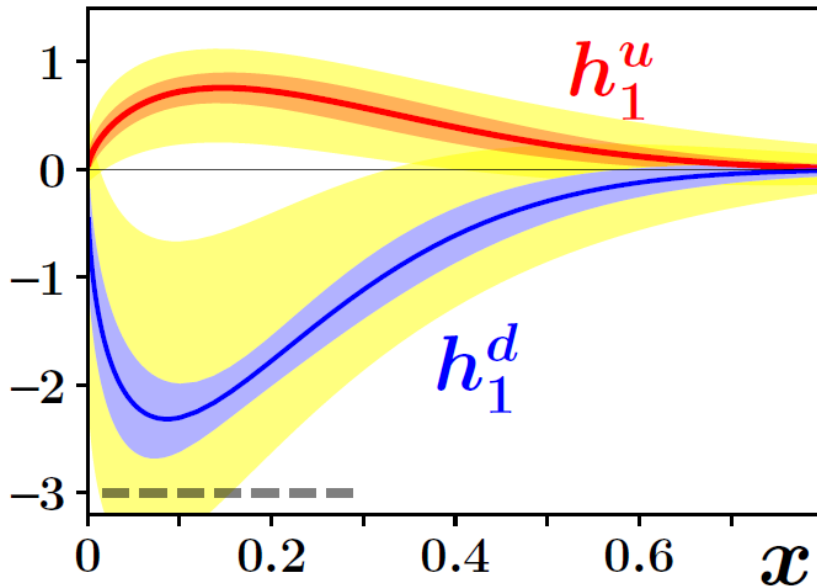


Lin, Melnitchouk, Prokudin, Sato, 1710.09858, Phys. Rev. Lett. 120, 152502 (2018)

# From Charges to PDFs

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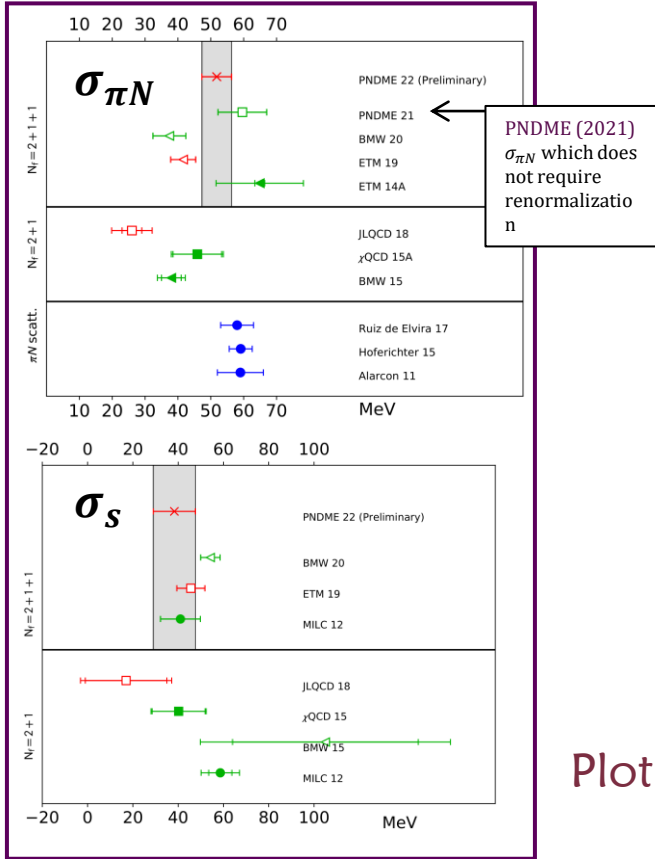
Lin, Melnitchouk, Prokudin, Sato, 1710.09858, Phys. Rev. Lett. 120, 152502 (2018)



# Nucleon Flavor Diagonal Charges

## Comparison with FLAG 2021 results

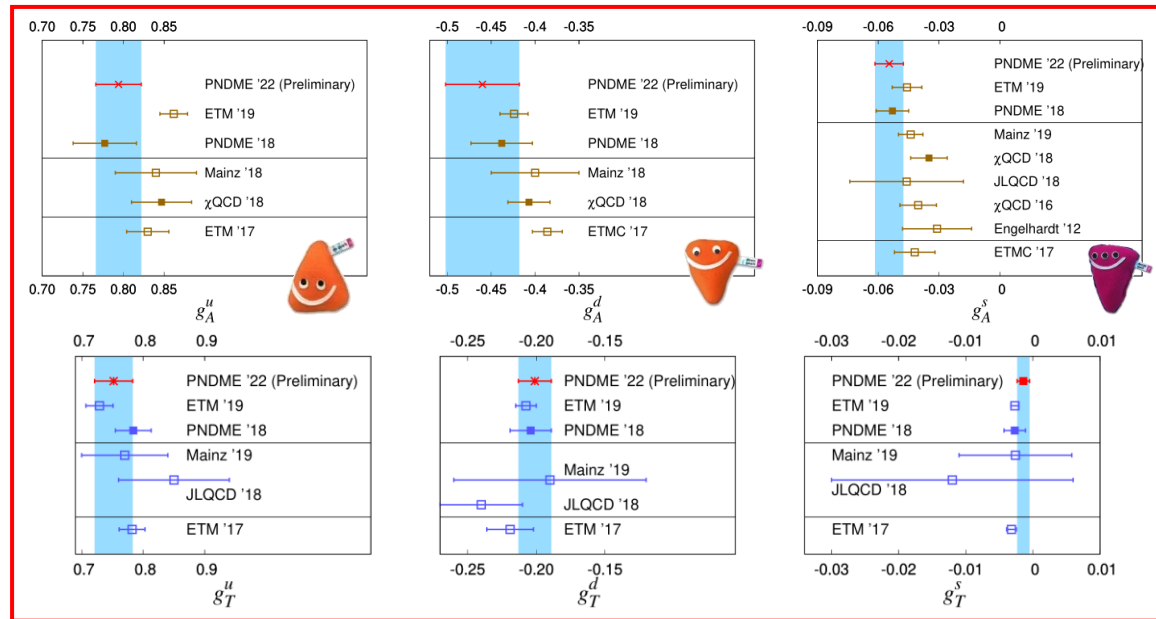
### Nucleon sigma terms (Scalar charges)



[PNDME collab., Lattice 2022 update, preliminary]

- Clover fermion on  $N_f = 2 + 1 + 1$  HISQ ensembles
- Flavor mixing calculated nonperturbatively
- **Chiral-Continuum extrapolation** including a data at  $M_\pi^{\text{Phys}}$

## Axial and Tensor charges



Plots by Sungwoo Park

# Hadron Structure

§ Structure function/distribution functions

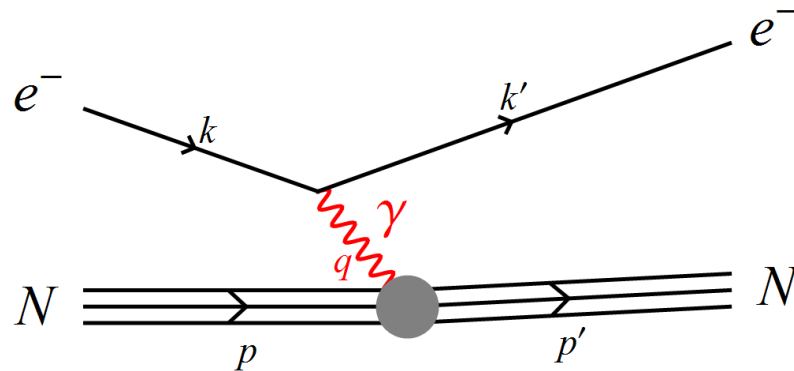
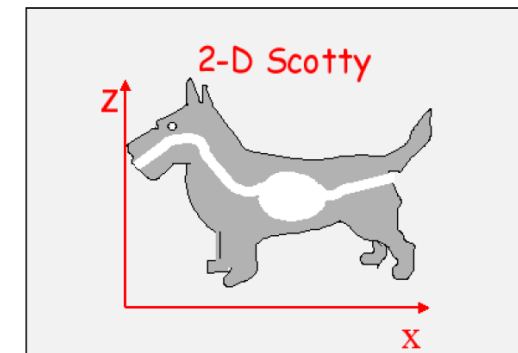
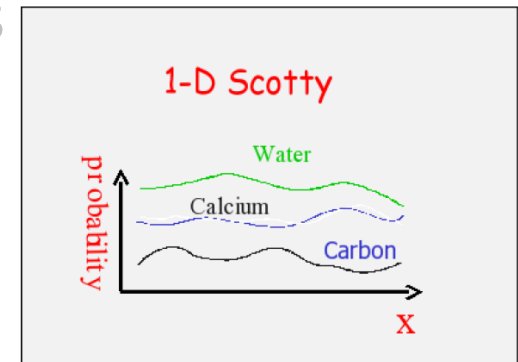
∞ deep inelastic scattering

∞  $\langle x^n \rangle_q$ ,  $\langle x^n \rangle_{\Delta q}$ ,  $\langle x^n \rangle_{\delta q}$

§ Form factors

∞ elastic scattering

∞  $F_1(Q^2)$ ,  $F_2(Q^2)$ ,  $G_A(Q^2)$ ,  $G_P(Q^2)$



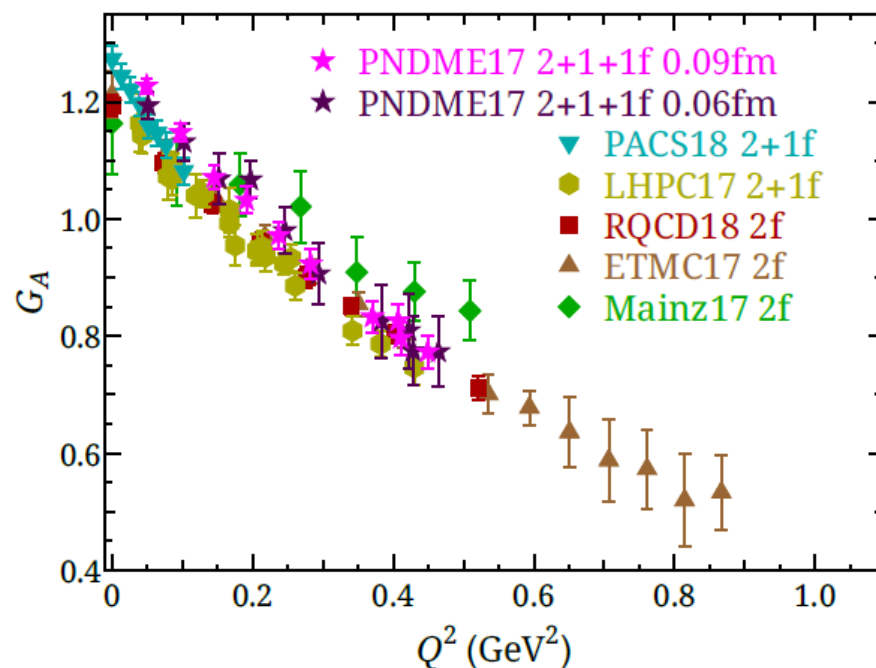
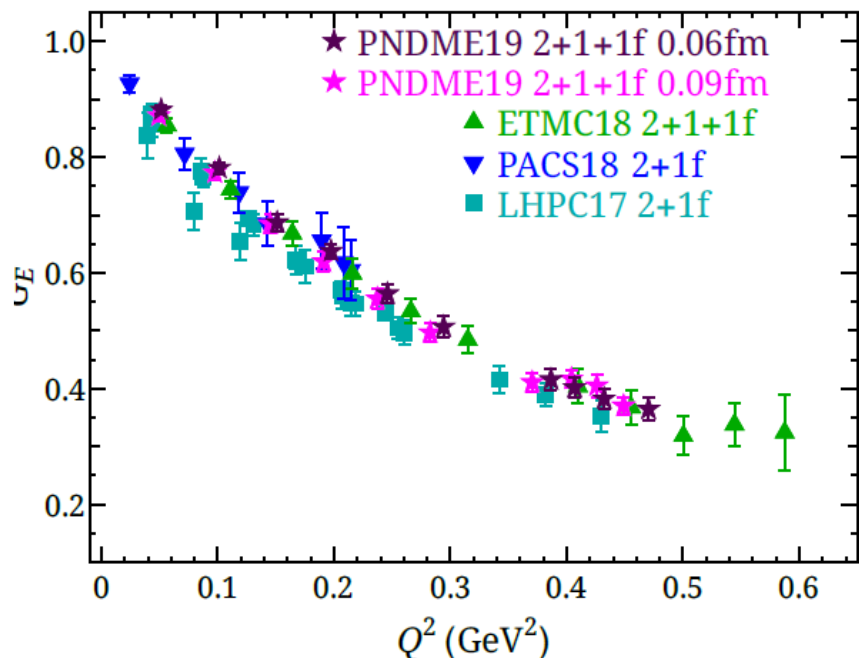
# Nucleon Form Factors

## § Nucleon *isovector* electromagnetic and axial form factor

∞ Many existing LQCD works in the past few decades

∞ Worldwide:

Increasingly many ensembles available at physical pion mass

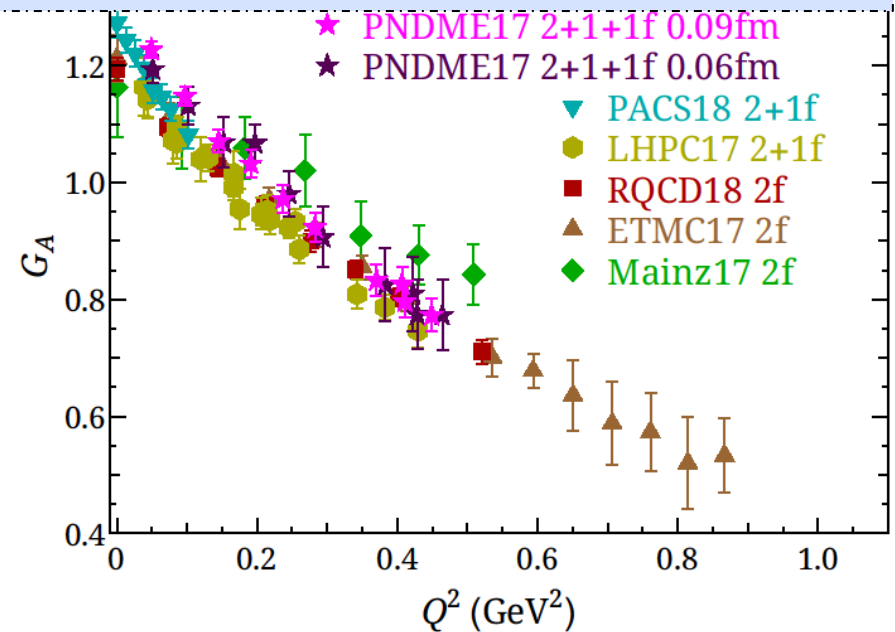
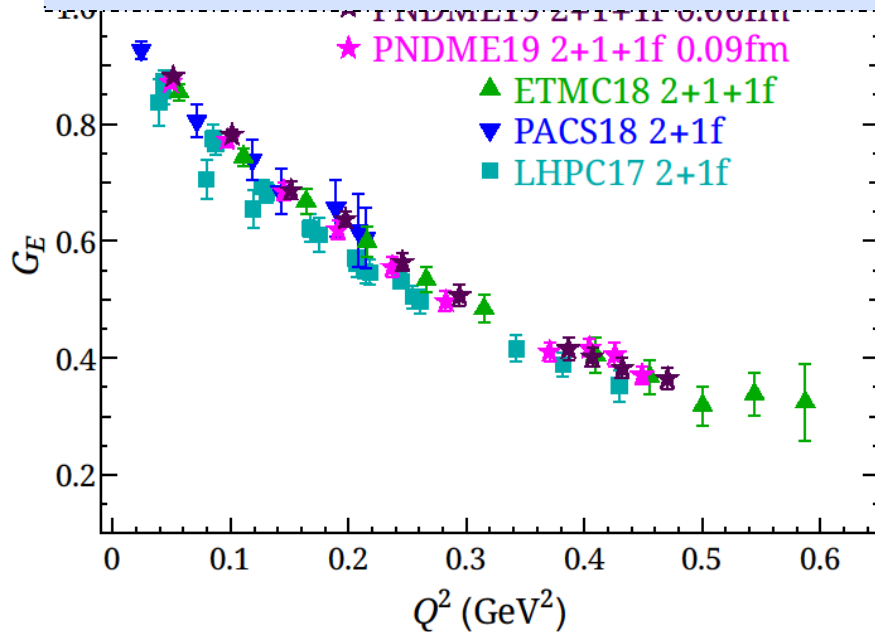


HL, Int. J. Mod. Phys. A 35 (2020) 11n12, 2030006

# Form Factors

## § Nucleon isovector electromagnetic and axial form factor

Exercise: 4) Given what you learned about the time dependence of the 3pt correlator, what do you expect to see on the ratio plots of the form-factor data?



$$\langle N(\mathbf{p}_f) | V_\mu(\mathbf{q}) | N(\mathbf{p}_i) \rangle = \bar{u}_N(\mathbf{p}_f) \left( F_1(Q^2) \gamma_\mu + \sigma_{\mu\nu} \frac{F_2(Q^2)}{2M_N} \right) u_N(\mathbf{p}_i)$$

$$G_E(Q^2) = F_1(Q^2) - \frac{Q^2}{4M_N^2} F_2(Q^2)$$

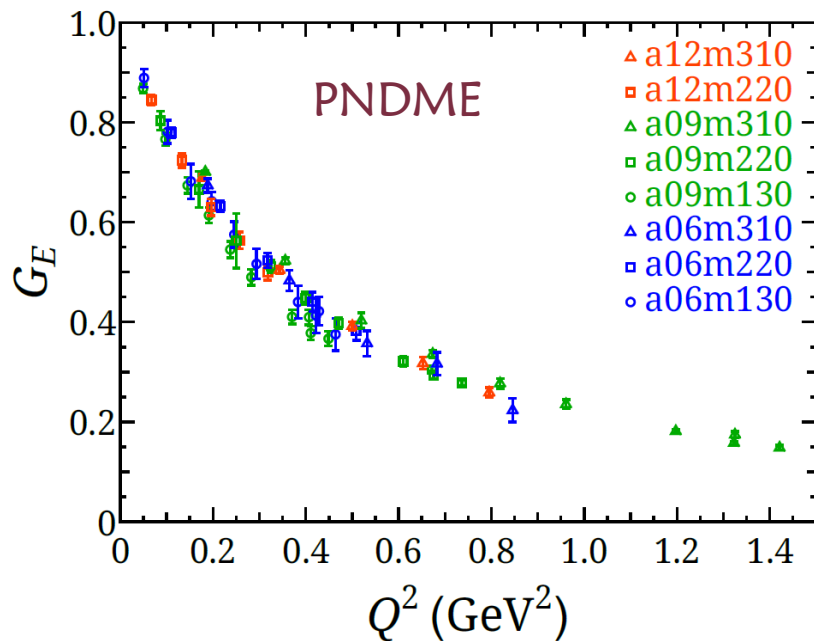
$$\langle N(\mathbf{p}_f) | A_\mu(\mathbf{q}) | N(\mathbf{p}_i) \rangle = \bar{u}_N(\mathbf{p}_f) \left( G_A(Q^2) \gamma_\mu + q_\mu \frac{\tilde{G}_P(Q^2)}{2M_N} \right) \gamma_5 u_N(\mathbf{p}_i)$$

HL, Int. J. Mod. Phys. A 35 (2020) 11n12, 2030006

# Nucleon Form Factors

## § Nucleon isovector electro

- ∞ Name of the game is now “p
- ∞ Example work supported by (Studying multiple lattice s control systematics)



USQCD Collaboration, Eur.Phys.J.A

The European Physical Journal  
volume 55 · number 11 · november · 2019

# EPJ A

Recognized by European Physical Society

## Hadrons and Nuclei

Inside: Topical Issue on Opportunities for Lattice Gauge Theory in the Era of Exascale Computing  
edited by William Detmold, Andreas Kronfeld, Ulf-G. Meißner

From: Lattice QCD and neutrino-nucleus scattering  
by USQCD Collaboration

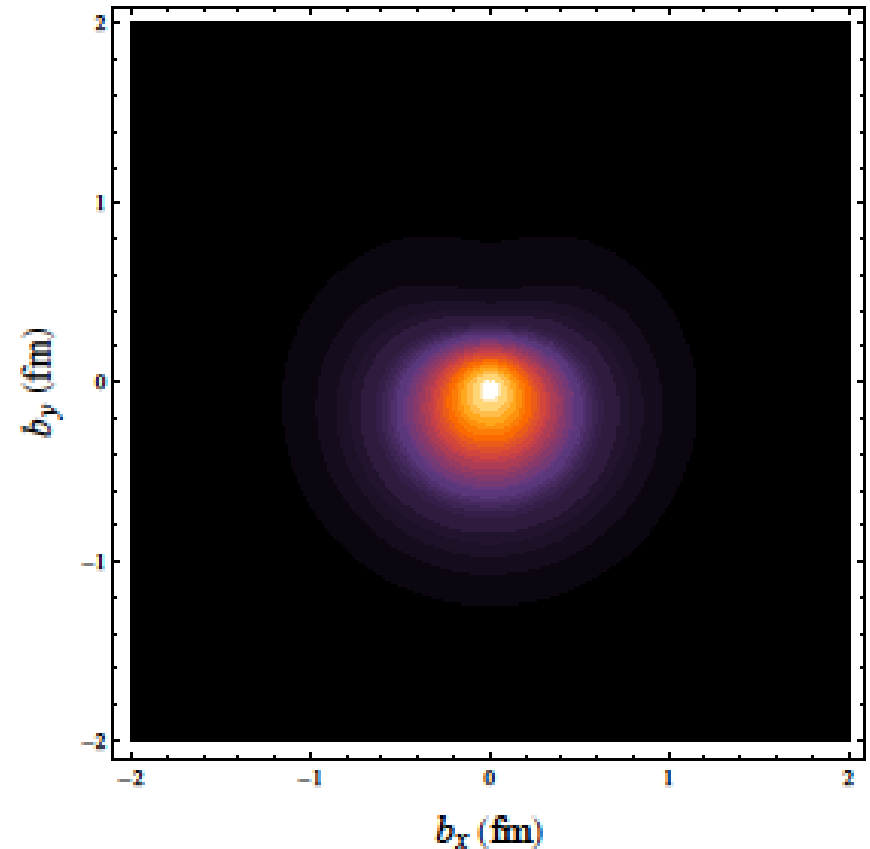
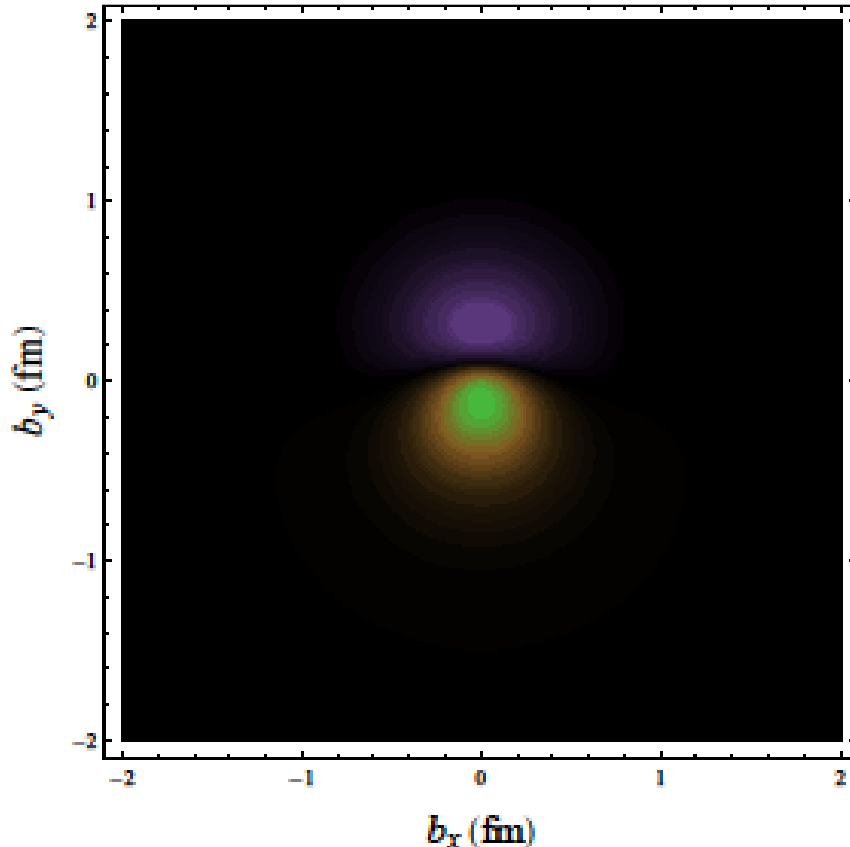
Società Italiana di Fisica

Springer

# Mapping Nucleon Picture

§ Fourier transform using large- $Q^2$  form factors to reveal transverse charge densities in a polarized nucleon

HWL, National Academies Press



# Hadron Structure

## § Structure function/distribution functions

∞ deep inelastic scattering

∞  $\langle x^n \rangle_q, \langle x^n \rangle_{\Delta q}, \langle x^n \rangle_{\delta q}$

## § Form factors

∞ elastic scattering

∞  $F_1(Q^2), F_2(Q^2), G_A(Q^2), G_P(Q^2)$

## § Generalized Parton Distribution

∞ DVCS

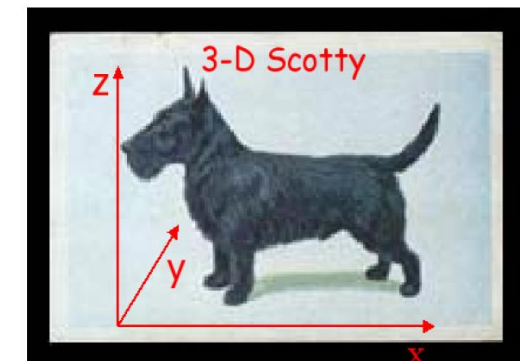
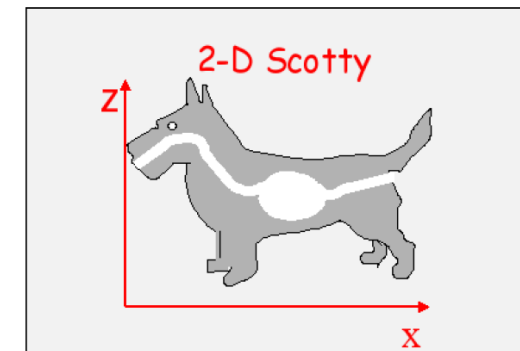
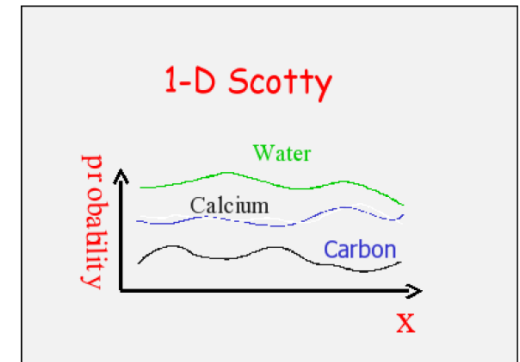
∞  $\langle x^{n-1} \rangle_q = A_{n0}(0), \langle x^{n-1} \rangle_{\Delta q} = \tilde{A}_{n0}(0),$

$\langle x^n \rangle_{\delta q} = A_{Tn0}(0),$

∞  $F_1(Q^2) = A_{10}(Q^2), F_2(Q^2) = B_{20}(Q^2),$

$G_A(Q^2) = \tilde{A}_{10}(Q^2), G_P(Q^2) = \tilde{B}_{10}(Q^2)$

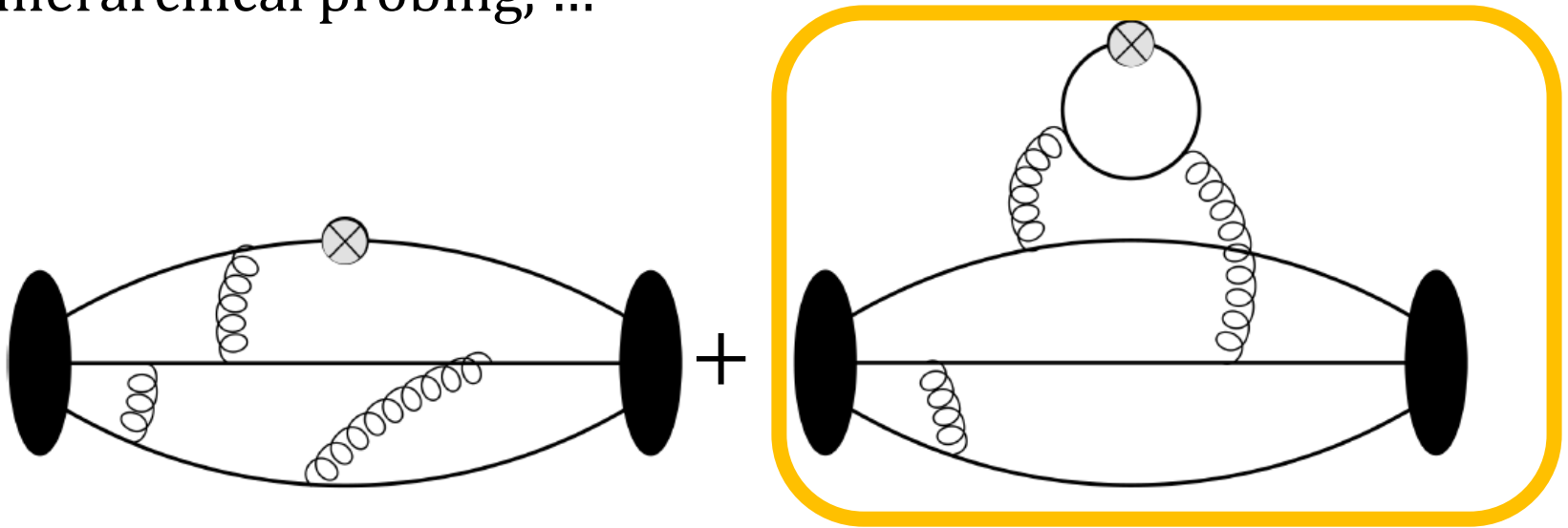
∞ Nucleon spin



# Disconnected Diagrams

## § Disconnected diagram

- ∞ Multiple ways to calculate this notorious contribution...
- ∞ Truncated solver, hopping-parameter expansion, hierarchical probing, ...





# Spin Decomposition

## § Orbital angular momentum from Ji definition

$$\vec{J}_g = \int d^3x [\vec{x} \times (\vec{E} \times \vec{B})]$$

$$\vec{J}_q = \int d^3x \psi^\dagger [\vec{\gamma} \gamma_5 + \vec{x} \times (-i\vec{D})] \psi = \frac{1}{2} (A_{20}^q + B_{20}^q)$$

quark spin

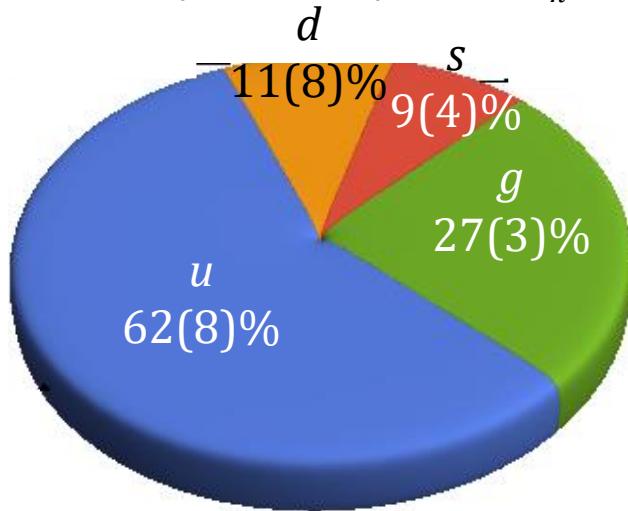
quark orbital  
angular momentum

obtained using  
GPD moment

## § Total quark and gluon contributions

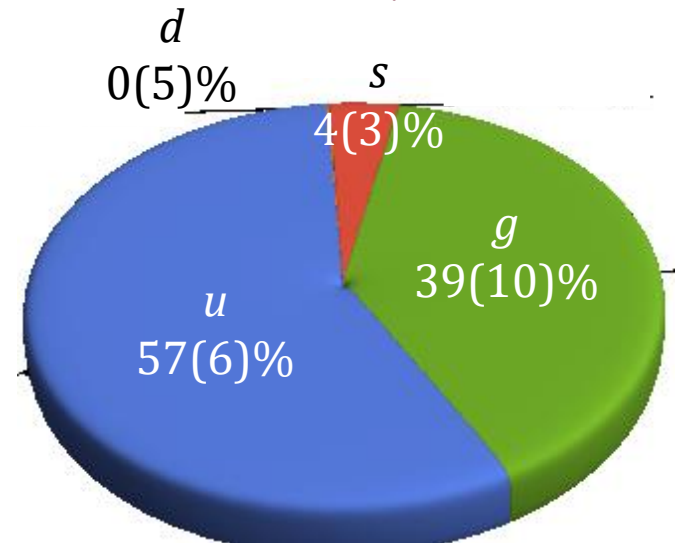
ETMC, 1706.02973

2f TM+clover, physical quark,  $M_\pi L < 3$

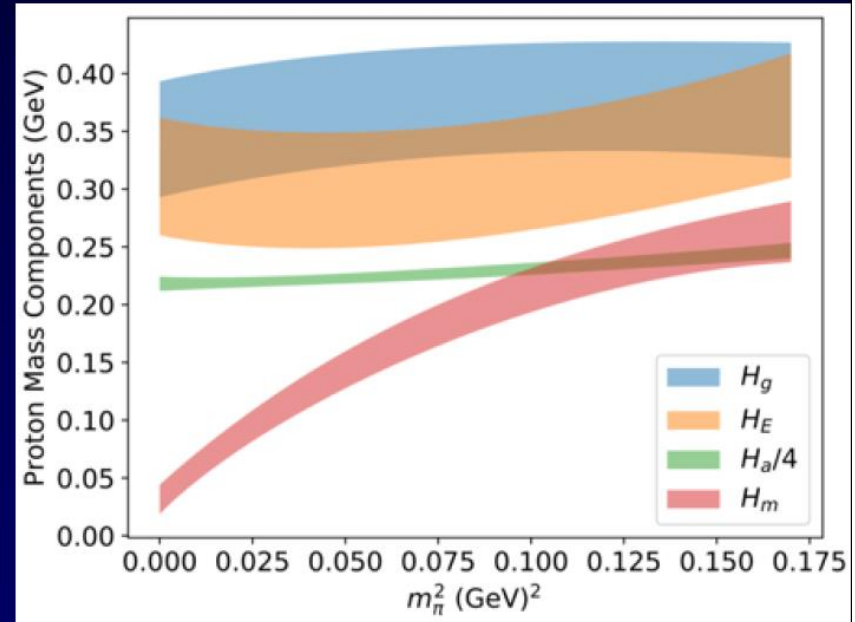
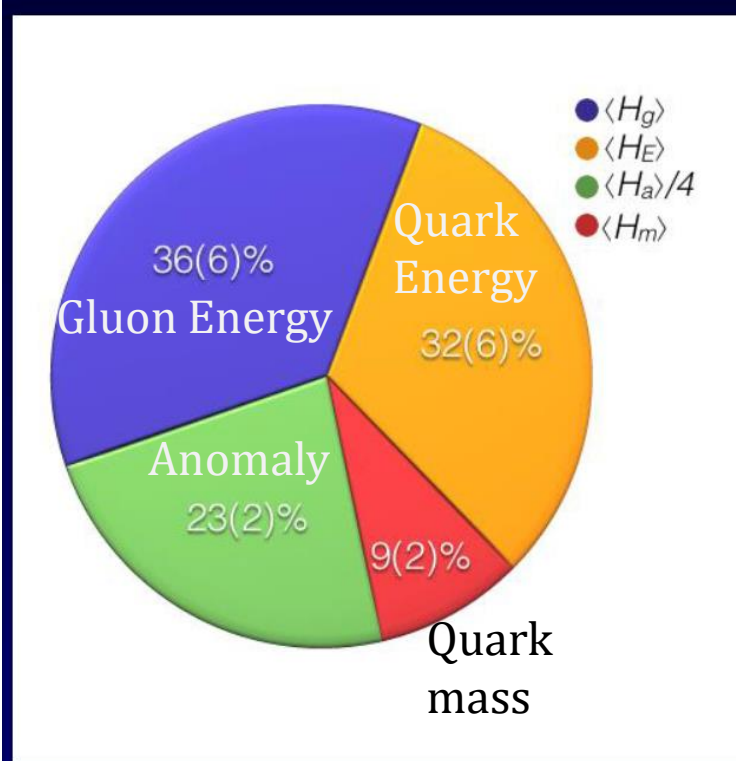


$\chi$ QCD, 1904.04138

2+1f Ov/DWF 400 MeV



# Proton Mass Decomposition



Y.B. Yang et al ( $\chi$ QCD), PRL 121, 212001 (2018)

Slide from Keh-Fei Liu

# Nucleons and New Physics

Many opportunities to probe new physics with nucleon inputs

## § Parton distribution functions for SM background

↪ Especially less known intrinsic strange/charm contribution

## § Dark matter detection Phys.Rev.D 89 (2014) 074505; ongoing work

↪ Popular candidates (e.g. SuSy neutralinos) exchange Higgs

## § Electric dipole moment Phys.Rev.Lett. 115 (2015) 21, 212002 (108 citations); Phys.Rev.D 98 (2018) 9, 091501 (43 citations)

↪ CP-violating effect, extremely small: in SM  $\approx 10^{-30}$  e-cm

## § Neutron beta decay PNDME, Phys.Rev.D 98 (2018) 034503 (84 citations); Phys.Rev.D 94 (2016) 5, 054508 (120 citations)

↪ Non- $V-A$  interactions to probe the existence of new particles (mediating new forces) with masses in the multi-TeV range

## § Nucleon (transition) form factors Phys. Rev. D 96, 114503 (67 citations)

↪ First-principles inputs for precision neutrino physics

Many of these are supported by P5 recommendations or ongoing Snowmass efforts

# Beta Decays & BSM

§ Given precision  $g_{S,T}$  and  $O_{\text{BSM}}$ , predict new-physics scales

Low-Energy

Expt  $\rightarrow$

$$O_{\text{BSM}} = f_O(\epsilon_{S,T} g_{S,T})$$

Precision LQCD input  
( $m_\pi \rightarrow 140$  MeV,  $a \rightarrow 0$ )  $\leftarrow$

$$\epsilon_{S,T} \propto \Lambda_{S,T}^{-2}$$

Upcoming precision

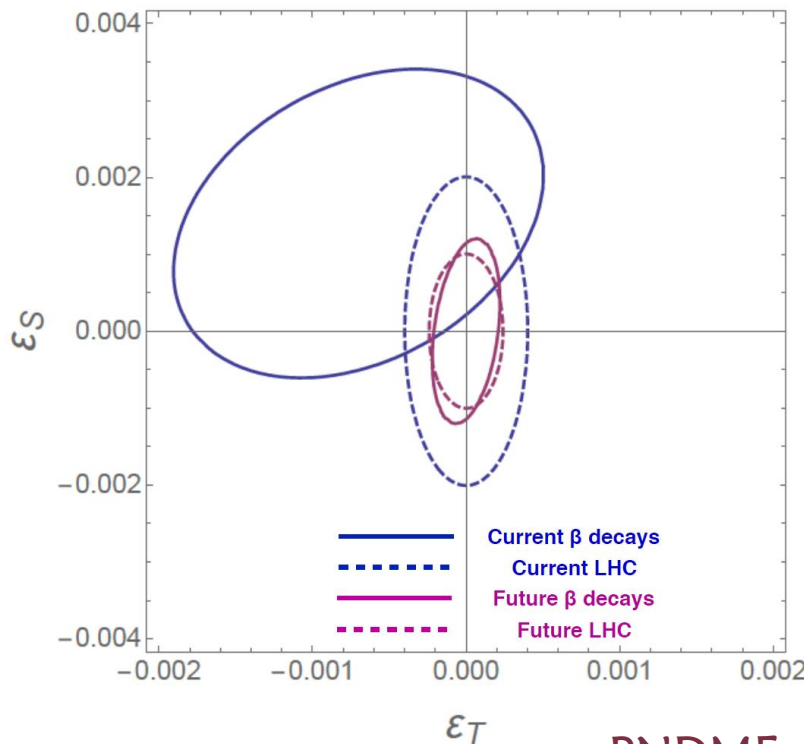
low-energy experiments

LANL/ ORNL UCN neutron  
decay exp't

$$|B_1 - b|_{\text{BSM}} < 10^{-3}$$

$$|b|_{\text{BSM}} < 10^{-3}$$

CENPA:  ${}^6\text{He}(b_{\text{GT}})$  at  $10^{-3}$



Plots by Vincenzo Cirigliano

PNDME, PRD85 054512 (2012);  
1306.5435; 1606.07049; 1806.09006

# Electric Dipole Moment

## § Why do we care?

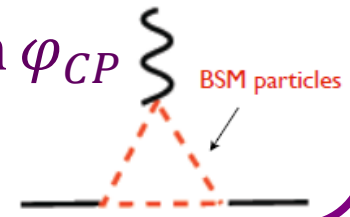
- ↻ CP-violating effect  $\Rightarrow$  Key ingredient for baryogenesis  
 $\Rightarrow$  Why matter exists
- ↻ Extremely small in SM:  $\approx 10^{-31}$  e-cm (expect to probe  $10^{-28}$  soon)
- ↻ Good candidate to constrain BSM models

§ Lagrangian  $L = L_{\text{QCD}}^{\text{CP Even}} + L_{\ominus} + L_{\text{quark}}^{\text{dim-5}} + L_{\text{chromo-quark}}^{\text{dim-5}} + \dots$

- ↻ If experiment sees signal before SM background  
 $\Rightarrow$  new physics  
 $\Rightarrow$  quark EDM (our focus here)

Induced by a variety of BSM scenarios


$$d_i \propto \frac{m_i}{\Lambda^2} \sin \varphi_{CP}$$

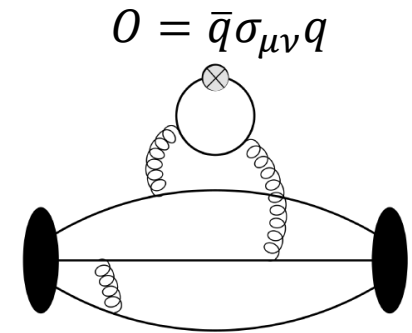


§ Lattice community are working on various contributions

# Electric Dipole Moment

§ Quark EDM ( $d_q$ ) in nucleon comes from

$$d_N = d_u g_T^{(n,u)} + d_d g_T^{(n,d)} + d_s g_T^{(n,s)}$$




§ Extrapolate to the continuum limit

PNDME, 1806.09006, 1808.07597

$$g_T^u = 0.784(28), g_T^d = -204(11), g_T^s = -0.0027(16)$$

§ Implications for new physics?

Wells, 2003;

∞ Take split SUSY for example

Arkani-Hamed and Dimopoulos, 2004;  
Giudice and Romanino, 2004

∞ Using our lattice inputs, we can derive an upper limit for the neutron EDM in split SUSY

$$|d_n| < 4 \times 10^{-28} e \cdot \text{cm}$$

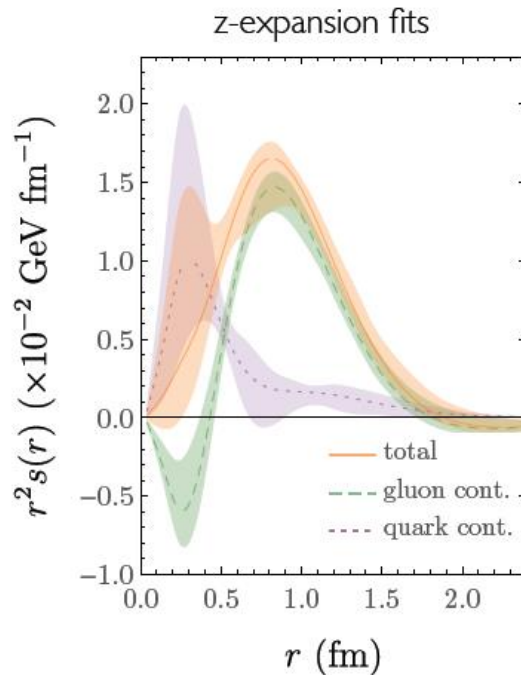
using  $|d_e| < 8.7 \times 10^{-29} e \cdot \text{cm}$  with 90% confidence

ACME Coll., Science Vol. 343 no. 6168 pp. 269-272 (2014)

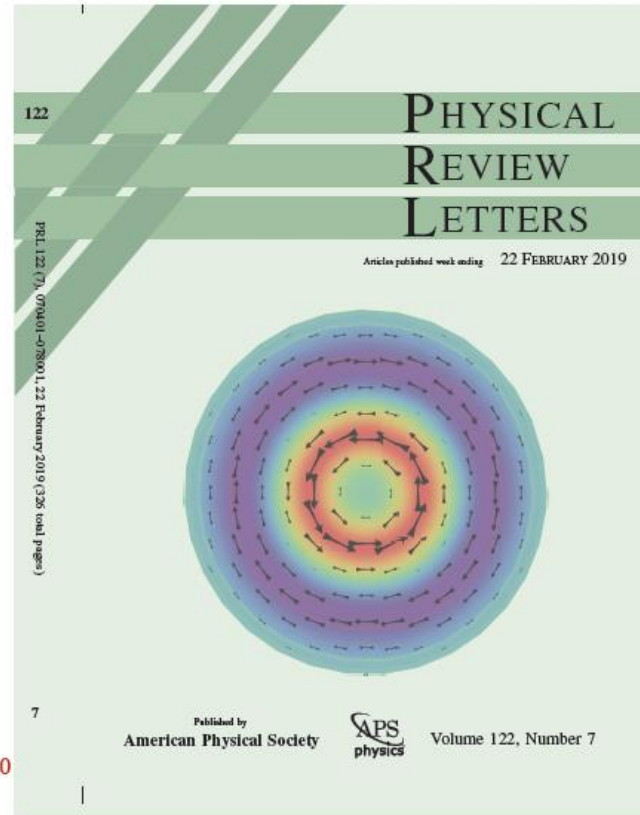
PNDME, 1806.09006, 1808.07597

# Gluon GFF

## LQCD proton shear



Gluon GFFs: Shanahan, Detmold, PRD 99, 014511, PRL 122 072003 (20)  
Quark GFFs: P. Hägler et al. (LHPC), PRD77, 094502 (2008)  
Expt quark GFFs (BEG): Burkert et al, Nature 557, 396 (2018)



Slide by Phiala Shanahan @ Lattice PDF workshop

# *Recent Lattice PDFs Progress*

Biased selected/highlighted results

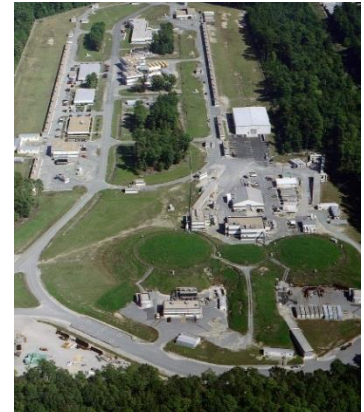
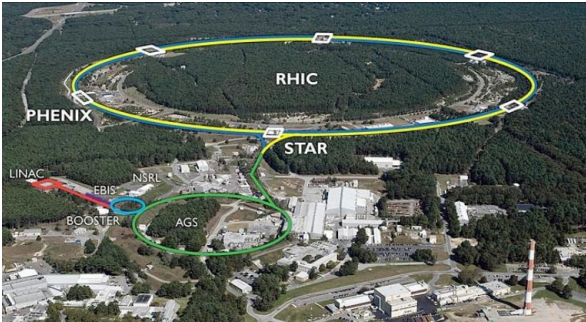




# Parton Distribution Functions

§ PDFs are universal quark/gluon distributions of nucleon

∞ Many ongoing/planned experiments  
(BNL, JLab, J-PARC, COMPASS, GSI, EIC, LHeC, ...)

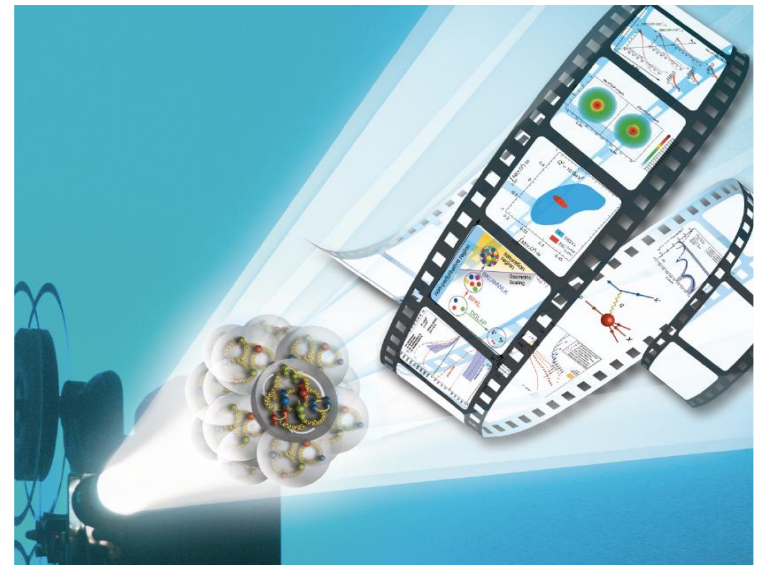


**Electron Ion Collider:  
The Next QCD Frontier**

Imaging of the proton

*How are the **sea** quarks and gluons,  
and their spins, distributed in space and  
momentum inside the nucleon?*

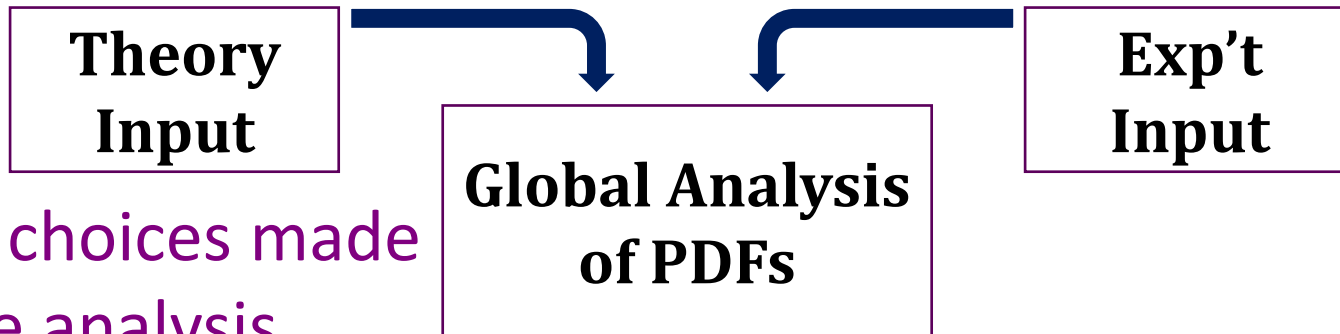
EIC White Paper, 1212.1701



# Global Analysis

§ Experiments cover diverse kinematics of parton variables

⇒ Global analysis takes advantage of all data sets



§ Some choices made for the analysis

- ⇒ Choice of data sets and kinematic cuts
- ⇒ Strong coupling constant  $\alpha_s(M_Z)$
- ⇒ How to parametrize the distribution

$$xf(x, \mu_0) = a_0 x^{a_1} (1 - x)^{a_2} P(x)$$

⇒ Assumptions imposed

SU(3) flavor symmetry, charge symmetry, strange and sea distributions

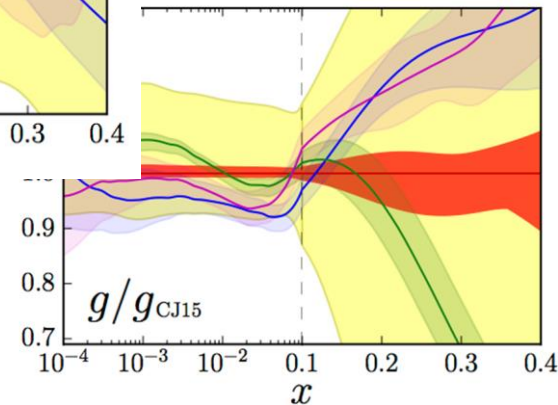
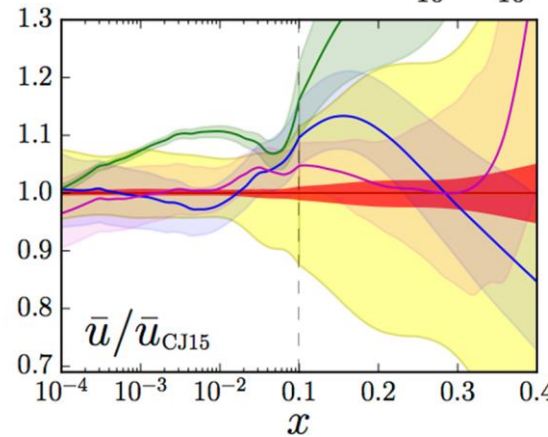
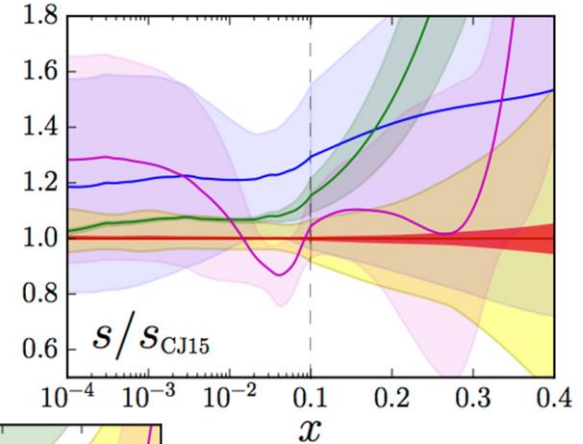
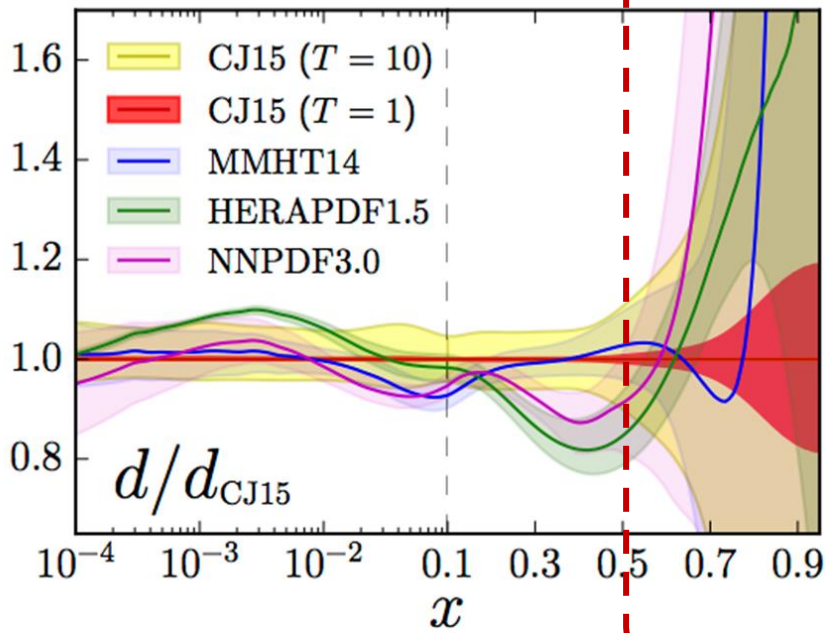
$$s = \bar{s} = \kappa(\bar{u} + \bar{d})$$

# Global Analysis

§ Discrepancies appear when data is scarce

§ Many groups have tackled the analysis

∞ CTEQ, MSTW, ABM, JR, NNPDF, etc.



CTEQ-JLAB

<https://www.jlab.org/theory/cj/>

# PDFs on the Lattice

§ Traditional lattice calculations rely on operator product expansion, only provide moments

	+	$\langle x^{n-1} \rangle_q = \int_{-1}^1 dx x^{n-1} q(x)$	most well known
spin-averaged/unpolarized			
	-	$\langle x^{n-1} \rangle_{\Delta q} = \int_{-1}^1 dx x^{n-1} \Delta q(x)$	
spin-dependent longitudinally polarized			
	-	$\langle x^{n-1} \rangle_{\delta q} = \int_{-1}^1 dx x^{n-1} \delta q(x)$	very poorly known
spin-dependent transversely polarized			



§ True distribution can only be recovered with all moments

# PDFs on the Lattice

## § Limited to the lowest few moments

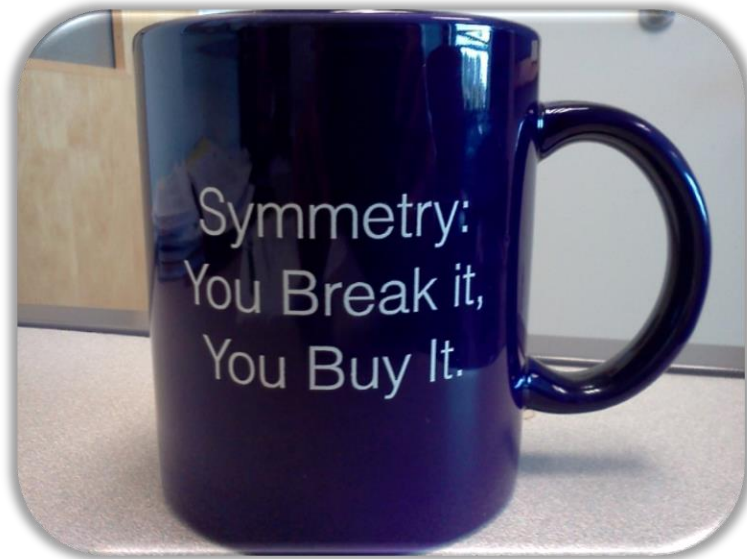
- ↪ For higher moments, all ops mix with lower-dimension ops
- ↪ Novel proposals to overcome this problem

## § Relative error grows in higher moments

- ↪ Calculation would be costly
- ↪ Hard to separate valence contrib. from sea

W. Detmold and C. Lin,  
Phys. Rev. D73 (2006)  
014501

Z. Davoudi and M. J.  
Savage, Phys. Rev. D86  
(2012) 054505



# Beyond Traditional Moments?

§ Longstanding obstacle!

§ Holy grail of structure calculations

§ Applies to many structure quantities:

∞ Generalized parton distributions (GPDs)

∞ Transverse-momentum distributions (TMD)

∞ Meson distribution amplitudes...

∞ Wigner distribution



"Marvelously  
zany humor."  
— NEWSWEEK

# A NEW HOPE

*It is a period of war and economic uncertainty.*

*Turmoil has engulfed the galactic republics.*

*Basic truths at foundation of the human civilization are disputed by the dark forces of the evil empire.*

*A small group of QCD Knights from United Federation of Physicists has gathered in a remote location on the third planet of a star called Sol on the inner edge of the Orion-Cygnus arm of the galaxy.*

*The QCD Knights are the only ones who can tame the power of the Strong Force, responsible for holding atomic nuclei together, for giving mass and shape to matter in the Universe.*

*They carry secret plans to build the most powerful*

# *Bjorken- $x$ Dependent Hadron Structure*

Biased selected results, highlighting work  
done by MSU students/postdocs

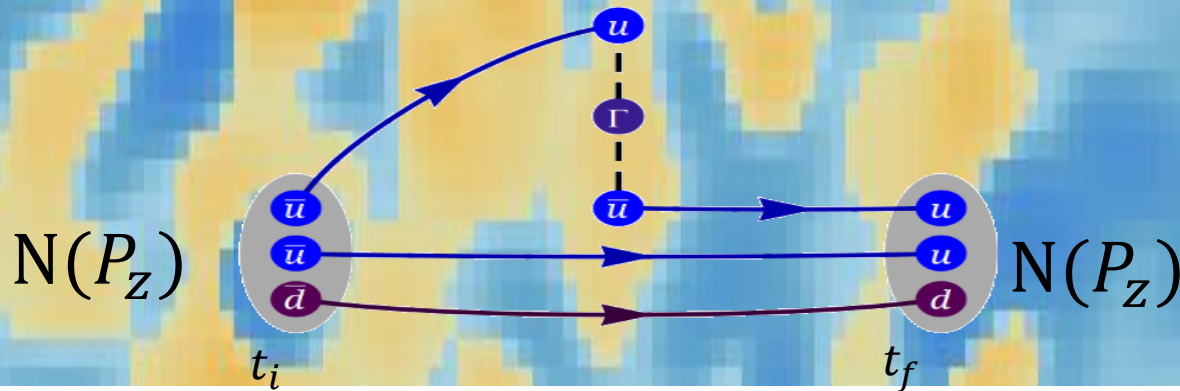




# Lattice Parton Method

## § Large-momentum effective theory (LaMET)/quasi-PDF

(X. Ji, 2013; See 2004.03543 for review)



## § Compute quasi-distribution via

$$\tilde{q}(x, \mu, P_z) = \int \frac{dz}{4\pi} e^{-izk_z} \left\langle P \left| \bar{\psi}(z) \Gamma \exp\left(-ig \int_0^z dz' A_z(z')\right) \psi(0) \right| P \right\rangle$$

## § Recover true distribution (take $P_z \rightarrow \infty$ limit)

$$\tilde{q}(x, \mu, P_z) = \int_{-\infty}^{\infty} \frac{dy}{|y|} C\left(\frac{x}{y}, \frac{\mu}{P_z}\right) \mathbf{q}(y, \mu) + \mathcal{O}\left(\frac{M_N^2}{P_z^2}, \frac{\Lambda_{\text{QCD}}^2}{(xP_z)^2}, \frac{\Lambda_{\text{QCD}}^2}{((1-x)P_z)^2}\right)$$

X. Xiong et al., 1310.7471; J.-W. Chen et al, 1603.06664

# Lattice Parton Method

## § Short-distance factorization (SDF)

∞ pseudo-PDF method (A. Radyushkin, 2017)

∞ Hadronic tensor currents

(Liu et al., hep-ph/9806491, ... 1603.07352)

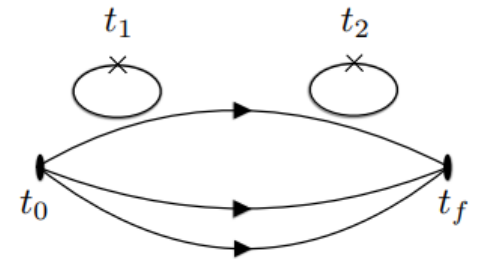
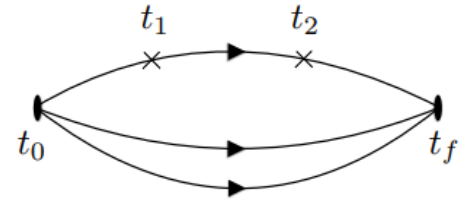
∞ Lattice cross-section method (LCS)

(Y Ma and J. Qiu, 2014, 2017)

∞ Euclidean correlation functions

(RQCD, 1709.04325)

∞ Compton amplitude approach (QCDSF, 1703.01153)



Quantities  
that can be  
calculated  
on the lattice  
today

=  $\Sigma$

Wanted  
PDFs,  
GPDs,  
etc.

$\times$

pQCD-  
calculated  
kernel

# Lattice Parton Calculations

## § Rapid developments!

HL, Few Body Syst. 64 (2023) 3, 58

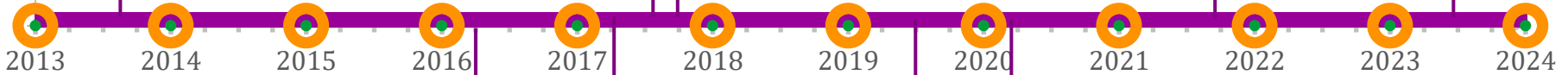
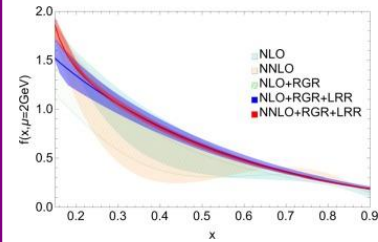
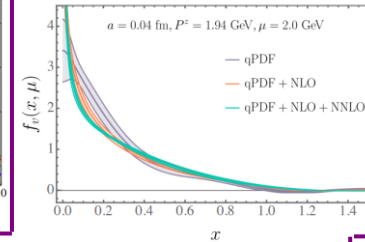
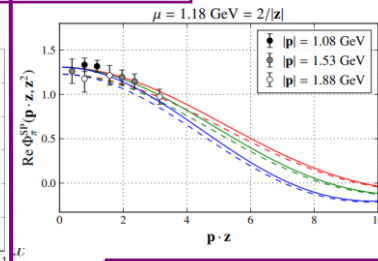
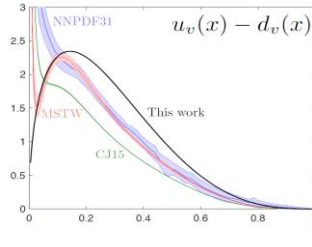
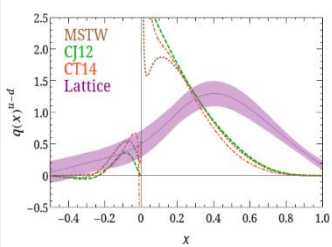
**First unpol. PDF lattice calculation**

**First lattice pseudo-PDFs**

**Euclidean correlation functions**

**1<sup>st</sup> NNLO PDF**

**1<sup>st</sup> PDF w/ LRR+RGR**

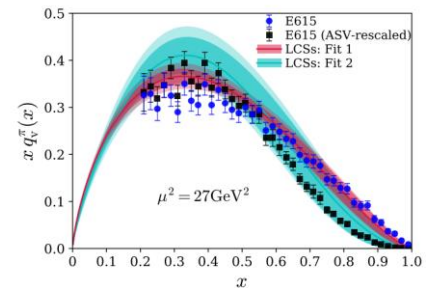
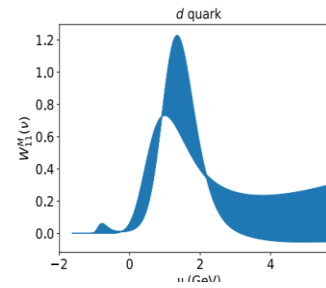
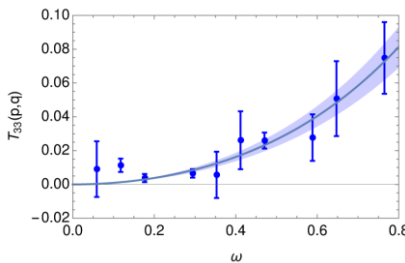
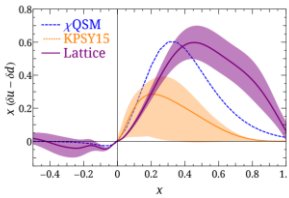
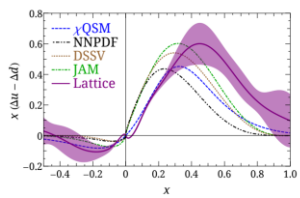


**Pol. PDFs and mass corrections**

**Compton amplitude**

**Hadronic tensor**

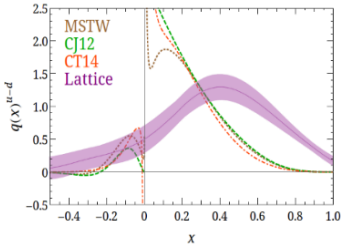
**LCS**



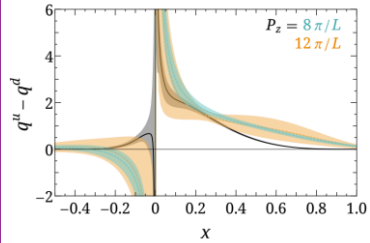
# Lattice Parton Calculations

## § Physics quantity milestones

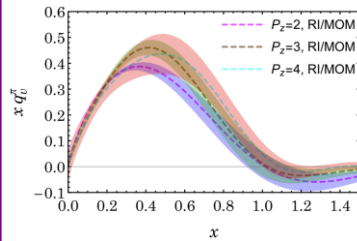
### First unpol. lattice PDF



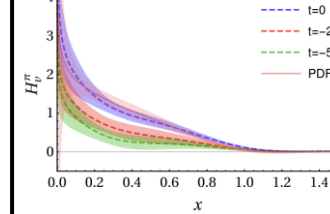
### First PDFs at $M_\pi^{\text{phys}}$



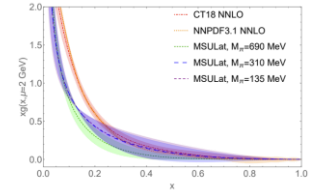
### Pion v-PDF



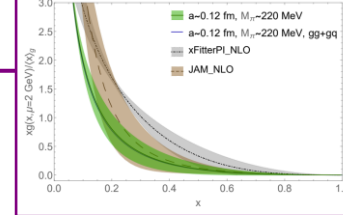
### 1st GPD ( $\pi$ )



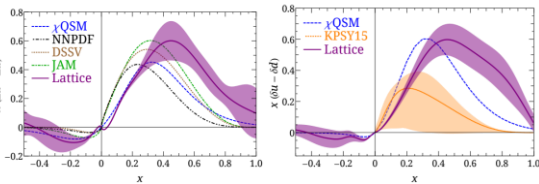
### $N$ $g$ -PDF



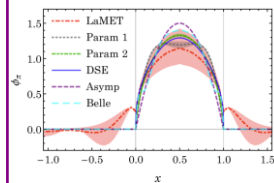
### $\pi$ $g$ -PDF



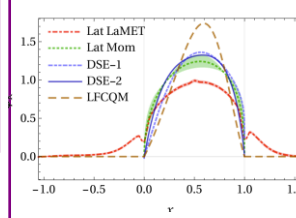
### Pol. PDFs and mass corrections



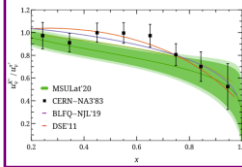
### Pion DA



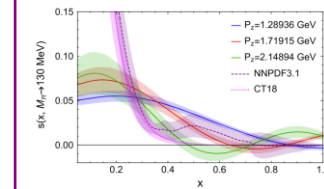
### Kaon DA



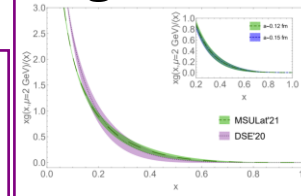
### $K$ PDF



### $s, c$ PDF



### Kaon $g$ -PDF

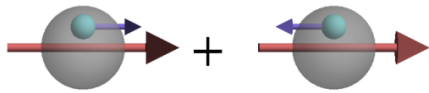


HL, Few Body Syst. 64 (2023) 3, 58

# Lattice Example Results

## § Summary of physical pion mass PDFs results

unpolarized



$$u(x) - d(x)$$

longitudinally polarized

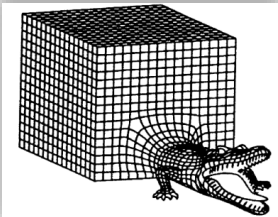
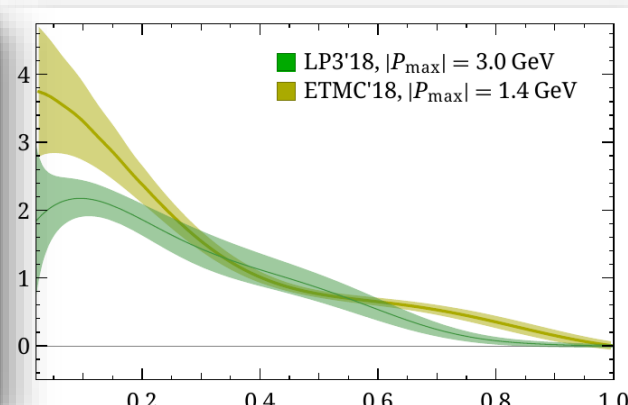
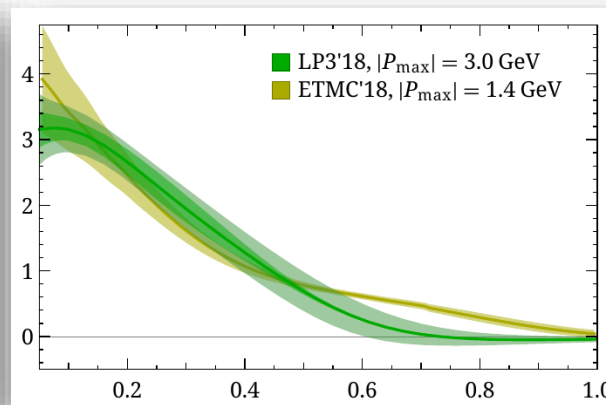
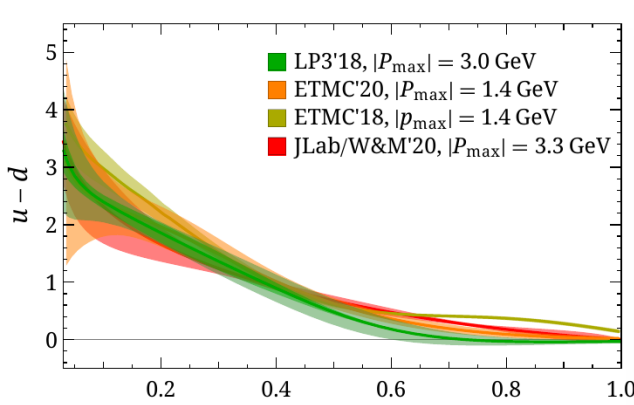


$$\Delta u(x) - \Delta d(x)$$

transversely polarized



$$\delta u(x) - \delta d(x)$$



Finite volume,  
Discretization,

...

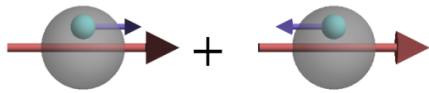


2006.08636 (PDFLattice2019)

# Lattice Example Results

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unpolarized



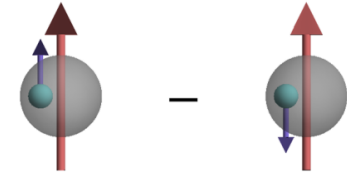
$$u(x) - d(x)$$

longitudinally polarized

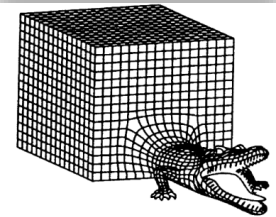
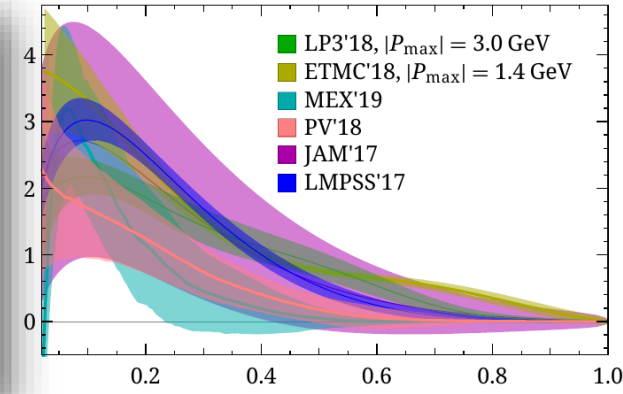
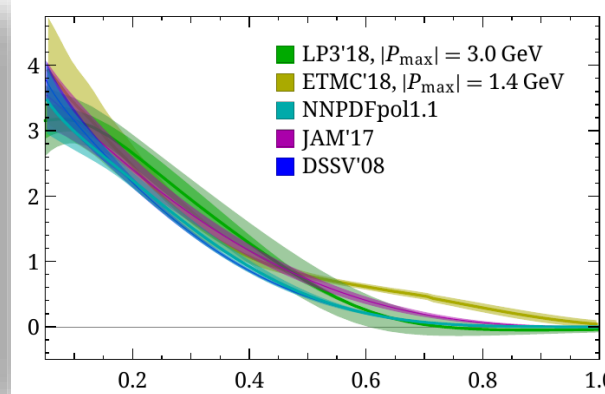
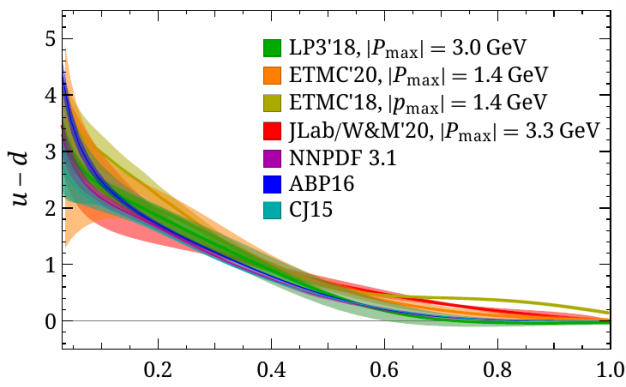


$$\Delta u(x) - \Delta d(x)$$

transversely polarized



$$\delta u(x) - \delta d(x)$$



Finite volume,  
Discretization,

...



2006.08636 (PDFLattice2019)

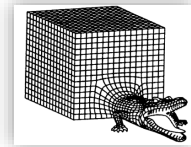
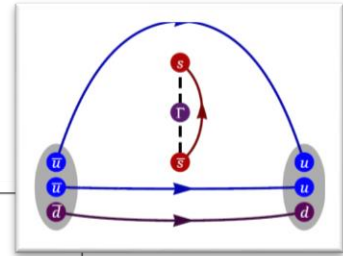
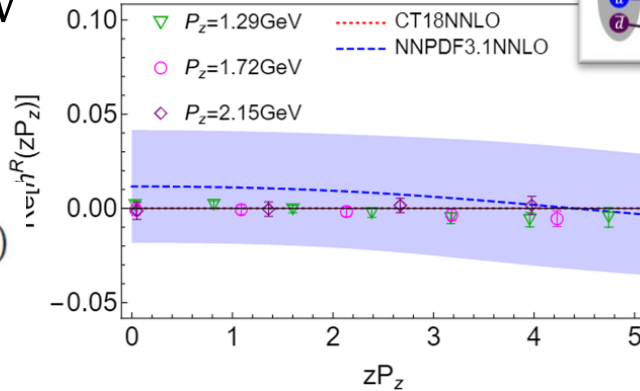
# First Lattice Strange PDF

## § Results by MSULat/quasi-PDF method

- ☞ Clover on 2+1+1 HISQ, 0.12-fm 310-MeV QCD vacuum
- ☞ Extrapolated to  $M_\pi \approx 140$  MeV

2005.01124, R. Zhang et al  
(MSULat)

$$\text{Re}[h(z)] \propto \int dx (s(x) - \bar{s}(x)) \cos(xzP_z)$$



Quantities that can be calculated on the lattice



# First Lattice Strange PDF

## § Results by MSULat/quasi-PDF method

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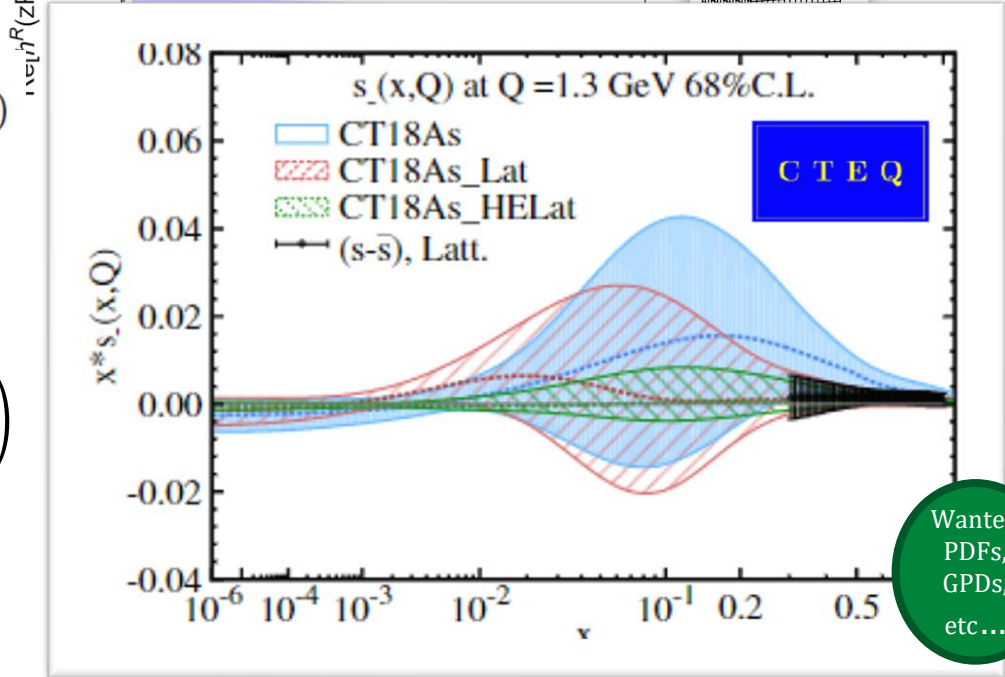
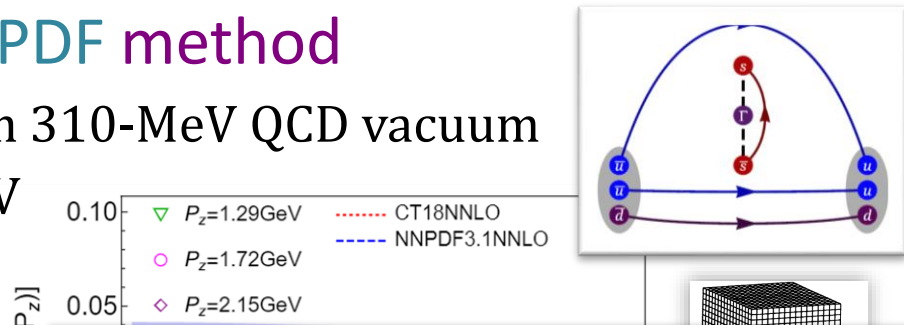
2005.01124, R. Zhang et al (MSULat)

$$\text{Re}[h(z)] \propto \int dx (s(x) - \bar{s}(x)) \cos(xzP_z)$$

## § From quasi-PDF to PDF

$$\tilde{f}_q(x, P_z) = \int_{-1}^1 \frac{dy}{|y|} f_q(y) C_{q/q}(x, y, P_z, \mu) + O\left(\frac{\Lambda_{\text{QCD}}^2}{x^2 P_z^2}, \frac{\Lambda_{\text{QCD}}^2}{(1-x)^2 P_z^2}\right)$$

T. Hou, HL, M. Yan, C. Yuan, 2204.07944



§ The strangeness asymmetry  $s(x, Q) - \bar{s}(x, Q)$  at  $x > 0.2$  is difficult to measure, but can be predicted in lattice QCD



# First Lattice Charm PDF

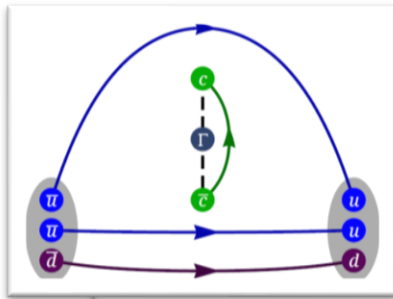


§ Large uncertainties in global PDFs

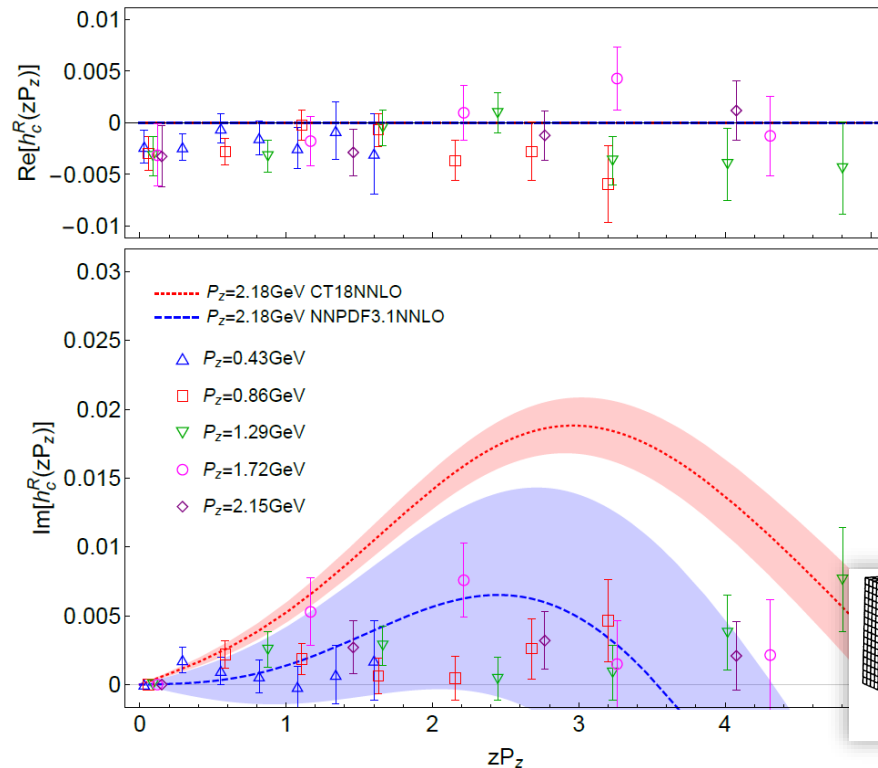
§ Results by MSULat/quasi-PDF method

☞ Clover on 2+1+1 HISQ 0.12-fm 310-MeV QCD vacuum

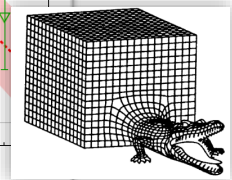
2005.01124, R. Zhang et al (MSULat)



- suggest a symmetric  $c - \bar{c}$  distribution
- much smaller than strange PDF



Quantities that can be calculated on the lattice



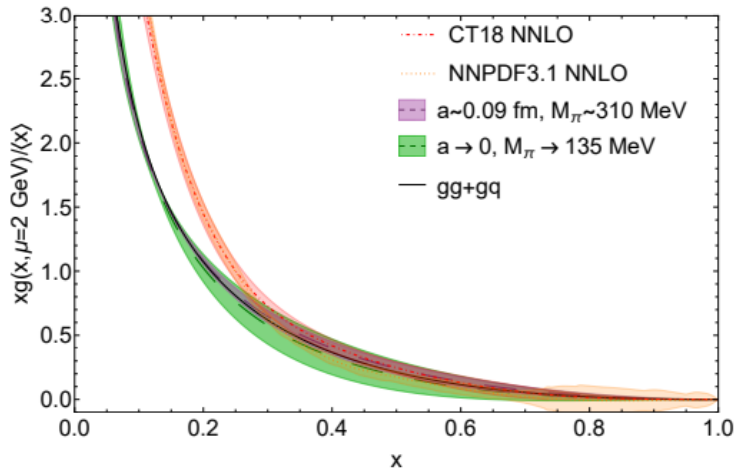
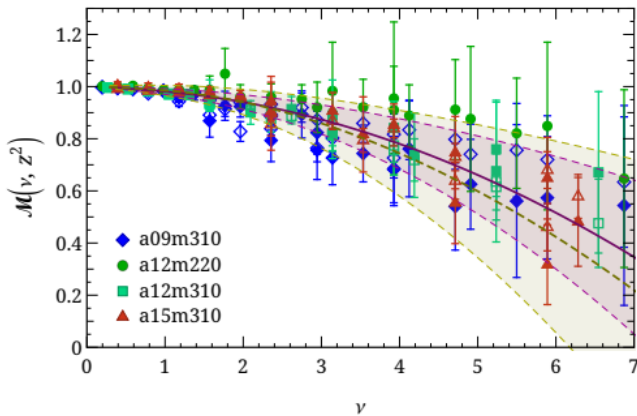
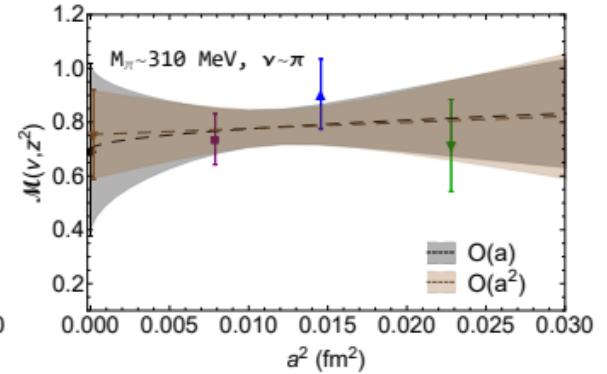
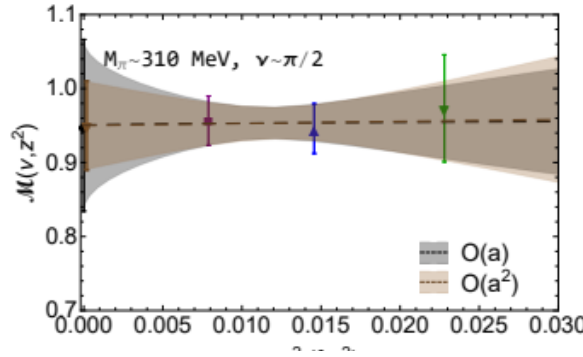
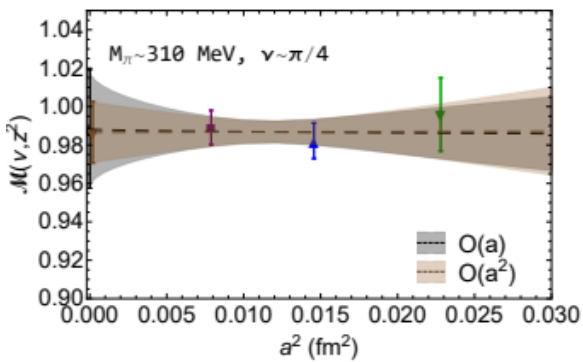
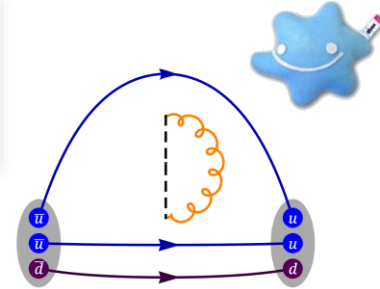
# Gluon PDF in Nucleon

§ Continuum Gluon PDF w/ pseudo-PDF

∞ 2+1+1 HISQ {0.09, 0.12, 0.15} fm,

[220,310,700]-MeV pion,  $10^5$ - $10^6$  statistics

[arXiv:2210.09985](https://arxiv.org/abs/2210.09985)



G: Bill Good

# Nucleon Tomography

§ Assuming we live in the Marvel Universe

∞ The special quantum tunnel allows us to shrink to the size particle to sub-nucleon scale ( $< 10^{-15}\text{m}$ )



§ What would it look like to travel inside the nucleon?

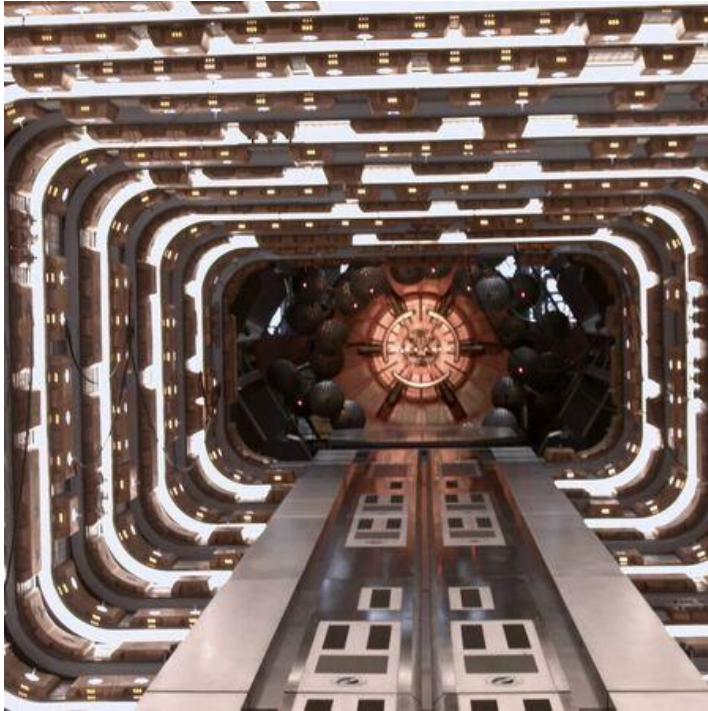
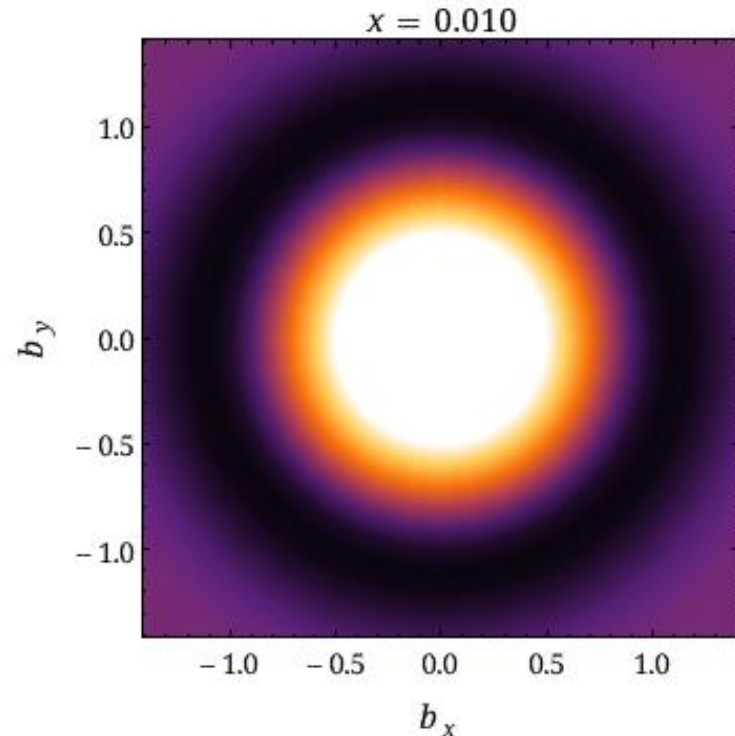


Image credit: Marvel Studios



Thanks to Cottrell Scholar Award from RCSA

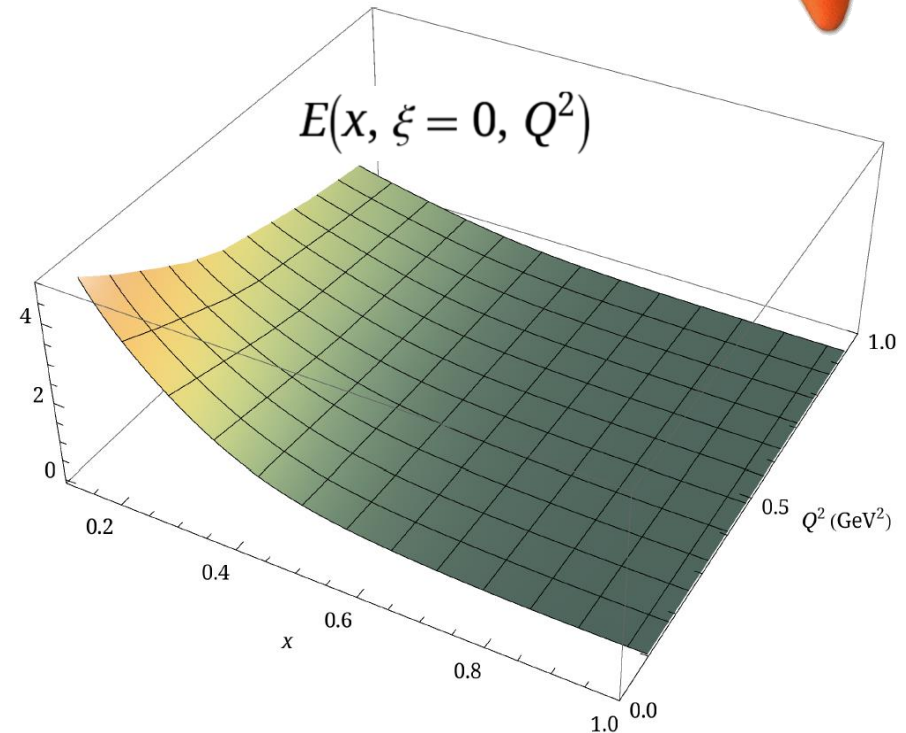
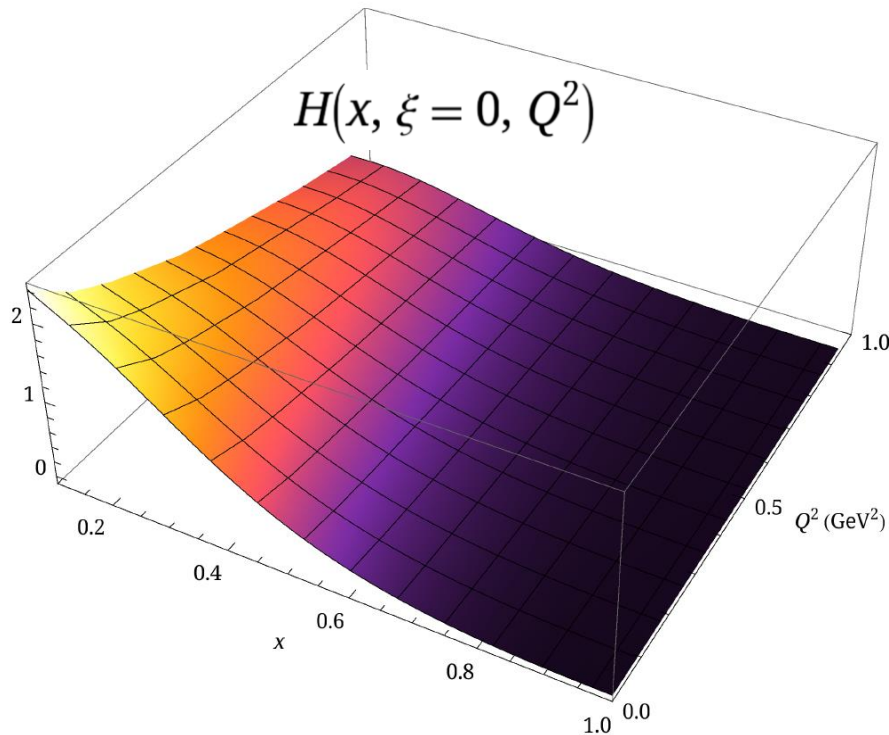
# Isvector Nucleon GPDs

§ Nucleon GPD using quasi-PDFs at **physical pion mass**

∞ MSULat: clover/2+1+1 HISQ

0.09 fm, 135-MeV pion mass,  $P_z \approx 2$  GeV

∞  $\xi = 0$  isovector nucleon GPD results



HL, Phys.Rev.Lett. 127 (2021) 18, 182001

# Isvector Nucleon GPDs

## § Nucleon GPD using quasi-PDFs at physical pion mass

∞ Lattice details: clover/2+1+1 HISQ (MSULat)

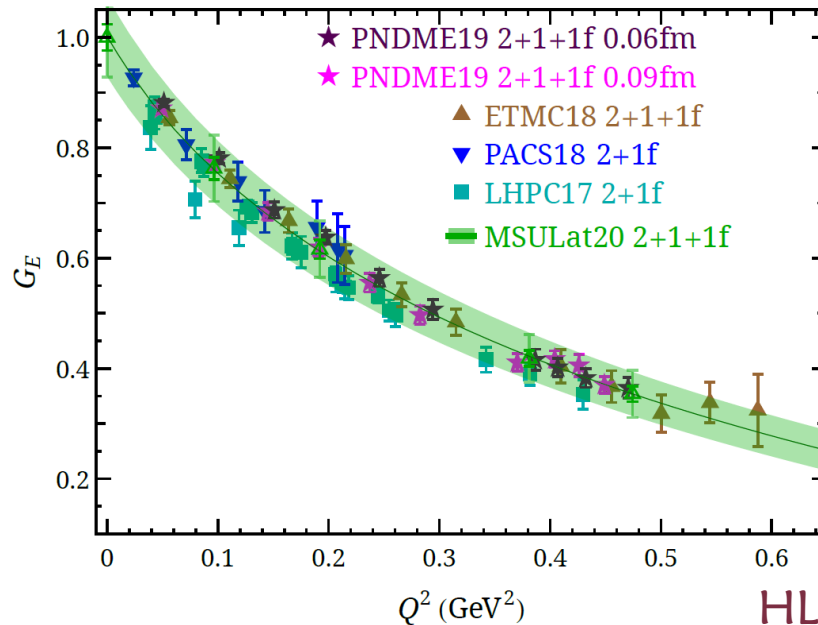
0.09 fm, **135-MeV** pion mass,  $P_z \approx 2$  GeV

∞  $\xi = 0$  isovector nucleon GPD results

$$\int_{-1}^{+1} dx x^{n-1} \text{[3D plot]} = \sum_{i=0, \text{even}}^{n-1} (-2\xi)^i A_{ni}^q(t) + (-2\xi)^n C_{n0}^q(t) \Big|_{n \text{ even}}$$



$n = 1$



HL, Phys.Rev.Lett. 127 (2021) 18, 182001

# Nucleon GPDs

## § Nucleon GPD using quasi-PDFs at physical pion mass

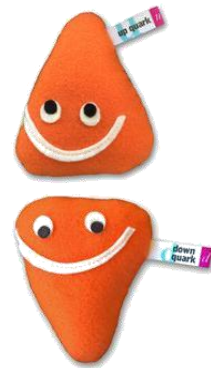
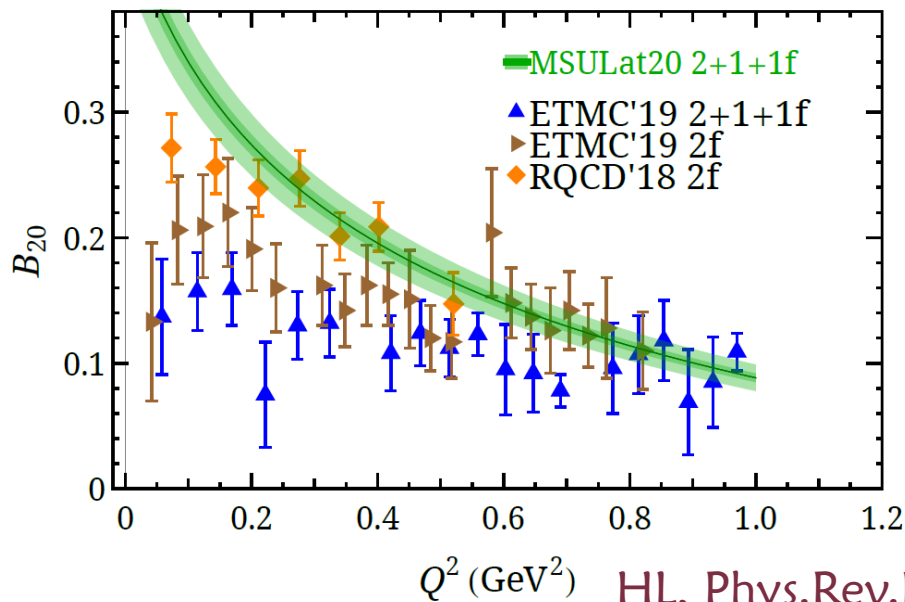
☞ Lattice details: clover/2+1+1 HISQ (MSULat)

0.09 fm, **135-MeV** pion mass,  $P_z \approx 2$  GeV

☞  $\xi = 0$  isovector nucleon GPD results

$$\int_{-1}^{+1} dx x^{n-1} \text{ (3D plot of } x \text{ vs } Q^2 \text{)} = \sum_{i=0, \text{even}}^{n-1} (-2\xi)^i B_{ni}^q(t) - (-2\xi)^n C_{n0}^q(t) \Big|_{n \text{ even}}$$

$n = 2$

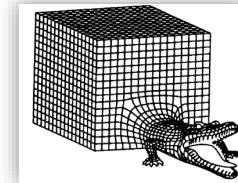


HL, Phys.Rev.Lett. 127 (2021) 18, 182001

# Nucleon Tomography

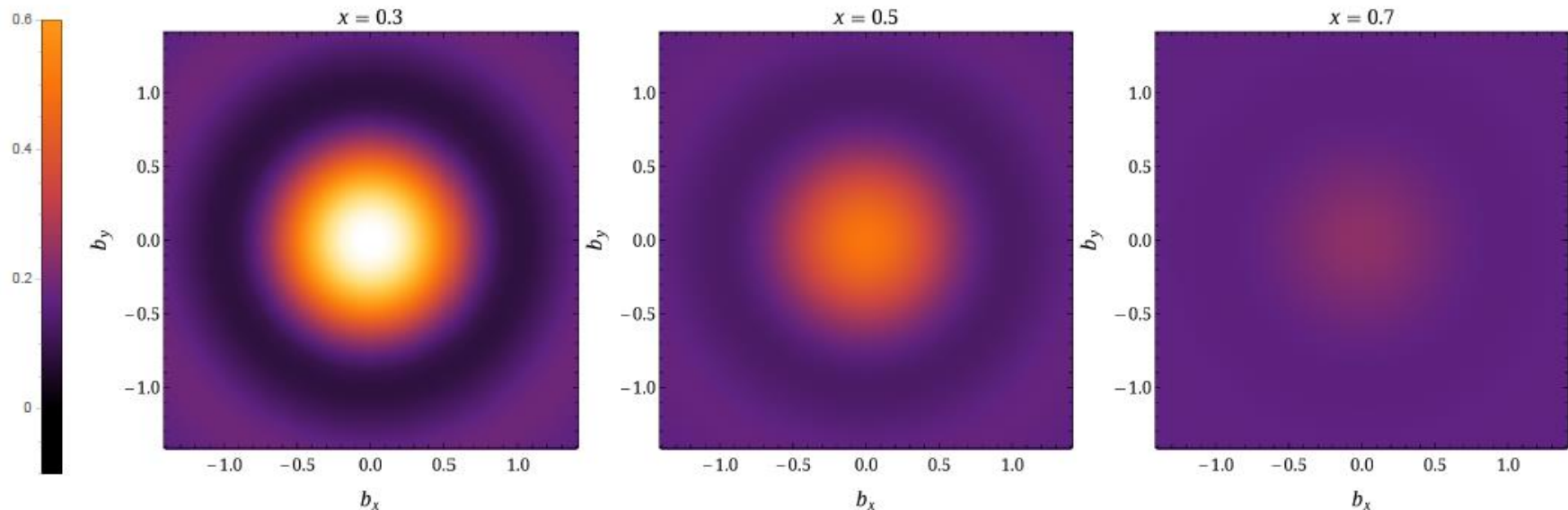
## § Nucleon GPD using quasi-PDFs at physical pion mass

- ∞ Lattice details: clover/2+1+1 HISQ  
0.09 fm, 135-MeV pion mass,  $P_z \approx 2$  GeV
- ∞  $\xi = 0$  isovector nucleon quasi-GPD results



finite-volume,  
discretization,

$$q(x, b) = \int \frac{d\vec{q}}{(2\pi)^2} H(x, \xi = 0, t = -\vec{q}^2) e^{i\vec{q} \cdot \vec{b}}$$



HL, Phys.Rev.Lett. 127 (2021) 18, 182001

# Challenges

## § Large momentum is essential

↪ With sufficient statistics nucleons may reach 5 GeV

## § Renormalization of linear divergence

↪ Wilson-line ops have linear divergences that must be subtracted

## § Methods for signal-to-noise improvement

↪ Gluonic observables, new ideas for large momentum

## § Inverse problems PDF extraction in SDF

↪ Remove the model/preconditioner-choice dependence

## § Reaching long-range correlations in LaMET

↪ For small- $x$  physics, new methods for calculating longer-range correlations must be developed

Whitepaper: Lattice QCD Calculations of Parton Physics, 2202.07193



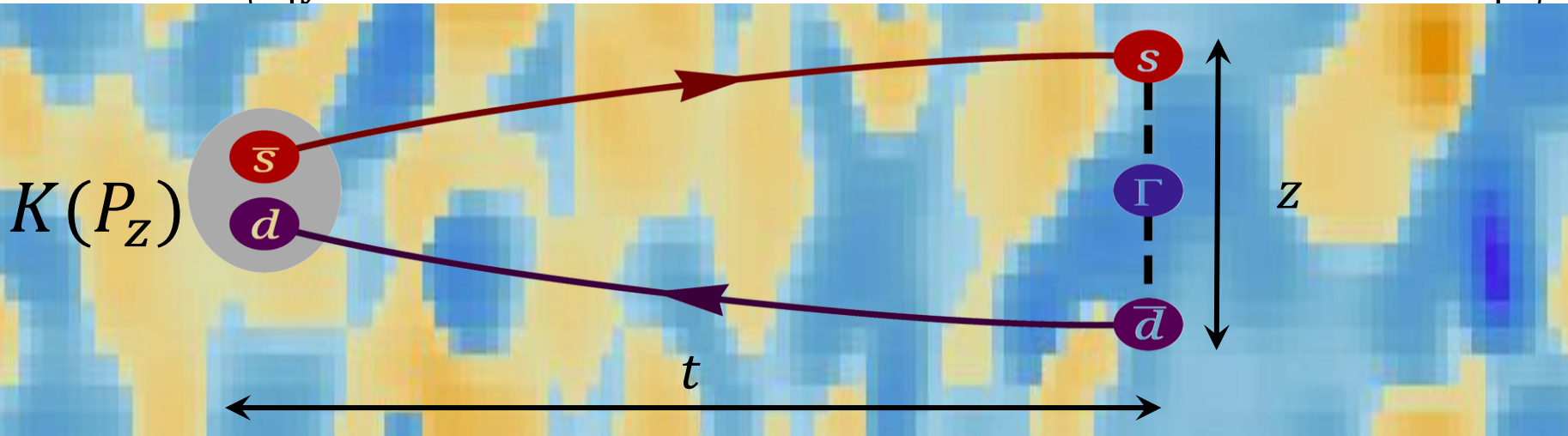
# Application on Inverse Problem



R. Zhang, C. Honkala, et al. (MSULat), 2005.13955

## Example: Pion/Kaon Distribution Amplitude

$$C_M^{DA}(z, P, t) = \left\langle 0 \left| \int d^3y e^{i\vec{P}\cdot\vec{y}} \bar{\psi}_1(\vec{y}, t) \gamma_z \gamma_5 U(\vec{y}, \vec{y} + z \hat{z}) \psi_2(\vec{y} + z \hat{z}, t) \bar{\psi}_2(0, 0) \gamma_5 \psi_1(0, 0) \right| 0 \right\rangle$$

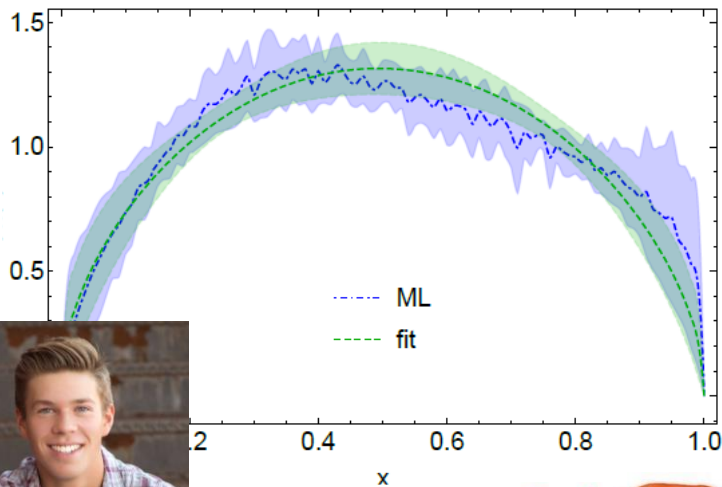


# Application on Inverse Problem



R. Zhang, C. Honkala, et al. (MSULat), 2005.13955

## Pion Distribution Amplitude

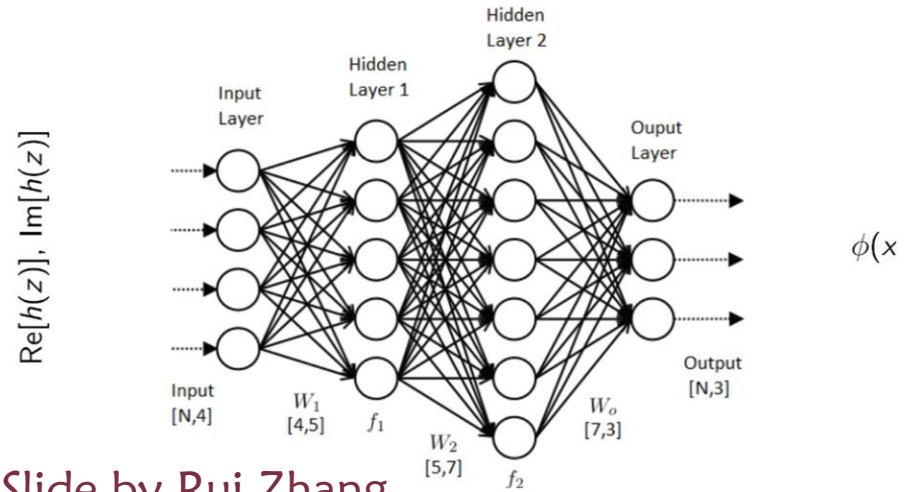


UG: Carson Honkala



### Machine Learning - A Promising Solution?

Machine learning models are effective in extracting complicated dependence of the output data on input data.



Slide by Rui Zhang

# Summary and Outlook

§ Exciting era using LQCD to study pion and nucleon structure

↪ Well-studied systematics → precision structures

§ Overcoming longstanding obstacles

↪ Bjorken- $x$  dependence of parton distributions are widely studied with LaMET and its variants

§ Precision and progress are limited on resources

↪ Challenges = new opportunities quantities

§ In the future

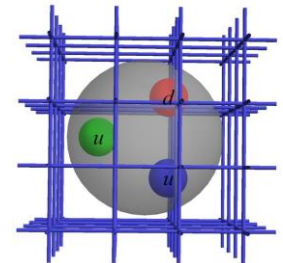
**Theory  
Input**



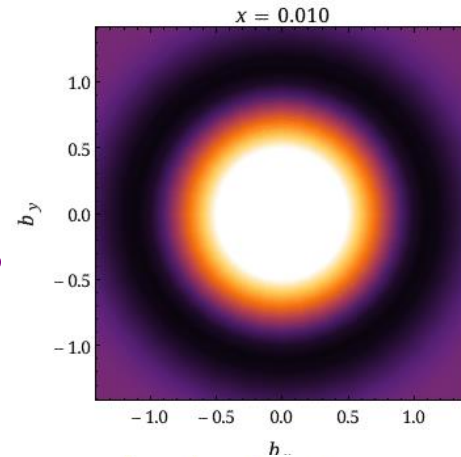
**Exp't  
Input**

**Global Analysis  
of GPDs**

+

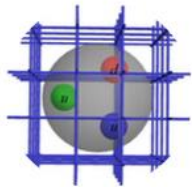
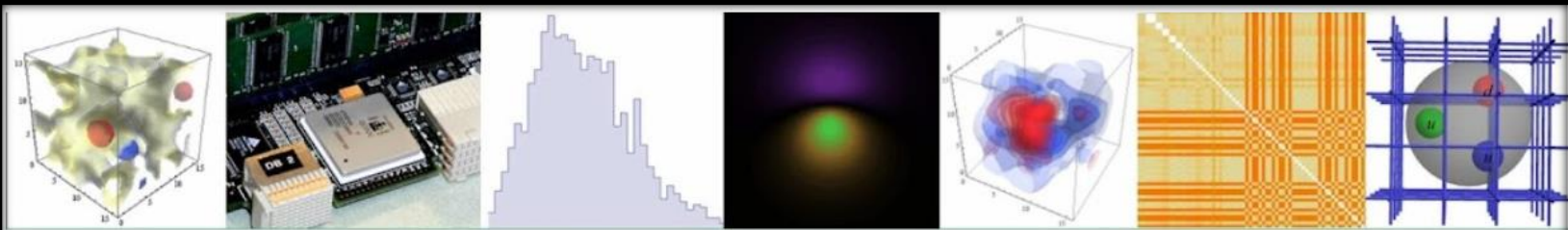


**EXCLAIM**



Thanks to MILC collaboration for sharing their 2+1+1 HISQ lattices & USQCD/NSF/DOE for computational resources  
The work of HL is sponsored by NSF under grant PHY 2209424 & 1653405, DOE under DE-SC0024053 & RCSA Cottrell Scholar Award

# Interested in learning more?



Lattice QCD Education

@latticeqcdeducation8264 381 subscribers 52 videos

Educational videos for people who want to get started doing lattice-QCD re... >



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**Jackknife Example**

§ Our complete data set:

$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
-------	-------	-------	-------	-------	-------	-------	-------	-------	----------	---------	-------

§ Our best estimate of the mean:

$$\bar{x} = \frac{f(x_1) + f(x_2) + f(x_3) + \dots + f(x_N)}{N}$$

§ Create new sets by dropping one value from each:

$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$
$x_1$	$x_2$	$x_3$	$x_4$	$x_5$	$x_6$	$x_7$	$x_8$	$x_9$	$x_{10}$	$\dots$	$x_N$

§ Now we have  $N$  new estimates of the mean:

$$\bar{x}_1 = \frac{f(x_2) + f(x_3) + \dots + f(x_N)}{N-1}, \bar{x}_2 = \frac{f(x_1) + f(x_3) + \dots + f(x_N)}{N-1}, \dots, \bar{x}_N = \frac{f(x_1) + f(x_2) + \dots + f(x_{N-1})}{N-1}$$

5:06

Jackknife Tutorial  
15K views · 3 years ago

Dynamics:

- Sum over elementary squares, "plaquettes"

$$U_p = U_{1,2} U_{2,3} U_{3,4} U_{4,1}$$

- like a "curl"  $\vec{\nabla} \times \vec{A} = \vec{B}$
- flux through corresponding plaquette:

$$S = \int d^4x (E^2 + B^2) \rightarrow \sum_p (1 - \frac{1}{3} \text{Re} \text{Tr} U_p)$$

27:35

Introduction to Lattice QCD (Michael Creutz):  
Lecture I  
5K views · 2 years ago

Potential has  $N_f^2 - 1$  "flat" directions

- one for each generator of  $SU(N_f)$

Excitations in the  $\pi$  direction don't raise energy

- pions become massless

28:36

Introduction to Lattice QCD (Michael Creutz):  
Lecture II  
1.1K views · 2 years ago

# *If time allows...*

## Learning by Doing!

§ Prior Python with Jupyter notebook experience required

↻ How many of you have not used these before?

§ Form groups of about 4 students

§ Introduce yourself to the other students in your group  
(5-10 mins)

↻ Name and preferred pronoun

↻ If you could hang out with any cartoon character,  
who would you choose and why?

↻ **Make sure you listen to each other!**

# Hands-on Exercise

## § Hands-on exercise

- ↪ Always starts with “pre-class” notebook before moving on to “in-class” one
- ↪ Set01: Jackknife analysis (Mon)
- ↪ Set02: Calculating proton masses (Mon)
- ↪ Set03: Extracting nucleon charges (Tue)



Every baby  
knows the

# scientific method!

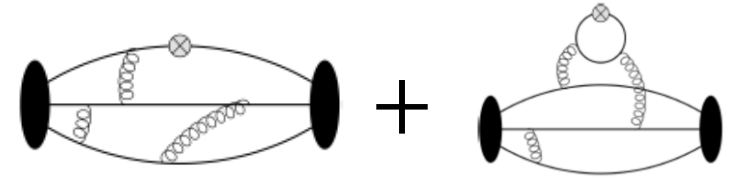


# *Backup Slides*



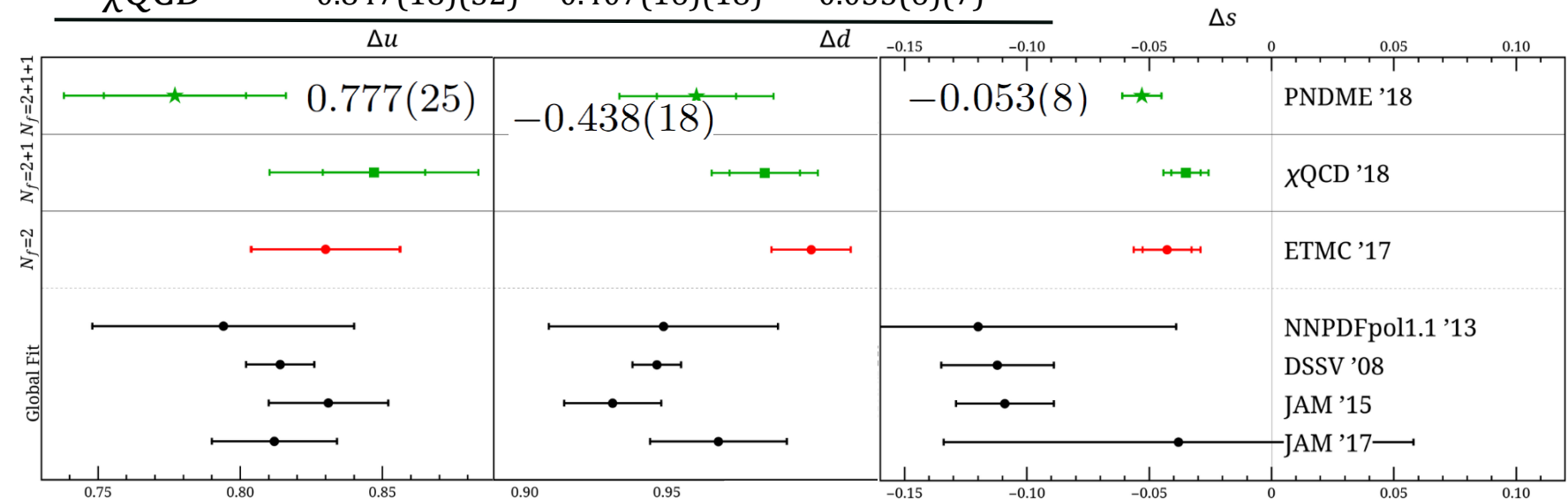
# Quark Spin Contribution

## § Total quark contributions



PNDME	$g_A^u \equiv \Delta u$	$g_A^d \equiv \Delta d$	$g_A^s \equiv \Delta s$
Connected	0.895(21)	-0.320(12)	
Disconnected	-0.118(14)	-0.118(14)	-0.053(8)
Sum	0.777(25)	-0.438(18)	-0.053(8)
ETMC	0.830(26)	-0.386(18)	-0.042(10)(2)
$\chi$ QCD	0.847(18)(32)	-0.407(16)(18)	-0.035(6)(7)

Difference caused by  $\Delta q^{\text{disc}}$



$$\sum_{q=u,d,s} \left(\frac{1}{2} \Delta q\right) = 0.143(31)$$

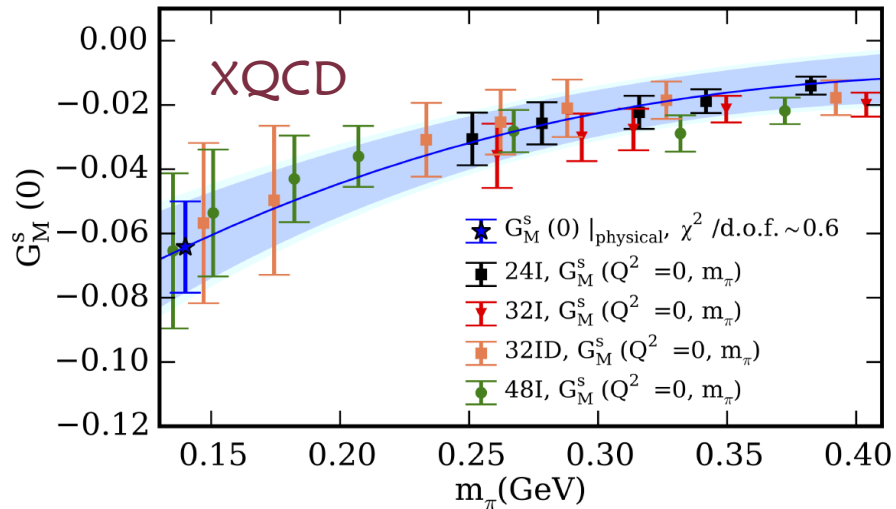
PNDME, 1806.09006, 1806.10604

The work of HL is supported by NSF CAREER Award under grant PHY 1653405

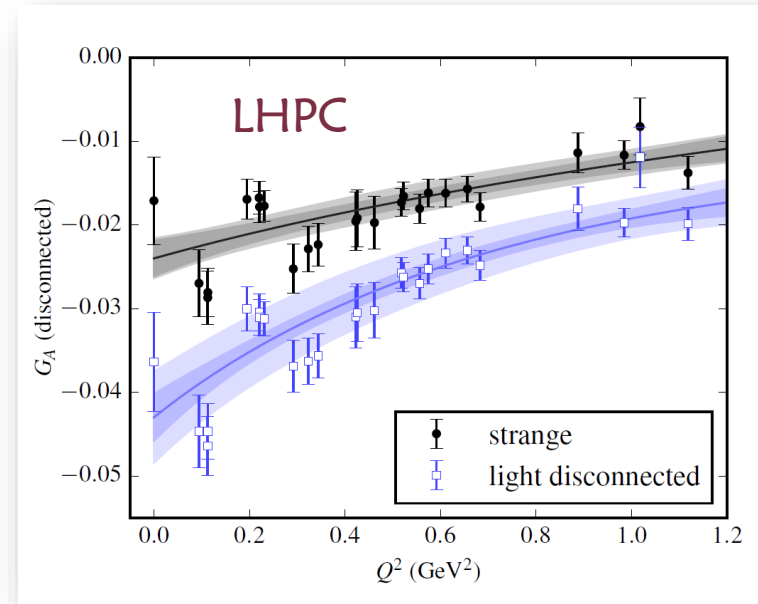


# Other Form Factors

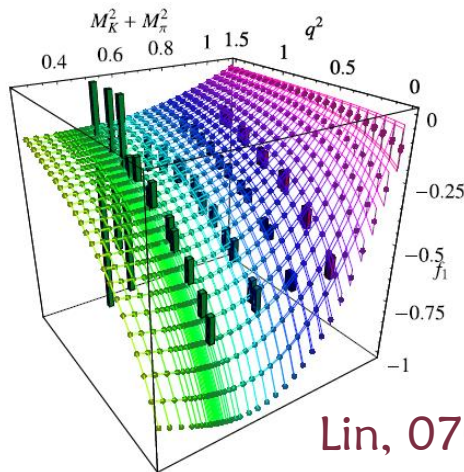
## § Toward flavor-dependent nucleon form factor



Sufian et al., Phys. Rev. Lett. 118, 042001 (2017)



Green et al., Phys. Rev. D 95, 114502 (2017)



## § Hyperon form factor

$$\begin{aligned} \bar{\nu}_\mu p &\rightarrow \mu^+ \Lambda^0 \\ \bar{\nu}_\mu n &\rightarrow \mu^+ \Sigma^- \\ \bar{\nu}_\mu p &\rightarrow \mu^+ \Sigma^0 \end{aligned}$$

Lin, 0707.3844 [hep-lat]

# First Lattice Charm PDF

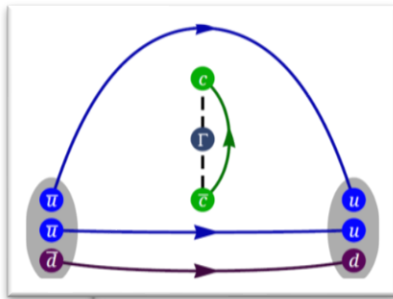
§ Large uncertainties in global PDFs

§ Results by MSULat/quasi-PDF method

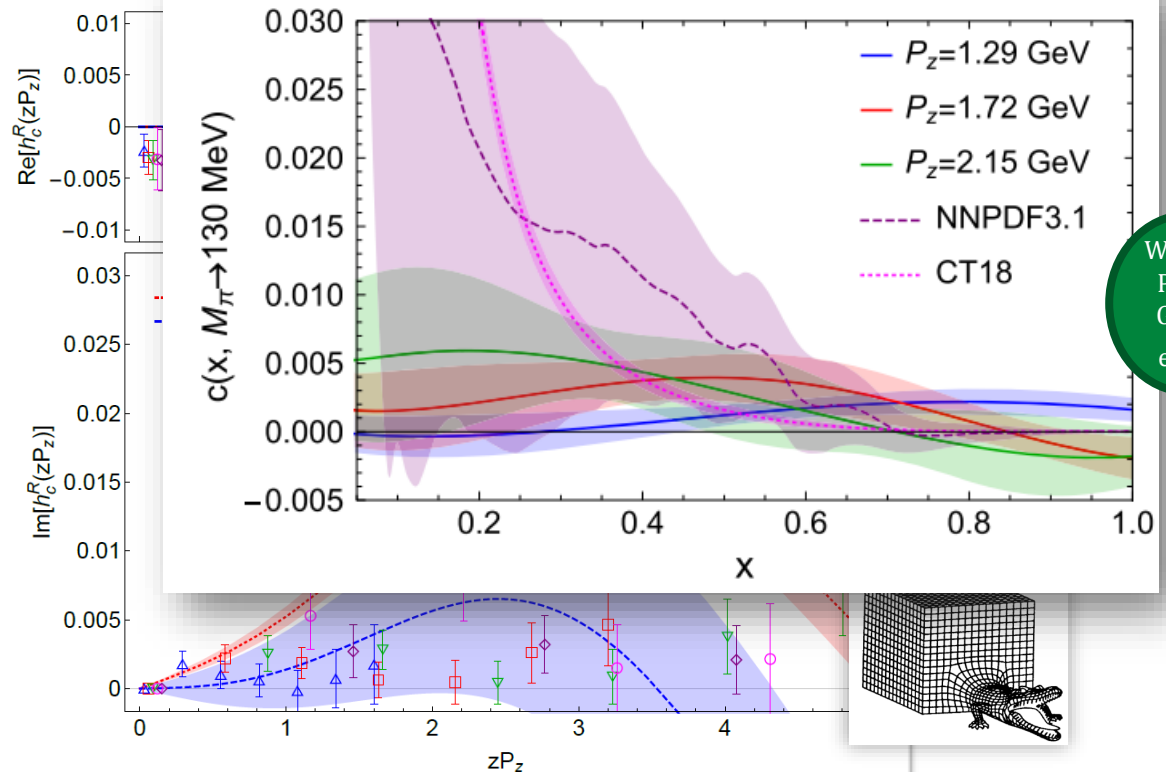
↻ Clover on 2+1+1 HISQ 0.12-fm 310-MeV QCD vacuum



2005.01124, R. Zhang et al (MSULat)



- suggest a symmetric  $c - \bar{c}$  distribution
- much smaller than strange PDF



Wanted PDFs, GPDs, etc...

# Moments of PDFs

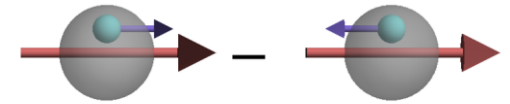
§ PDG-like rating system or average

§ LatticePDF Workshop

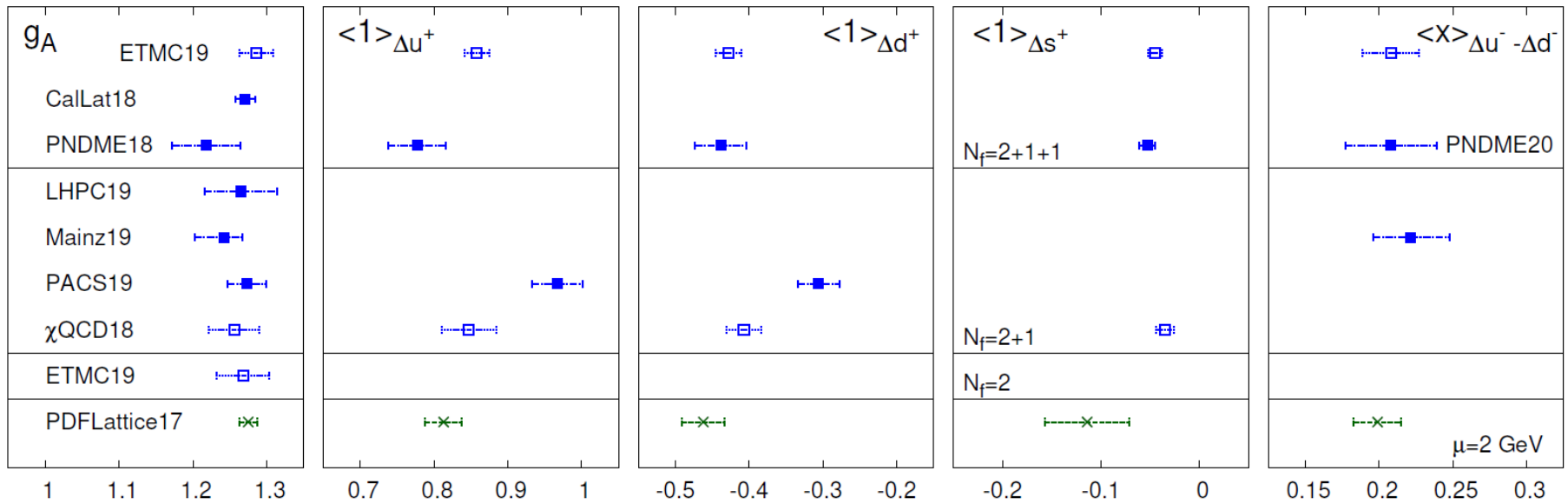
↻ Lattice representatives came together and devised a rating system

§ Recent lattice QCD/global fit status

$$\langle x^{n-1} \rangle_{\Delta q} = \int_{-1}^1 dx x^{n-1} \Delta q(x)$$



LatticePDF Report, 1711.07916,2006.08636

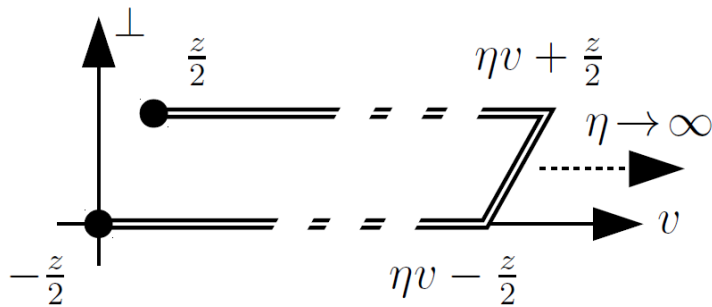


# Orbital Angular Momentum

## § Two definitions: Ji vs Jaffe & Manohar

$$\vec{L}_q^{\text{Ji}} = \int d^3x q^\dagger [\vec{x} \times i\vec{D}] q$$

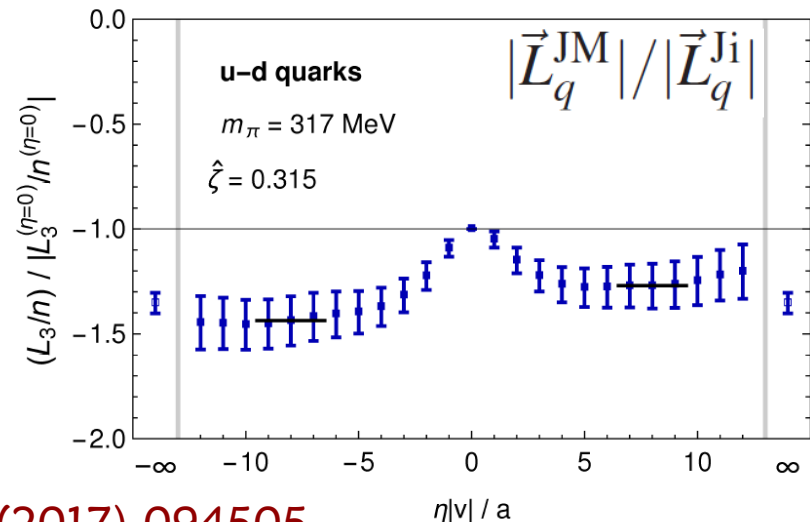
$$\vec{L}_q^{\text{JM}} = \int d^3x q^\dagger [\vec{x} \times i\vec{\nabla}] q,$$



- $\eta = 0$  gives Ji's OAM
- Staple  $\eta \rightarrow \infty$  gives Jaffe-Manohar OAM
- Difference is accumulated torque from final-state interaction

## § First result carried out by M. Engelhardt

2+1f clover at 518-MeV pion mass



Phys. Rev. D95 (2017) 094505

# New Interactions in Beta Decays

§ Neutron beta decay could be sensitive to new interactions:

$$H_{\text{eff}} = G_F \left( J_{V-A}^{\text{lept}} \times J_{V-A}^{\text{quark}} + \sum_i \varepsilon_i^{\text{BSM}} \hat{O}_i^{\text{lept}} \times \hat{O}_i^{\text{quark}} \right)$$

$$\hat{O}_S = \bar{u}d \times \bar{e}(1 - \gamma_5)\nu_e \quad \rightarrow \quad g_S = \langle n | \bar{u}d | p \rangle$$

$$\hat{O}_T = \bar{u}\sigma_{\mu\nu}d \times \bar{e}\sigma^{\mu\nu}(1 - \gamma_5)\nu_e \quad \rightarrow \quad g_T = \langle n | \bar{u}\sigma_{\mu\nu}d | p \rangle$$

∴  $\varepsilon_S$  and  $\varepsilon_T$  are related to the masses of the new TeV-scale particles  
(just like  $G_F \propto M_{W,Z}^{-2}$ )

# New Interactions in Beta Decays

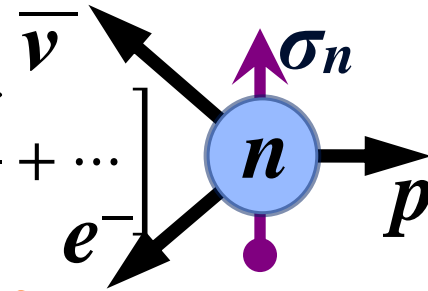
§ Neutron beta decay could be related to new interactions:

$$H_{\text{eff}} = G_F \left( J_{V-A}^{\text{lept}} \times J_{V-A}^{\text{quark}} + \sum_i \varepsilon_i^{\text{BSM}} \hat{O}_i^{\text{lept}} \times \hat{O}_i^{\text{quark}} \right)$$

∞  $\varepsilon_S$  and  $\varepsilon_T$  are related to the masses of the new TeV-scale particles

∞ Parameters sensitive to new physics

$$d\Gamma \propto F(E_e) \left[ 1 + A \frac{\vec{\sigma}_n \cdot \vec{p}_e}{E_e} + b \frac{m_e}{E_e} + \left( B_0 + B_1 \frac{m_e}{E_e} \right) \frac{\vec{\sigma}_n \cdot \vec{p}_\nu}{E_\nu} + \dots \right]$$



**Fierz interference term:**

Deviations from the leading-order  $e^-$  spectrum

Energy-dependent part of the **neutrino asymmetry parameter** with neutron spin

$$\{b, B\}_{\text{BSM}} = f_0(\varepsilon_{S,T} g_{S,T})$$

Precision LQCD input  
( $m_\pi \approx 140$  MeV,  $a \rightarrow 0$ )

$$\varepsilon_{S,T} \propto \Lambda_{S,T}^{-2}$$

# New Interactions in Beta Decays

§ Neutron beta decay could be related to new interactions:

$$H_{\text{eff}} = G_F \left( J_{V-A}^{\text{lept}} \times J_{V-A}^{\text{quark}} + \sum_i \varepsilon_i^{\text{BSM}} \hat{O}_i^{\text{lept}} \times \hat{O}_i^{\text{quark}} \right)$$

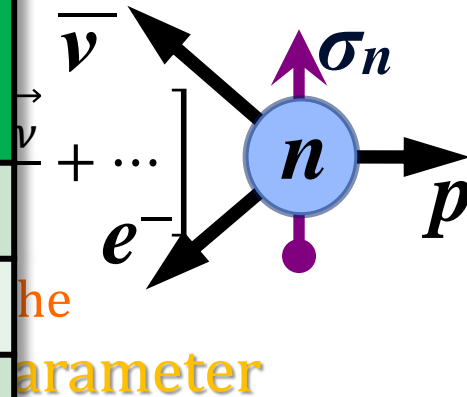
∞  $\varepsilon_S$  and  $\varepsilon_T$  are related to the masses of the new TeV-scale particles

∞ Parameters

$$d\Gamma \propto F(E_e)$$

Fierz inter  
Deviations f  
leading-ord

Ongoing and Future Experiments	Expected Precision
UCNb & UCNB at LANL	$10^{-3}$ to $10^{-4}$
Nab at ORNL	$10^{-3}$
FRMII in Munich, ...	
CENPA ${}^6\text{He}(b_{\text{GT}})$	$10^{-3}$ to $10^{-4}$



$$\{O, D\}_{\text{BSM}} \sim \mathcal{O}(\varepsilon_S, T, \gamma_S, T)$$

$$\varepsilon_{S,T} \propto \Lambda_{S,T}^{-2}$$

precision LQCD input  
( $m_\pi \approx 140$  MeV,  $a \rightarrow 0$ )

# Beta Decays & BSM

§ Given precision  $g_{S,T}$  and  $O_{\text{BSM}}$ , predict new-physics scales

Low-Energy

Expt  $\rightarrow$

$$O_{\text{BSM}} = f_O(\epsilon_{S,T} g_{S,T})$$

Precision LQCD input  
( $m_\pi \rightarrow 140$  MeV,  $a \rightarrow 0$ )  $\leftarrow$

$$\epsilon_{S,T} \propto \Lambda_{S,T}^{-2}$$

Upcoming precision

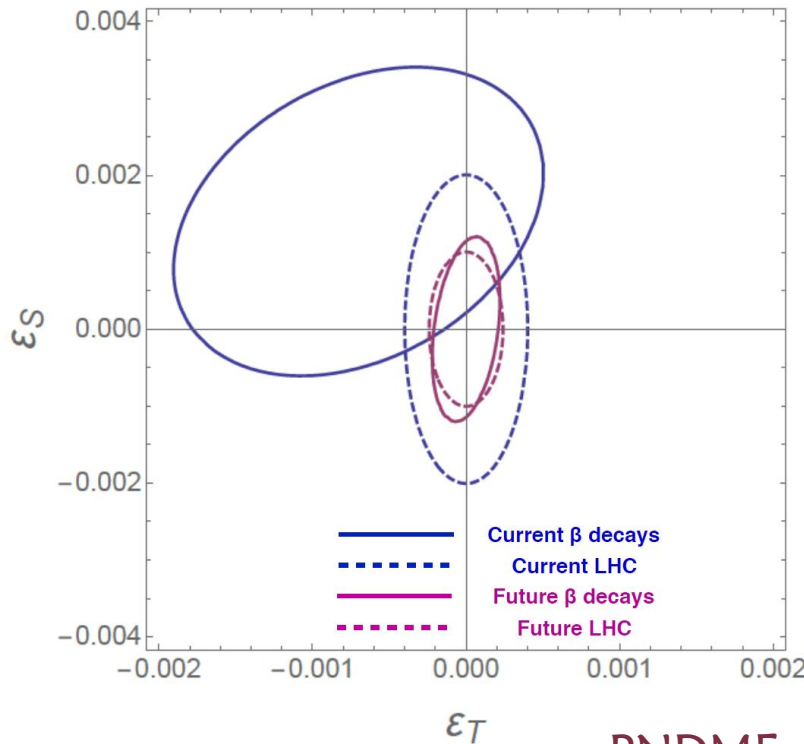
low-energy experiments

LANL/ ORNL UCN neutron  
decay exp't

$$|B_1 - b|_{\text{BSM}} < 10^{-3}$$

$$|b|_{\text{BSM}} < 10^{-3}$$

CENPA:  ${}^6\text{He}(b_{\text{GT}})$  at  $10^{-3}$



Plots by Vincenzo Cirigliano

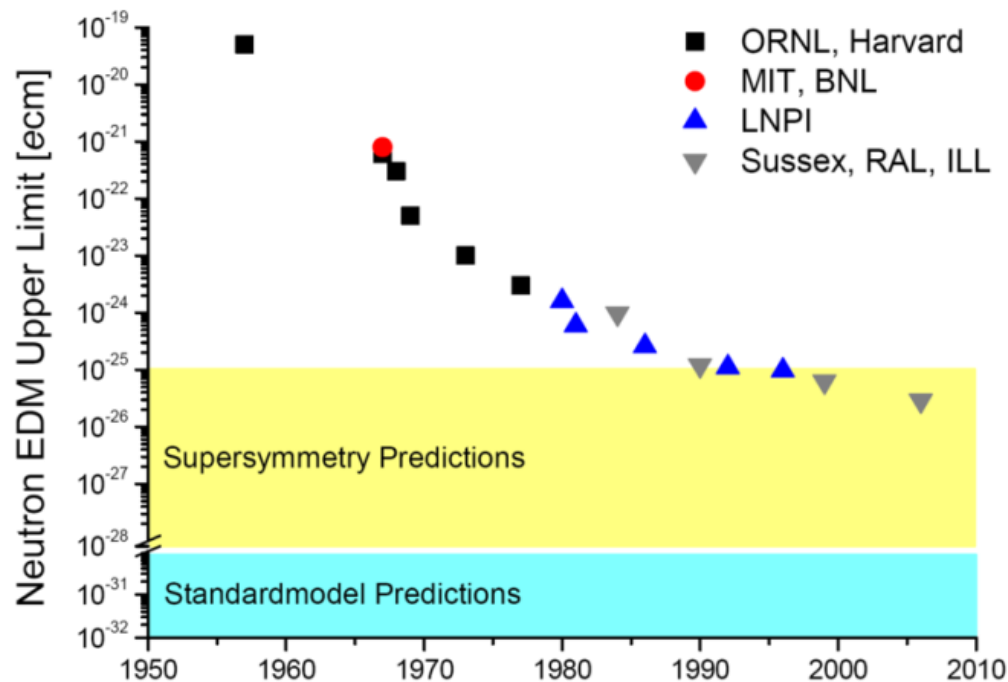
PNDME, PRD85 054512 (2012);  
1306.5435; 1606.07049; 1806.09006



# Electric Dipole Moment

## § Why do we care?

- ↻ CP-violating effect  $\Rightarrow$  Key ingredient for baryogenesis  
 $\Rightarrow$  Why matter exists
- ↻ Extremely small in SM:  $\approx 10^{-31}$  e-cm (expect to probe  $10^{-28}$  soon)
- ↻ Good candidate to constrain BSM models



# Electric Dipole Moment

## § Why do we care?

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- ↻ Good candidate to constrain BSM models

## § Lattice community are working on various contributions

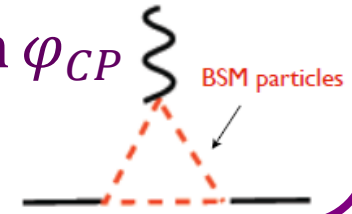
### § Lagrangian

$$L = L_{\text{QCD}}^{\text{CP Even}} + L_{\Theta} + L_{\text{quark}}^{\text{dim-5}} + L_{\text{chromo-quark}}^{\text{dim-5}} + \dots$$

- ↻ If experiment sees signal before SM background  
 $\Rightarrow$  new physics  
 $\Rightarrow$  quark EDM (our focus here)

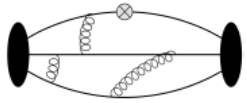
Induced by a variety of BSM scenarios

$$d_i \propto \frac{m_i}{\Lambda^2} \sin \varphi_{CP}$$



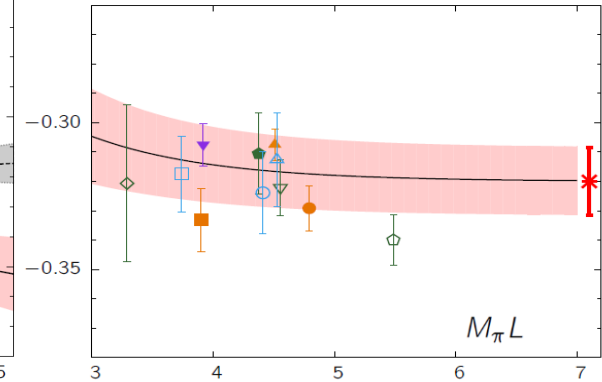
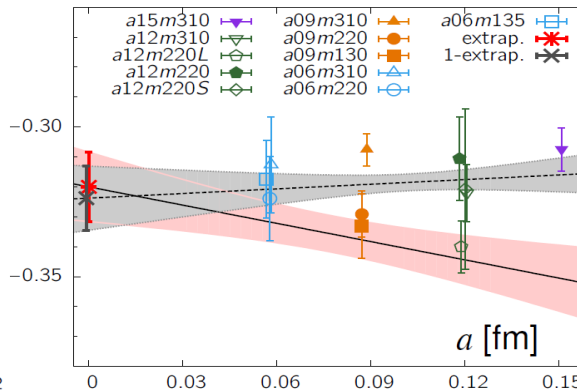
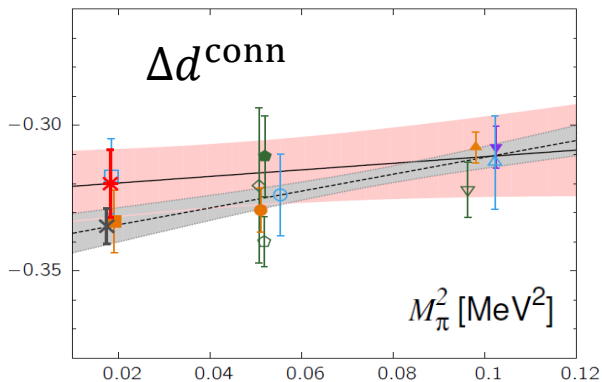
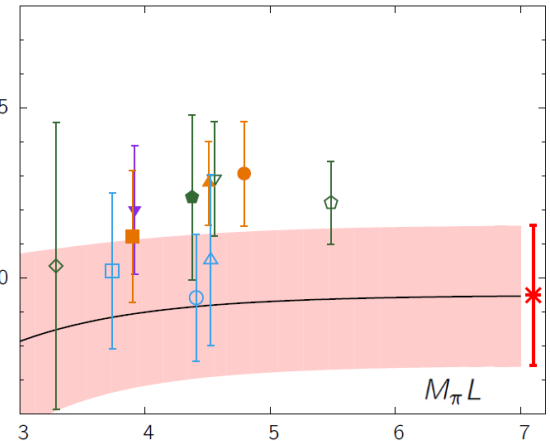
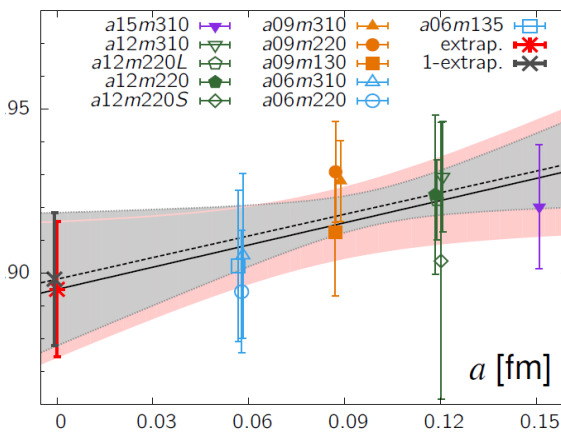
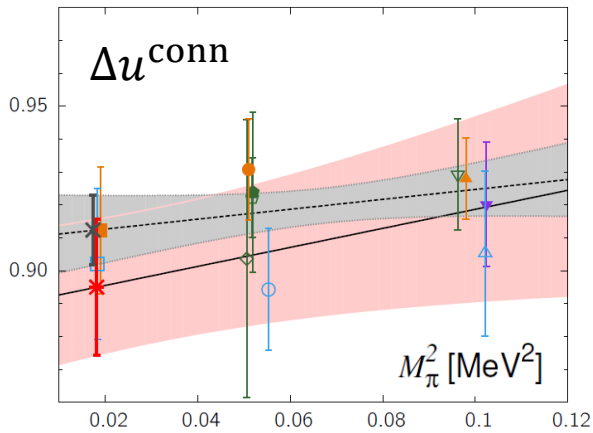
# Quark Spin

## § Up and down quark “connected” contribution



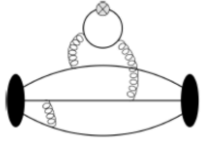
PNDME, 1806.09006, 1806.10604

$$\Delta q(a, m_\pi, L) = c_1 + c_2 m_\pi^2 + c_3 a + c_4 e^{-m_\pi L}$$



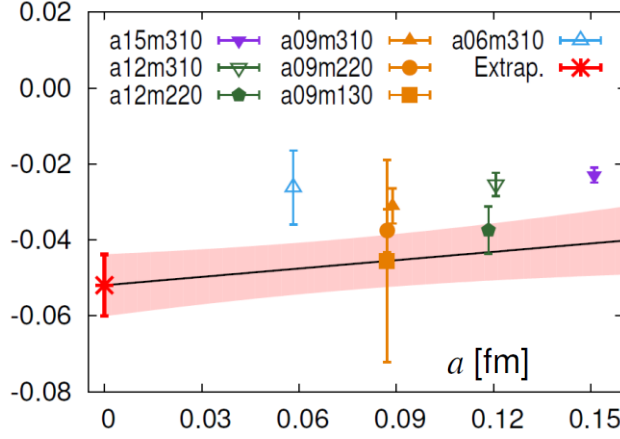
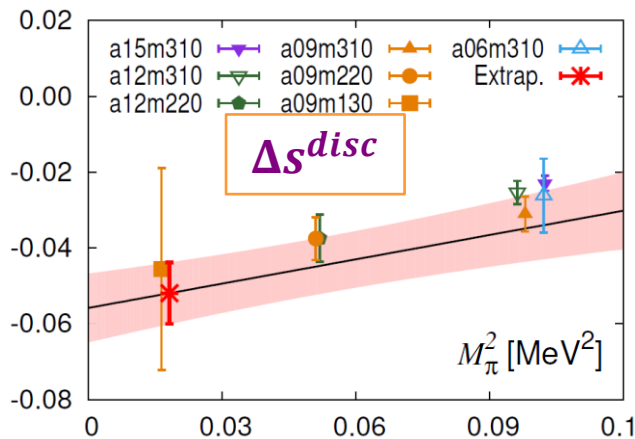
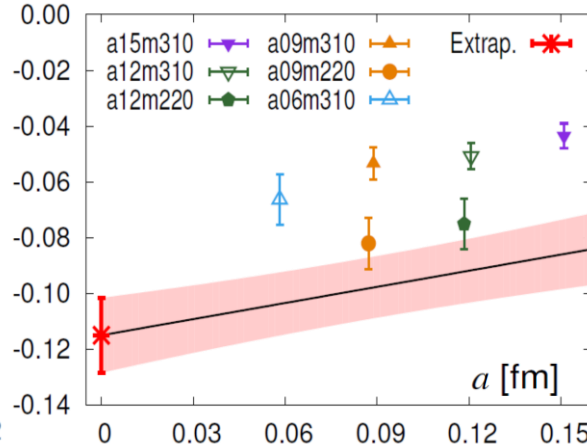
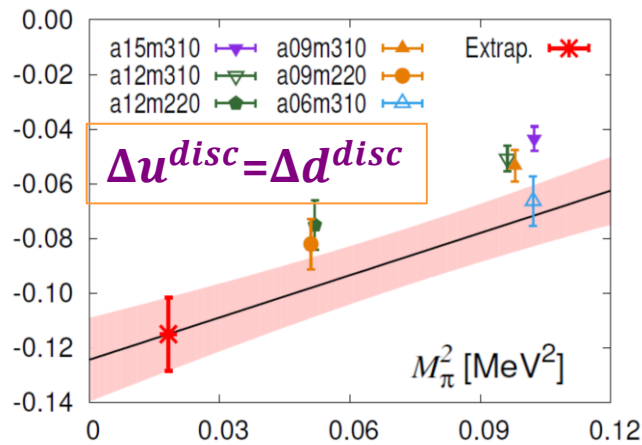
# Quark Spin

## § Up and down quark “disconnected” contribution



PNDME, 1806.09006, 1806.10604

$$\Delta q^{\text{disc}} = c_1 + c_2 m_\pi^2 + c_3 a + c_4 e^{-m_\pi L}$$



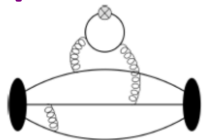
Anticipated pion-mass dependence

Unexpectedly strong lattice-spacing dependence!

Calculation at  $a \approx 0.09$  fm can have 50% change in  $\Delta u^{\text{disc}}$

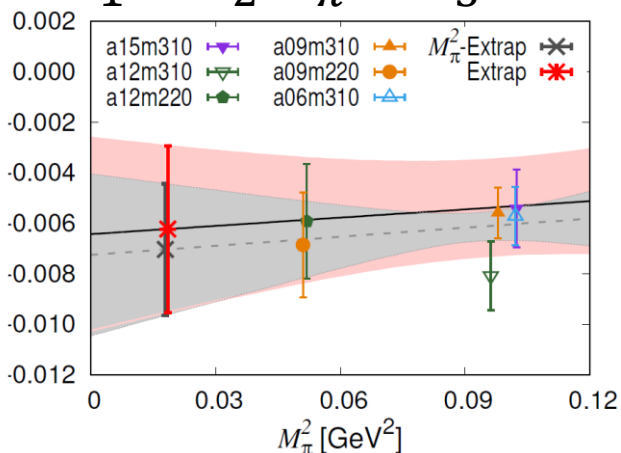
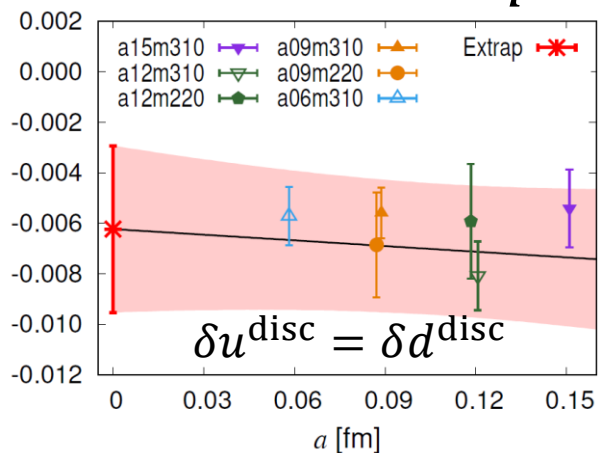
# $u/d/s$ Tensor Charges

## § Up and down quark “disconnected” contribution



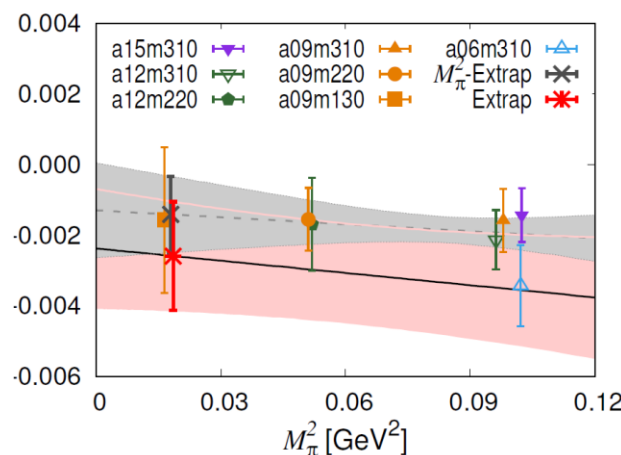
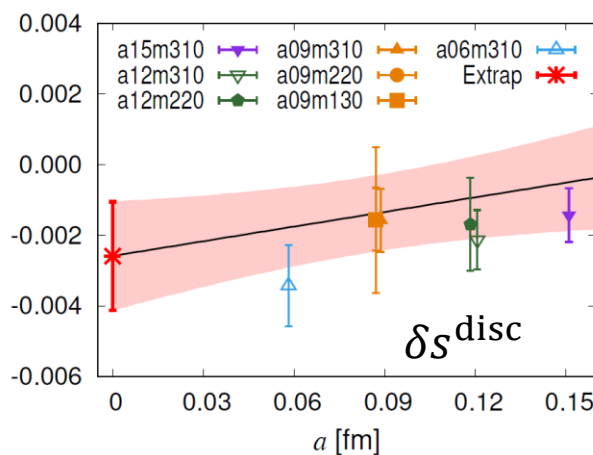
PNDME, 1806.09006, 1808.07597

$$\delta q^{\text{disc}} = c_1 + c_2 m_\pi^2 + c_3 a + c_4 e^{-m_\pi L}$$



First time in LQCD

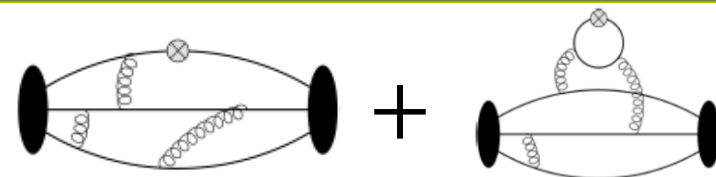
Mild dependence on  $a$  and pion mass!



Yong-Chull Jang

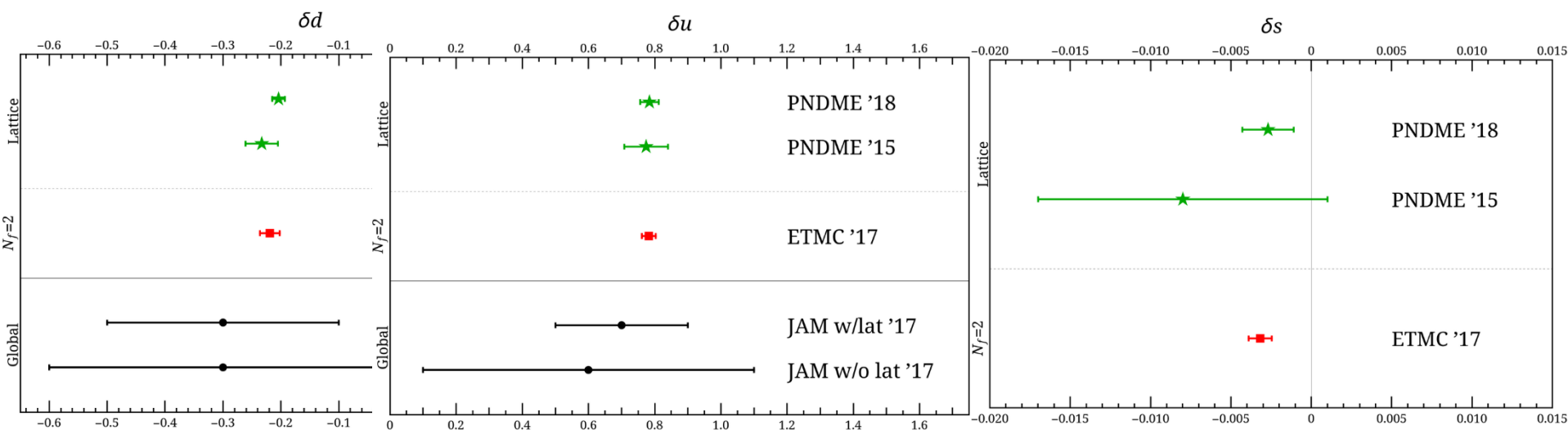
# $u/d/s$ Tensor Charges

## § Sum up both contributions



	$g_T^u$	$g_T^d$	$g_T^s$
Connected	0.790(27)	-0.198(10)	
Disconnected	-0.0064(33)	-0.0064(33)	-0.0027(16)
PDNME'18 (Sum)	0.784(28)	-0.204(11)	-0.0027(16)
ETMC'17 [14]	0.782(21)	-0.219(17)	-0.00319(72)
PNDME'15 [5]	0.774(66)	-0.233(28)	0.008(9)

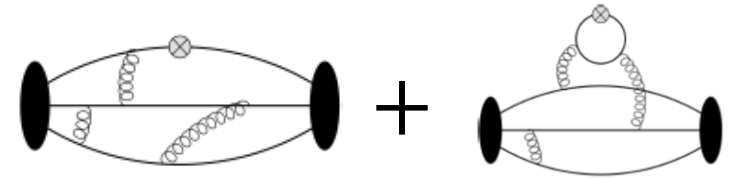
Calculation from one lattice ensemble only  
No cont. extrapolation errors



PNDME, 1806.09006, 1808.07597

# $u/d/s$ Tensor Charges

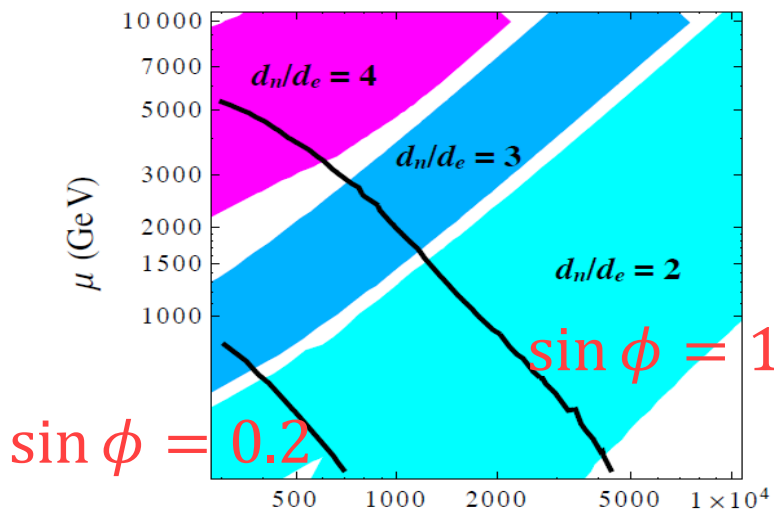
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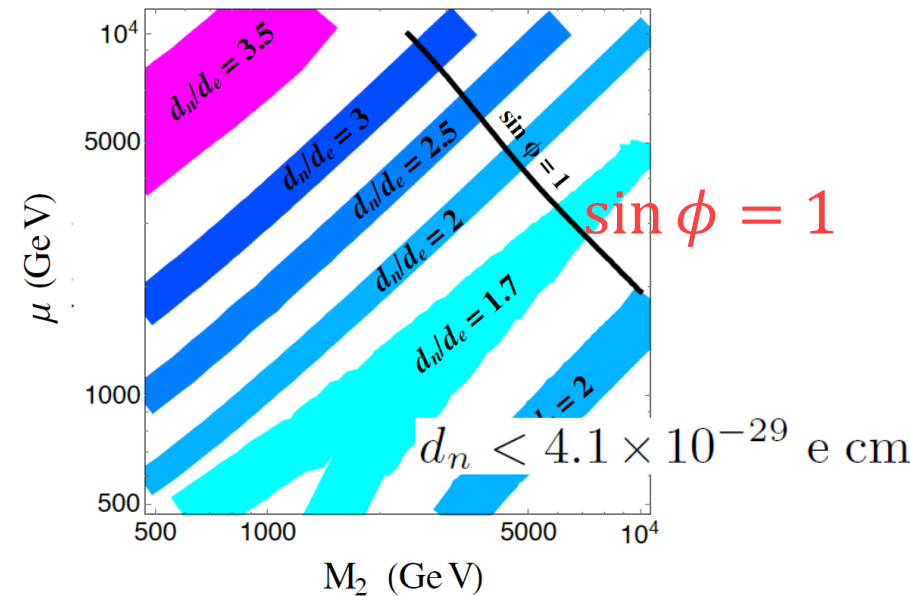
Calculation from one lattice ensemble only  
No cont. extrapolation errors

Set limits for nEDM in split-SUSY scenario with gaugino mass unification



$$d_n < 4 \times 10^{-28} e \cdot \text{cm} \quad M_2 \text{ (GeV)}$$

PNDME, 1506.04196; 1506.06411



$\sin \phi = 1$

$$d_n < 4.1 \times 10^{-29} e \text{ cm}$$

PNDME, 1806.09006, 1808.07597

# Precision Nucleon Couplings

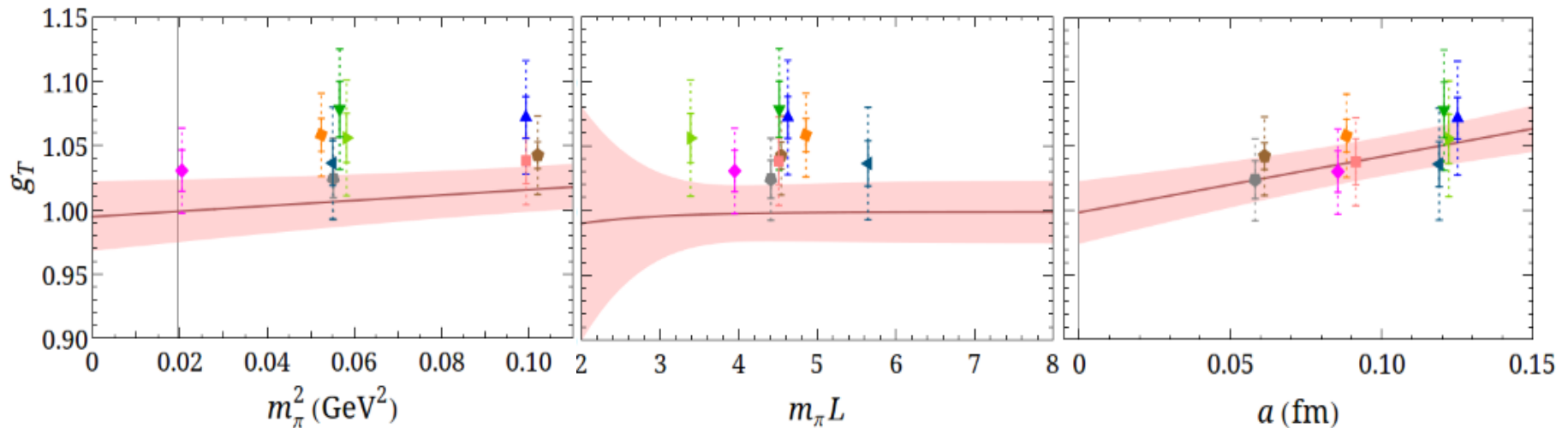
§  $g_T$ : zeroth moment of transversity  $\Gamma = \sigma_{\mu\nu}$

§ A state-of-the-art calculation (PNDME)

$$g_T = \int_{-1}^1 dx \delta q(x)$$

↻ Extrapolate to the physical limit

$$g_T(a, m_\pi, L) = c_1 + c_2 m_\pi^2 + c_3 a + c_4 e^{-m_\pi L}$$



First extrapolation to the physical limit  
of a nucleon matrix element!

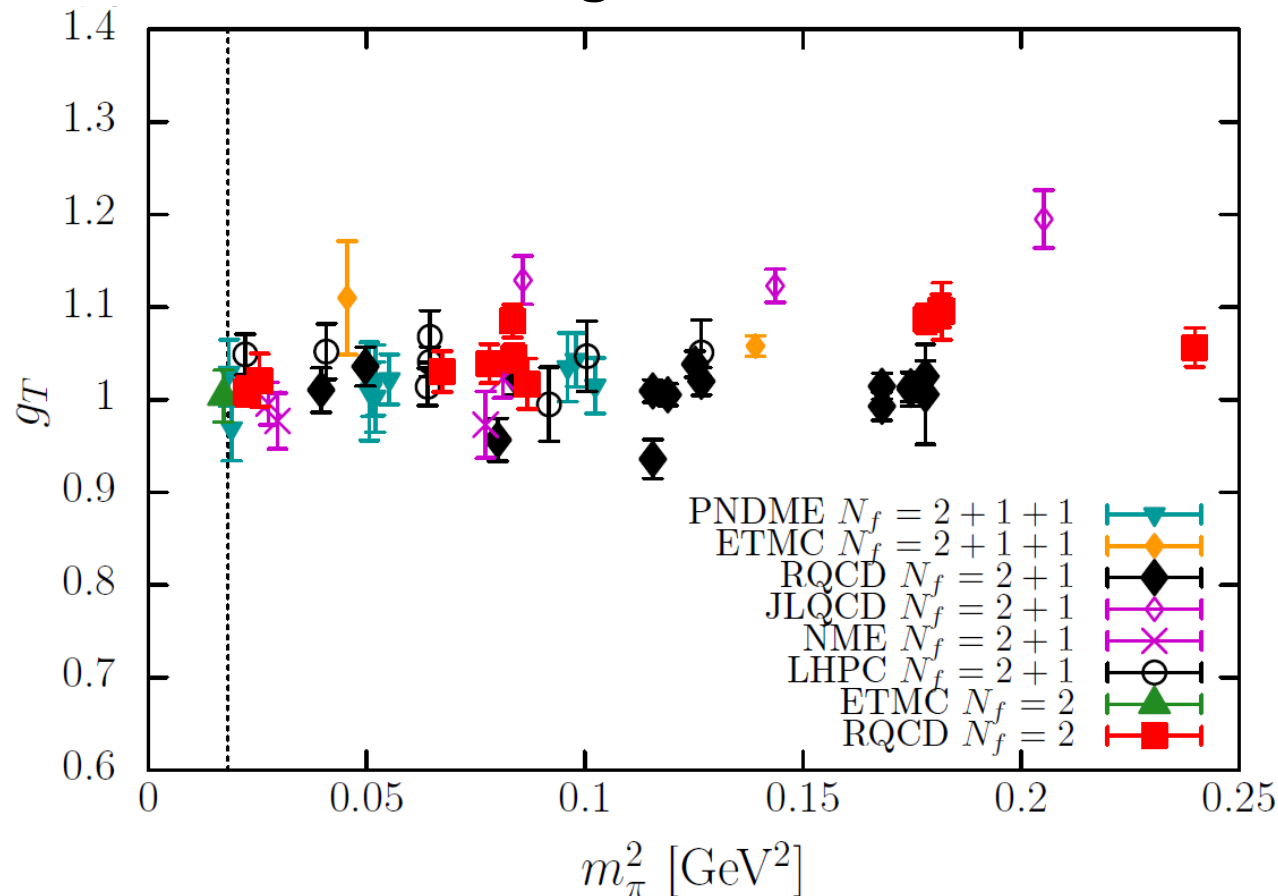


# Precision Nucleon Couplings

§ Usually more than one LQCD calculation

∞ For example, tensor charge

∞ Lattice results should agree in the continuum limit



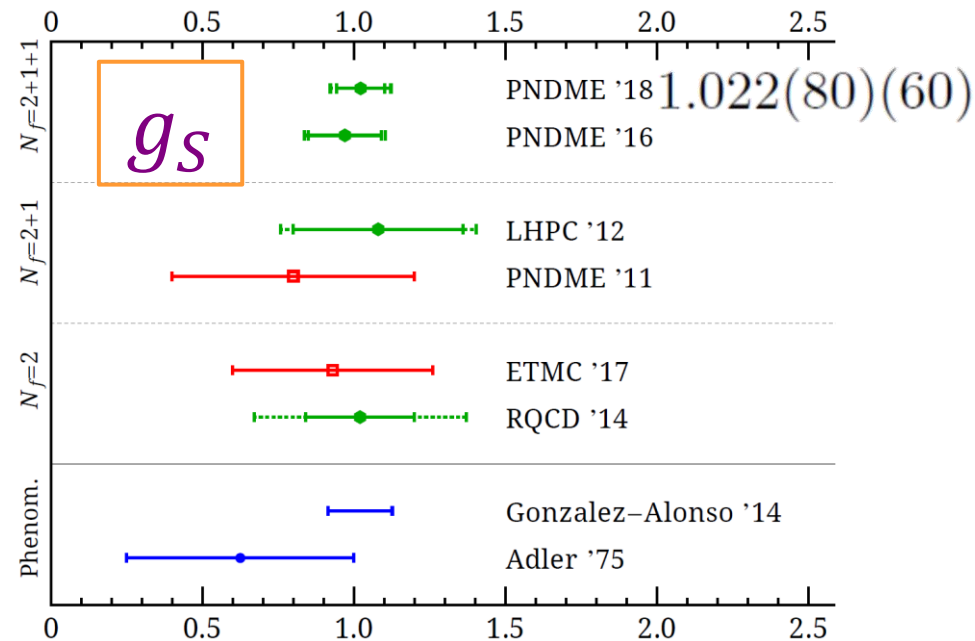
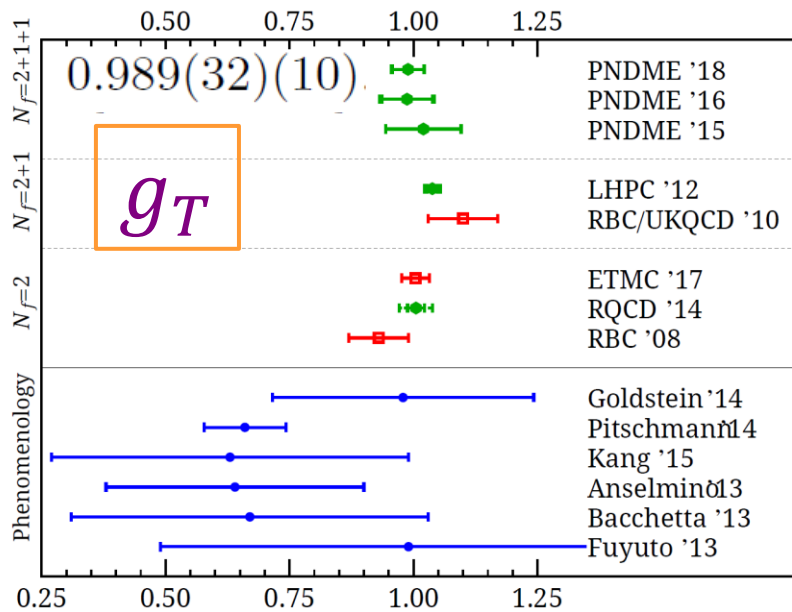
# Precision Nucleon Couplings

FLAG rating system PNDME, 1506.06411; 1606.07049

New: excited-state rating

Collaboration	Ref.	publication status	$N_f$	chiral extrapolation	continuum extrapolation	finite volume	excited state	renormalization	$g_T$
PNDME'15	This work	P	2+1+1	★	★	★	★	★	1.020(76) <sup>a</sup>
ETMC'13	[30]	C	2+1+1	■	○	○	■	★	1.11(3) <sup>b</sup>
LHPC'12	[28]	A	2+1	★	○	★	○	★	1.037(20) <sup>c</sup>
RBC/UKQCD'10	[29]	A	2+1	○	■	★	★	★	1.10(7) <sup>d</sup>
RQCD'14	[31]	P	2	★	★	★	○	★	1.005(17)(29)) <sup>e</sup>
ETMC'13	[30]	C	2	★	■	○	■	○	1.114(46) <sup>f</sup>
RBC'08	[32]	P	2	■	■	★	■	★	0.93(6) <sup>g</sup>

PNDME, 1806.09006

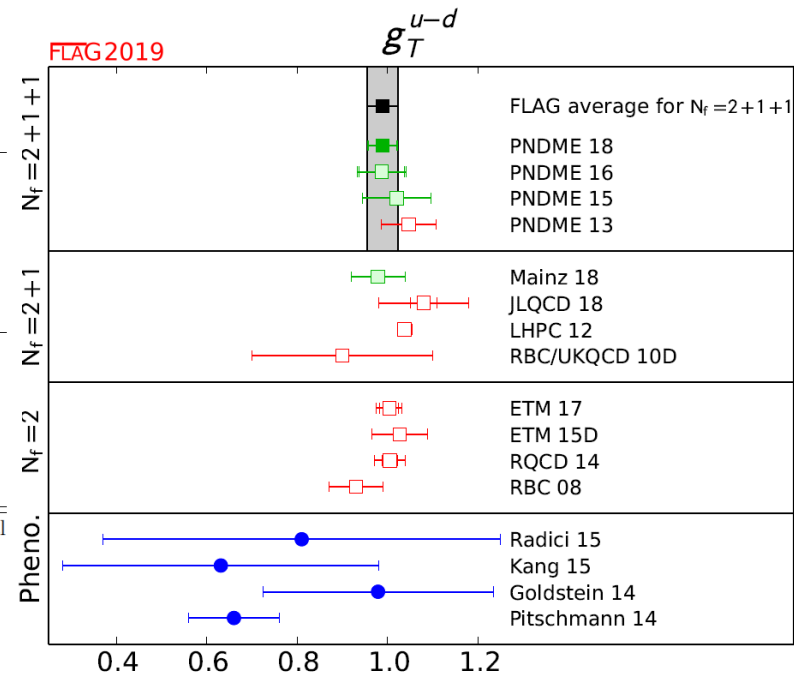


# FLAG 2019

§ Finally adopted by FLAG! <https://arxiv.org/pdf/1902.08191.pdf>

Collaboration	Ref.	$N_f$	publication status	continuum extrapolation	chiral extrapolation	finite volume	renormalization	excited states	$g_T^{u-d}$
PNDME 18	[84]	2+1+1	A	★ <sup>‡</sup>	★	★	★	★	0.989(32)(10)
PNDME 16	[830]	2+1+1	A	○ <sup>‡</sup>	★	★	★	★	0.987(51)(20)
PNDME 15	[828, 829]	2+1+1	A	○ <sup>‡</sup>	★	★	★	★	1.020(76)
PNDME 13	[827]	2+1+1	A	■ <sup>‡</sup>	■	★	★	★	1.047(61)
Mainz 18	[915]	2+1	C	★	○	★	★	★	0.979(60)
JLQCD 18	[839]	2+1	A	■	○	○	★	★	1.08(3)(3)(9)
LHPC 12	[920]	2+1	A	■ <sup>‡</sup>	★	★	★	★	1.038(11)(12)
RBC/UKQCD 10D	[834]	2+1	A	■	■	○	★	■	0.9(2)
ETM 17	[826]	2	A	■	○	○	★	★	1.004(21)(2)(19)
ETM 15D	[822]	2	A	■	○	○	★	★	1.027(62)
RQCD 14	[819]	2	A	○	★	★	★	■	1.005(17)(29)
RBC 08	[918]	2	A	■	■	■	★	■	0.93(6)

<sup>‡</sup> The rating takes into account that the action is not fully  $O(a)$  improved by requiring an additional lattice spacing.



# Total Quark Intrinsic Spin

Not “equal”: systematics are different?

- PNDME: 0.143(31)(36) (2+1+1 flavor clover-on-HISQ)
- ETMC: 0.201(17)(5) (2 flavor twisted mass)
- $\chi$ QCD: 0.202(13)(19) (2+1 flavor overlap-on-Domain Wall)

	$g_A^{u-d}$	$a \rightarrow 0$	$M_\pi$ MeV	$M_\pi L$	$Z_A$
PNDME $N_f = 2+1+1$	1.218(25)(30)	Yes 11 ensembles 0.15 – 0.06 fm	135 220 310	3.3 – 5.5	Assume $Z_A^S = Z_A^{NS}$
ETMC $N_f = 2$	1.212(40)	0.094 fm	130	2.93	Checked $Z_A^S = Z_A^{NS}$
$\chi$ QCD $N_f = 2+1$	1.254(16)(30)	“No” a variation 0.143 fm 0.11 fm 0.083 fm	171 337 302	3.97 4.53 4.06	Checked $Z_A^S = Z_A^{NS}$

In perturbation theory  $Z_A^S \neq Z_A^{NS}$  at 2 loops . ETMC &  $\chi$ QCD show a  $\sim 1\%$  difference

Slide from Rajan Gupta @ Spin 2018

# Proton Mass Decomposition

$$M = -\langle T_{44} \rangle = \langle H_m \rangle + \langle H_E \rangle (\mu) + \langle H_g \rangle (\mu) + \frac{1}{4} \langle H_a \rangle$$

$$M = -\langle T_{\mu\mu} \rangle = \langle H_m \rangle + \langle H_a \rangle$$

X. Ji, PRL74:1071 (1995)

quark mass

$$\langle H_m \rangle = \sum_{u,d,s,\dots} \int d^3x m \bar{\psi} \psi$$

$$\langle x \rangle_{q,g} = \int_0^1 dx x f_{q,g}(x) = -\frac{\langle N | \frac{4}{3} \bar{T}_{44}^{q,g} | N \rangle}{M \langle N | N \rangle}$$

quark energy

$$\langle H_E \rangle = \frac{3}{4} (\langle x \rangle_q M - \langle H_m \rangle)$$

$$\bar{T}_{44}^q = \int d^3x \bar{\psi} \frac{1}{2} \left( \gamma_4 \vec{D}_4 - \frac{1}{4} \sum_{i=0,1,2,3} \gamma_i \vec{D}_i \right) \psi$$

glue energy

$$\langle H_g \rangle = \frac{3}{4} \langle x \rangle_g M$$

$$\bar{T}_{44}^g = \int d^3x \frac{1}{2} (E^2 - B^2)$$

anomaly

$$\langle H_a \rangle = \langle H_g^a \rangle + \langle H_m^Y \rangle$$

$$\langle H_g^a \rangle = \int d^3x \frac{-\beta(g)}{g} (E^2 + B^2)$$

$$\langle H_m^Y \rangle = \sum_{u,d,s,\dots} \int d^3x \gamma_m m \bar{\psi} \psi$$

## Ingredients

- ◆ proton mass
- ◆ scalar charge
- ◆ momentum fractions (both quark and glue)
- ◆ renormalization of momentum fractions including mixing

x

Slide from Keh-Fei Liu

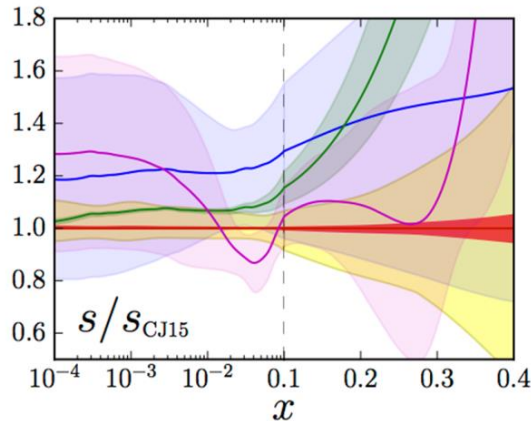
# First Lattice Strange PDF

## § Large uncertainties in global PDFs



∞ Assumptions imposed  
due to lack of precision data

$$s = \bar{s} = \kappa(\bar{u} + \bar{d})$$



- CJ15 ( $T = 10$ )
- CJ15 ( $T = 1$ )
- MMHT14
- HERAPDF1.5
- NNPDF3.0

CTEQ-JLAB <https://www.jlab.org/theory/cj/>

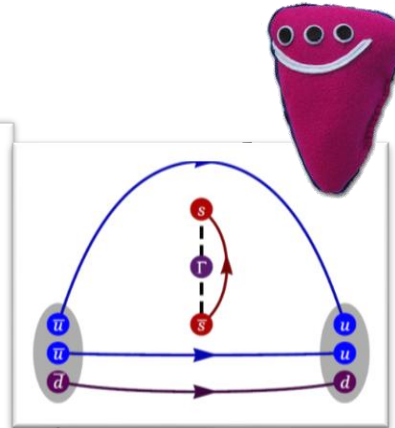
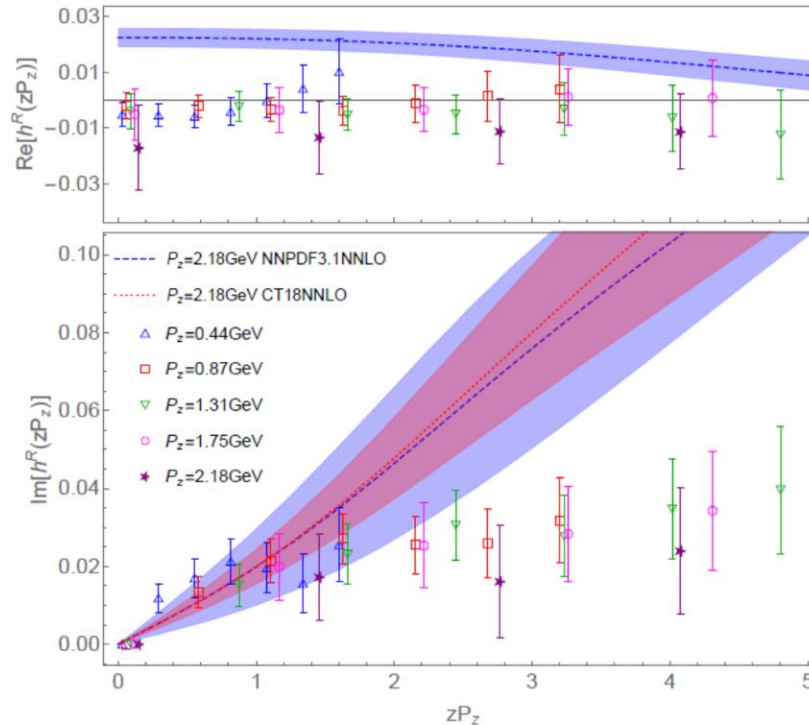
# First Lattice Strange PDF

## § Large uncertainties in global PDFs

$$\text{Re}[h(z)] \propto \int dx (s(x) - \bar{s}(x)) \cos(xzP_z)$$

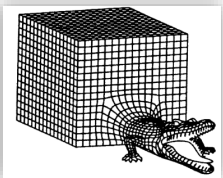
$$\text{Im}[h(z)] \propto \int dx (s(x) + \bar{s}(x)) \sin(xzP_z)$$

- symmetric  $s - \bar{s}$  distribution.
- smaller momentum fraction



G: Rui Zhang

Slide by Rui Zhang @ DNP 2020



# Gluon Helicity

§ Jaffe & Manohar, 1990  $\frac{1}{2} = \frac{1}{2} \Delta\Sigma + \Delta G + \mathcal{L}_q^z + \mathcal{L}_g^z$

§ Can be calculated through large-momentum frame

X. Ji et al., PRL. 111 (2013) 112002; 110 (2013) 262002; PRD 89, 085030 (2014)

$$S_G(P) S_z = \frac{\langle PS | \int d^3x (\vec{E} \times \vec{A}_{\text{phys}})_z | PS \rangle}{2E_P}$$

§ First results by  $\chi$ QCD

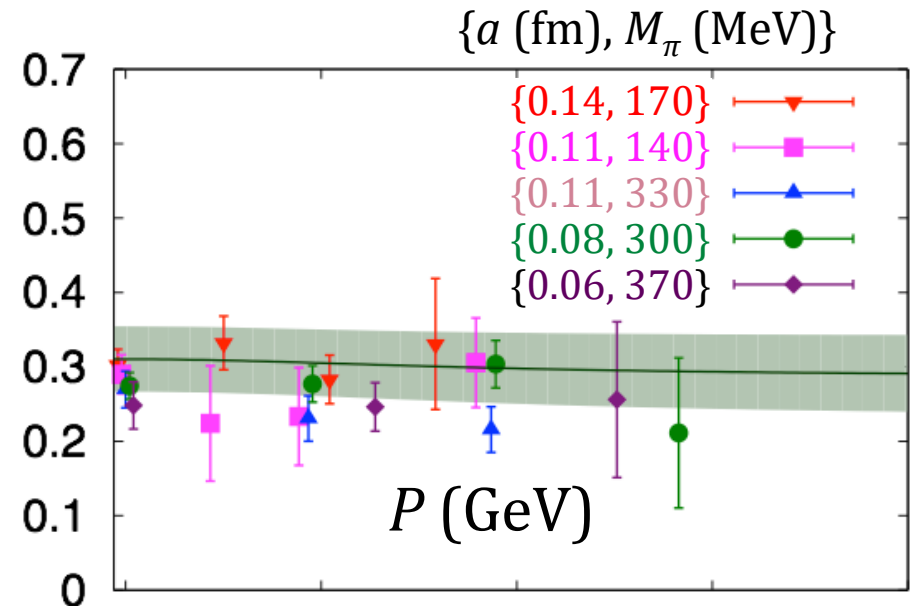
$$\begin{aligned} \Delta G(\mu^2 = 10 \text{ GeV}^2) \\ \approx S_G(\infty, \mu^2 = 10 \text{ GeV}^2) \\ = 0.287(55)(16) \end{aligned}$$

Yang et al, Phys. Rev. Lett. 118 (2017) 102001

∞ Future improvement to matching

§ Current limit

$$\infty \text{ DSSV14 } \int_{0.05}^1 dx \Delta G(10^2 \text{ GeV}, x) \approx [0.14, 0.24]$$

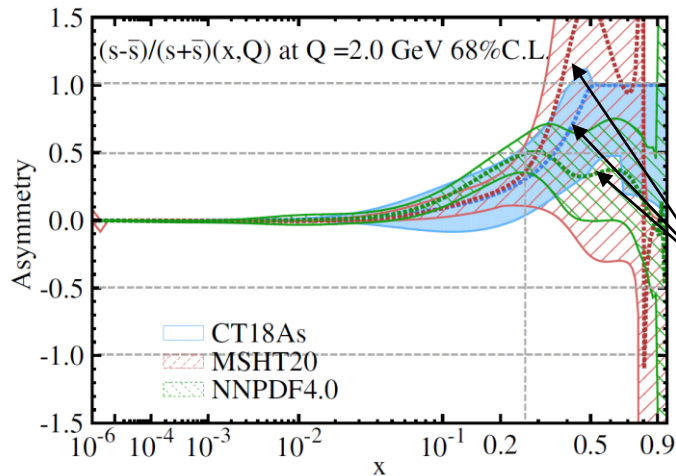




# Lattice Impact on Strange PDF

§ lattice QCD can constrain PDFs (polarized, meson, TMDs, GPDs,...) that are difficult to access in experiments

§ Example: the strangeness asymmetry  $s(x, Q) - \bar{s}(x, Q)$  at  $x > 0.2$  is difficult to measure (left), but can be predicted in lattice QCD (right)

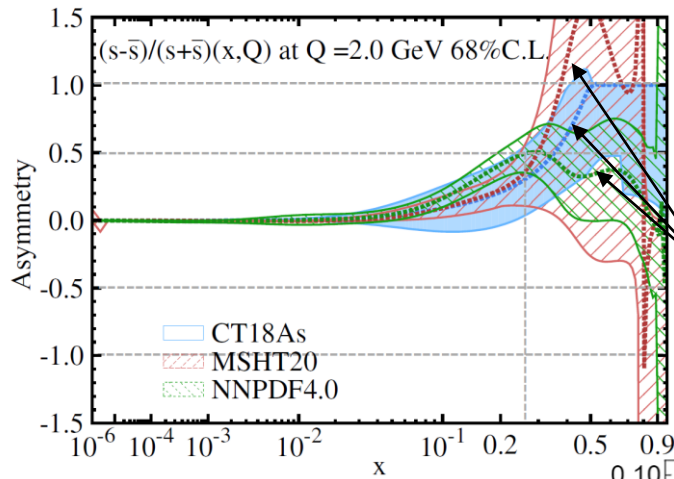


differences reflect the pulls of LHC and other experiments

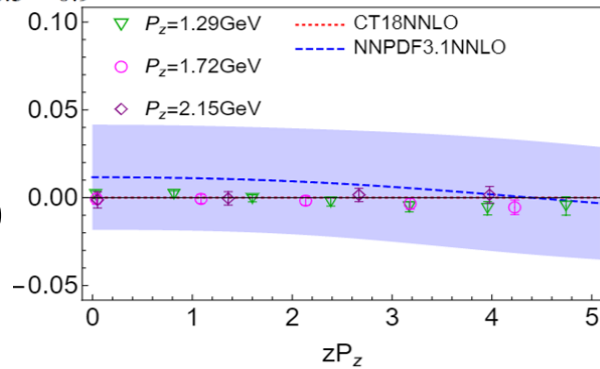
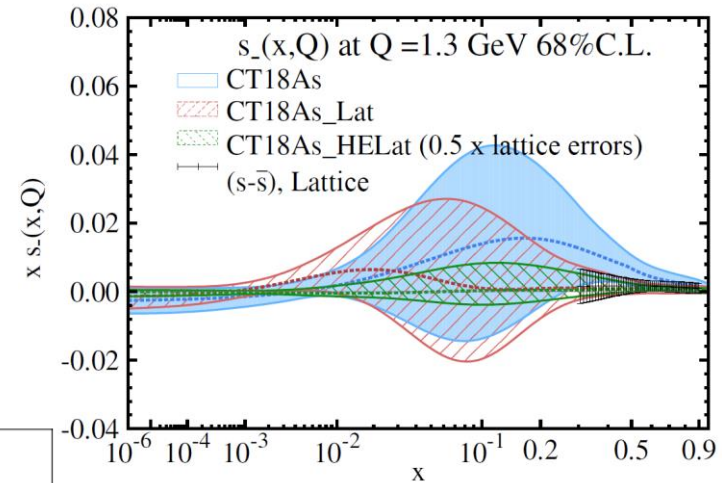
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differences reflect the pulls of LHC and other experiments



$$\text{Re}[h(z)] \propto \int dx (s(x) - \bar{s}(x)) \cos(xzP_z)$$

T. Hou, HL, M. Yan, C.-P. Yuan, 2204.07944